



A Measurement of g_2^p at Low Q^2

Melissa Cummings

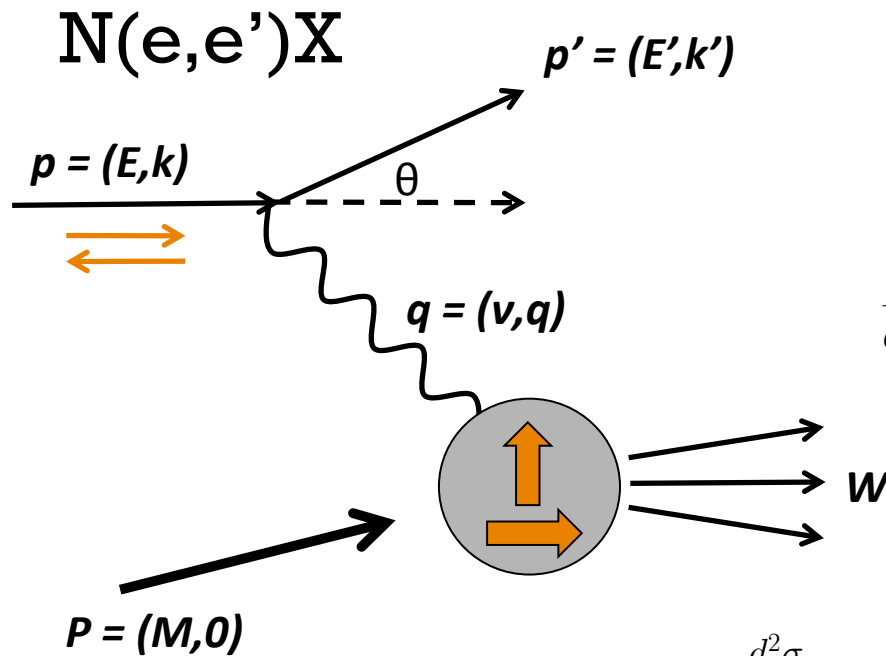
The College of William and Mary

On Behalf of the E08-027 Collaboration

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Inclusive Electron Scattering



Inclusive inelastic
unpolarized cross section:

$$\frac{d^2\sigma}{d\Omega dE'} = \sigma_{Mott} \left[\frac{1}{\nu} \underline{F_2(x, Q^2)} + \frac{2}{M} \underline{F_1(x, Q^2)} \tan^2 \frac{\theta}{2} \right]$$

Inclusive **polarized** cross sections:

$$\frac{d^2\sigma}{dE' d\Omega} (\downarrow\uparrow - \uparrow\uparrow) = \frac{4\alpha^2}{MQ^2} \frac{E'}{\nu E} \left[(E + E' \cos \theta) \underline{g_1(x, Q^2)} - \frac{Q^2}{\nu} \underline{g_2(x, Q^2)} \right]$$

$$\vec{q} = \vec{k} - \vec{k}'$$

$$Q^2 = -q^2 = 4EE' \sin^2 \frac{\theta}{2}$$

$$x = \frac{Q^2}{2M\nu}$$

$$\frac{d^2\sigma}{dE' d\Omega} (\downarrow\Rightarrow - \uparrow\Rightarrow) = \frac{4\alpha^2 \sin \theta}{MQ^2} \frac{E'^2}{\nu^2 E} \left[\nu \underline{g_1(x, Q^2)} + 2E \underline{g_2(x, Q^2)} \right]$$

g_1, g_2 are related to the
spin distribution

Quark-Parton Model

- Bjorken Scaling Limit:

$$\begin{array}{l} Q^2 \rightarrow \infty \\ \nu \rightarrow \infty \end{array} \quad \begin{array}{l} \text{such} \\ \text{that:} \end{array} \quad x = \frac{Q^2}{2M\nu} \quad \text{is finite}$$

- Structure functions can be written in terms of quark distribution functions:

$$F_1(x) = \frac{1}{2} \sum_f z_f^2 [q_f(x) + \bar{q}_f(x)]$$

$$F_2(x) = 2xF_1(x)$$

$$g_1(x) = \frac{1}{2} \sum_f z_f^2 [q_f(x) - \bar{q}_f(x)]$$

No simple
interpretation for g_2

g_2 includes contributions from
quark gluon interactions

What is g_2 ?

(Parton Model Description)

$$g_2(x, Q^2) = g_2^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$$

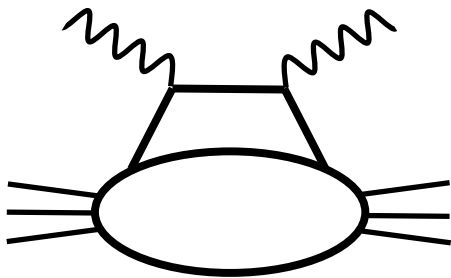
What is g_2 ?

$$g_2(x, Q^2) = g_2^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$$

Wandzura–Wilczek Relation:

$$g_2^{WW}(x, Q^2) = -g_1(x, Q^2) + \int_x^1 \frac{dy}{y} g_1(y, Q^2)$$

Leading twist-2 term



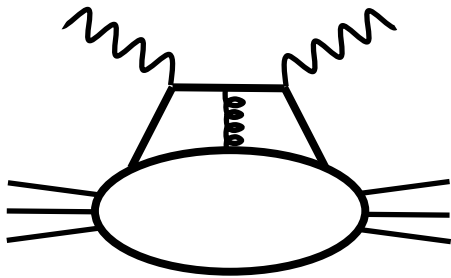
What is g_2 ?

$$g_2(x, Q^2) = g_2^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$$

$$\int_x^1 \frac{\partial}{\partial y} \left[\frac{m_q}{M} h_T(y, Q^2) + \zeta(y, Q^2) \right] \frac{dy}{y}$$

h_T : Arises from quark transverse polarization distribution

ζ : Arises from quark-gluon interactions (twist-3)



Measurements of g_2 and its Moments

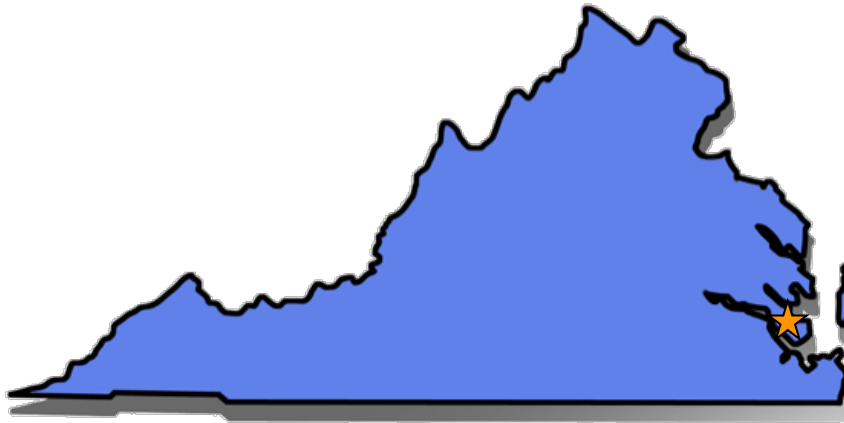
Measurements of g_2 require a transversely polarized target – more difficult experimentally

- 0th moment (no x -weighting): Burkhardt-Cottingham Sum Rule
 - Valid at all Q^2

$$\int_0^1 g_2(x, Q^2) dx = 0$$

- 2nd moment (x^2 weighting):
 - High Q^2 – d_2 , twist-3 color polarizability, test of lattice QCD
 - Low Q^2 – spin polarizabilities, test of χ PT

Measurements of g_2 and its Moments

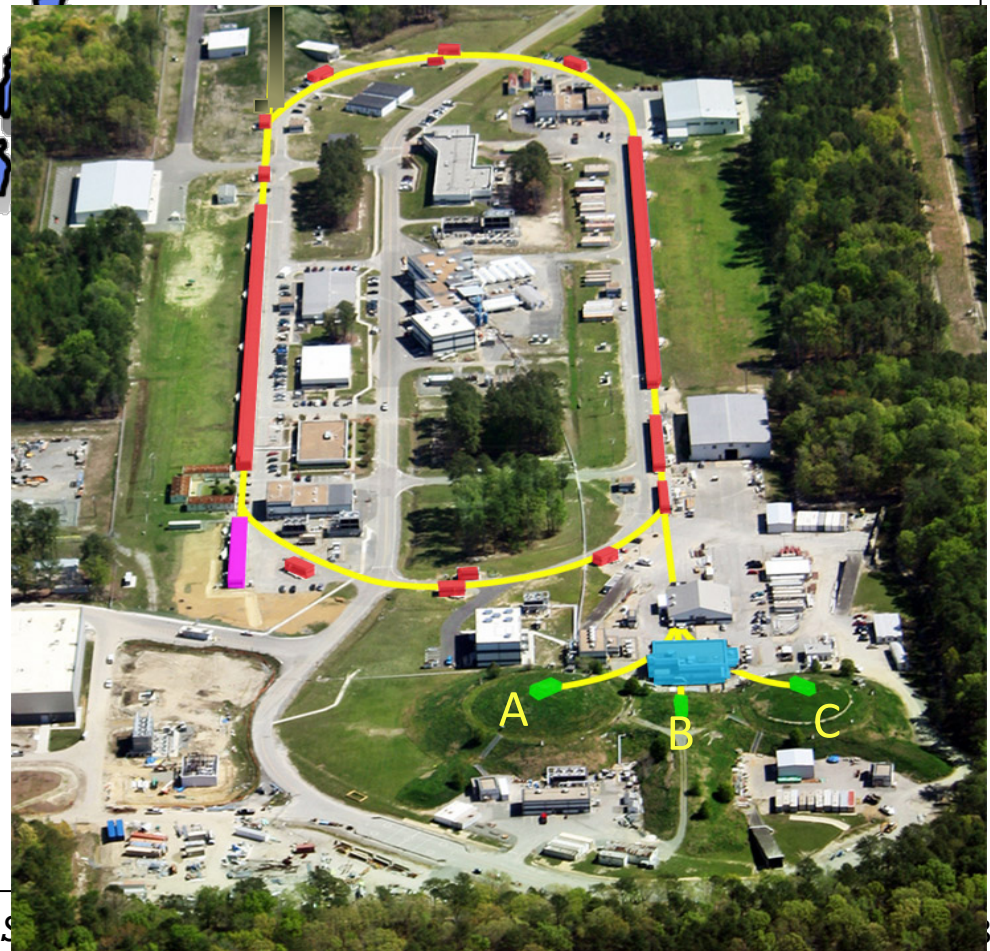


Jefferson Lab

CEBAF

- High intensity electron accelerator based on CW SRF technology
- $E_{\text{max}} = 6 \text{ GeV}$
- $I_{\text{max}} = 200 \mu\text{A}$
- $\text{Pol}_{\text{max}} = 85\%$

****Recently upgraded to 12 GeV****



Measurements of g_2 and its Moments

- Prior to measurements at JLab, first dedicated experiment was SLAC E155x
- g_2 Measurements on the **neutron** at JLab:
 - E97-103: $W > 2$ GeV, $Q^2 \approx 1$ GeV², $x \approx 0.2$, study higher twist (published)
 - E99-117: $W > 2$ GeV, high Q^2 (3-5 GeV²) (published)
 - E94-010: moments at low Q^2 (0.1-1 GeV²) (published)
 - E97-110: moments at very low Q^2 (0.02-0.3 GeV²) (analysis)
 - E01-012: moments at intermediate Q^2 (1-4 GeV²) (submitted)
 - E06-014: moments at high Q^2 (2-6 GeV²) (analysis)

Measurements on the **proton** at JLab:

- RSS: moments at intermediate Q^2 (1-2 GeV²) (published)
- SANE: moments at high Q^2 (2-6 GeV²) (analysis)
- • E08-027 (g_2^p): moments at very low Q^2 (0.02-0.2 GeV²) (analysis)

0th Moment of g_2

$$\Gamma_2 = \int_0^1 g_2(x, Q^2) dx = 0$$

Brown: SLAC E155x

Red: Hall C RSS

Black: Hall A E94-010

Green: Hall A E97-110 (preliminary)

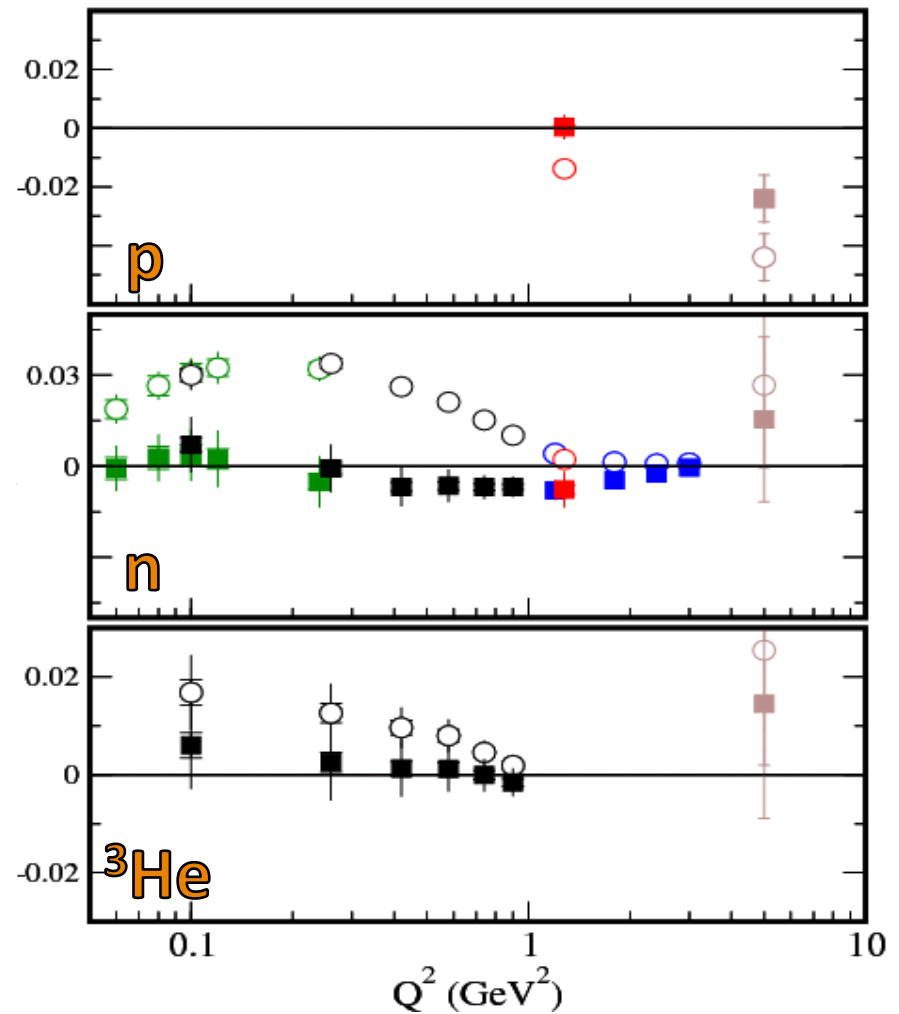
Blue: Hall A E01-012 (submitted)

BC Sum = Measured + Low x + Elastic

Measured: open circles

Low x: unmeasured low-x part of the integral – assume leading twist behavior

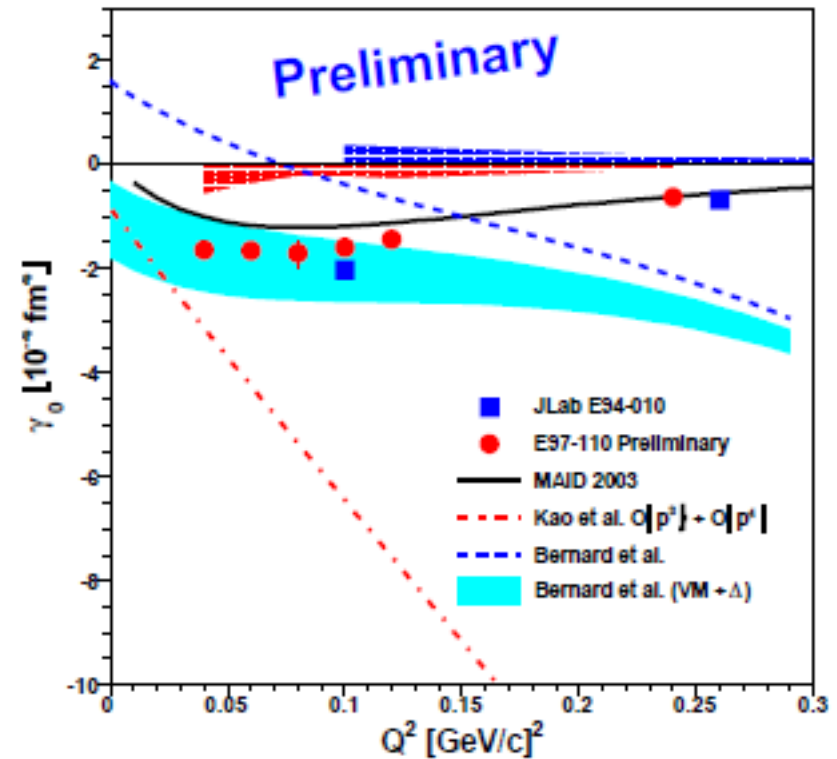
Elastic: obtained from well known Form Factors



2nd Moment: Spin Polarizabilities

(V. Sulkosky)

- Generalized spin polarizabilities γ_0 and δ_{LT} are a benchmark test of χ PT
- Difficulty is how to include the nucleon resonance contributions
- γ_0 is sensitive to resonances, δ_{LT} is not
- Neutron results for γ_0

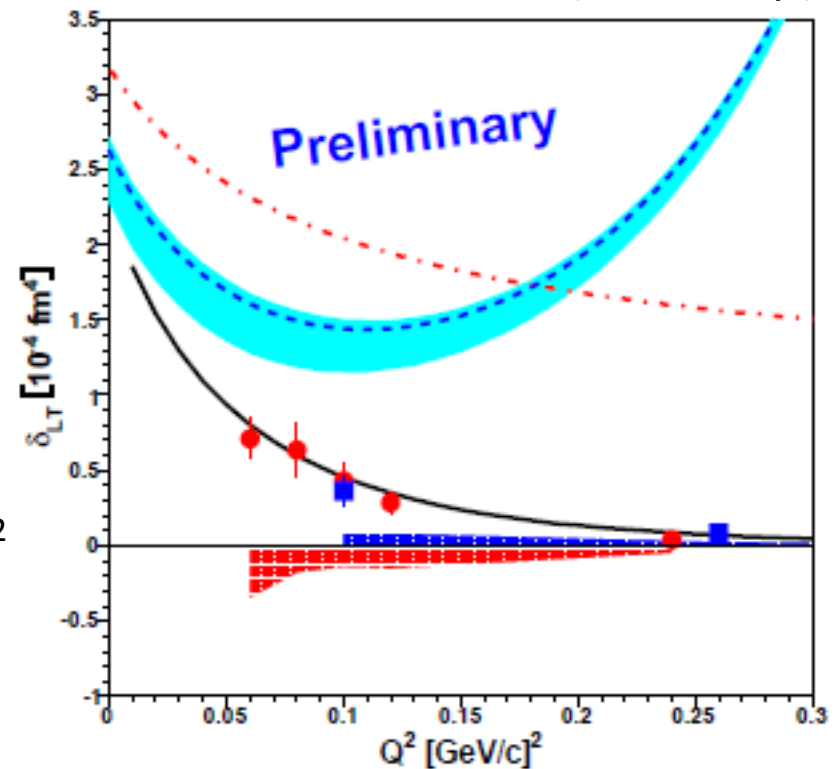


$$\gamma_0(Q^2) = \frac{16\alpha M^2}{Q^6} \int_0^{x_0} x^2 \left[g_1(x, Q^2) - \frac{4M^2}{Q^2} x^2 g_2(x, Q^2) \right] dx$$

2nd Moment: Spin Polarizabilities

(V. Sulkosky)

- Neutron results for δ_{LT}
- δ_{LT} is seen as a more suitable testing ground – insensitive to Δ -resonance
- Data is in significant disagreement with χ PT calculations
- g_2^p experiment will provide test at low Q^2



$$\delta_{LT}(Q^2) = \frac{16\alpha M^2}{Q^6} \int_0^{x_0} x^2 [g_1(x, Q^2) + g_2(x, Q^2)] dx$$

- JLab E94-010
- E97-110 Preliminary
- MAID 2003
- - - Kao et al. $|O|p^2 + O|p^4|$
- - - Bernard et al.
- Bernard et al. (VM + Δ)

Finite Size Effects

Hyperfine Splitting of Hydrogen

- Splitting is defined in terms of Fermi Energy E_f

$$\Delta_E = (1 + \delta) E_F$$

Where:

$$\delta = 1 + (\delta_{QED} + \delta_R + \delta_{small}) + \Delta_S$$

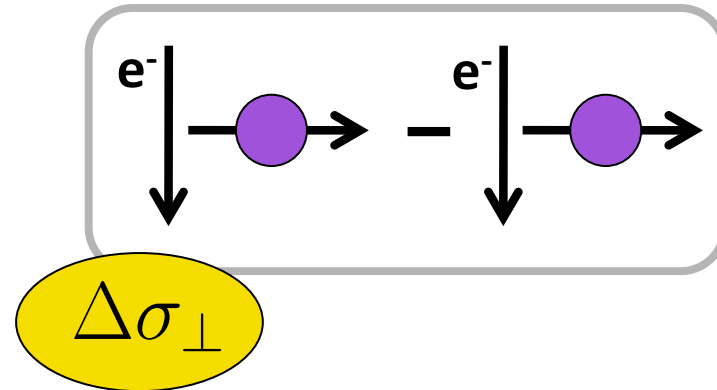
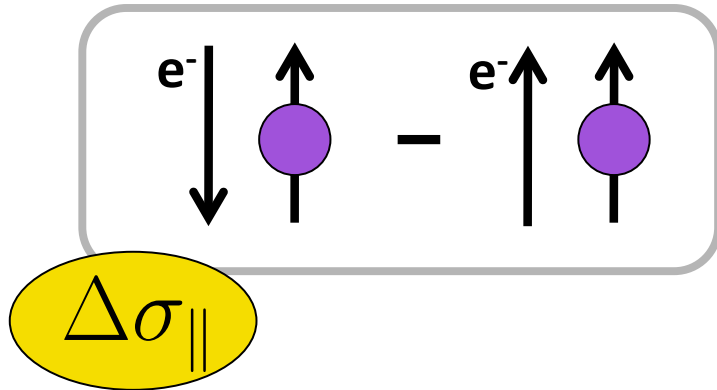
Correction for proton structure – contains contribution from g_2 (dominated by g_2 at low Q^2)

- Proton Charge Radius
 - Results from μ P disagrees with eP scattering result by $\sim 7\sigma$
 - Main uncertainties arise from proton polarizability and differing values of the Zemach radius

g_2^p Experiment at JLab (E08-027)

- Will provide the first measurement of g_2 for the proton at low to moderate Q^2
- Will provide insight on several outstanding physics puzzles:
 - BC sum rule
 - Discrepancy suggested for high- Q^2 data
 - δ_{LT} polarizability
 - χ PT calculations do not match data
 - Finite size effects:
 - Hydrogen hyperfine splitting: proton structure contributes to uncertainty
 - Proton charge radius: proton polarizability contributes to uncertainty
- Data was taken in Hall A in 2012 – analysis is currently underway

Experimental Technique



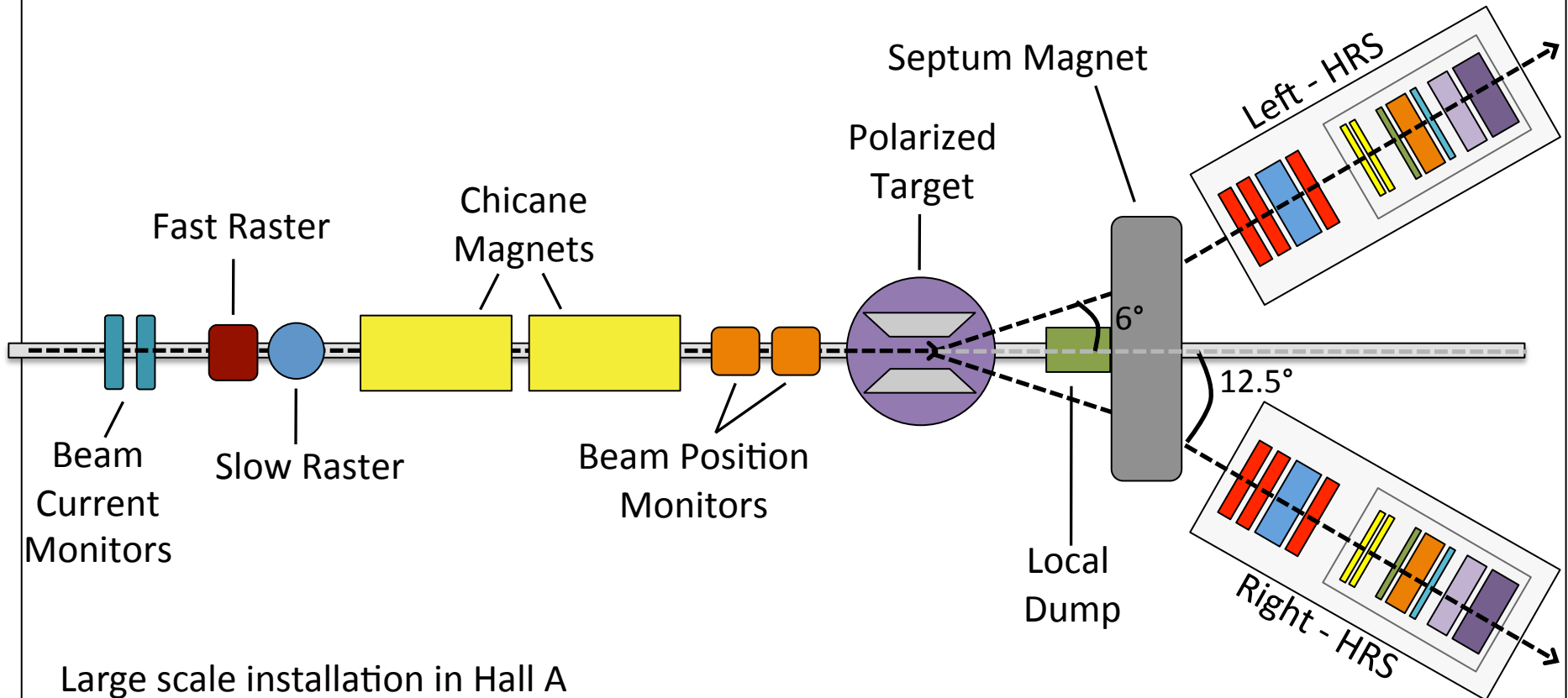
$$\frac{d^2\sigma}{dE'd\Omega}(\downarrow\uparrow - \uparrow\uparrow) = \frac{4\alpha^2 \sin\theta}{MQ^2} \frac{E'}{\nu E} \left[(E + E' \cos\theta) g_1(x, Q^2) - \frac{Q^2}{\nu} g_2(x, Q^2) \right]$$

$\Delta\sigma_{||}$ measured during EG4 experiment in Hall B: will extract g_1^p at low Q^2

$$\frac{d^2\sigma}{dE'd\Omega}(\downarrow\Rightarrow - \uparrow\Rightarrow) = \frac{4\alpha^2 \sin\theta}{MQ^2} \frac{E'^2}{\nu E} [\nu g_1(x, Q^2) + 2E g_2(x, Q^2)]$$

$\Delta\sigma_{\perp}$ obtained from g_2^p experiment

Experimental Setup

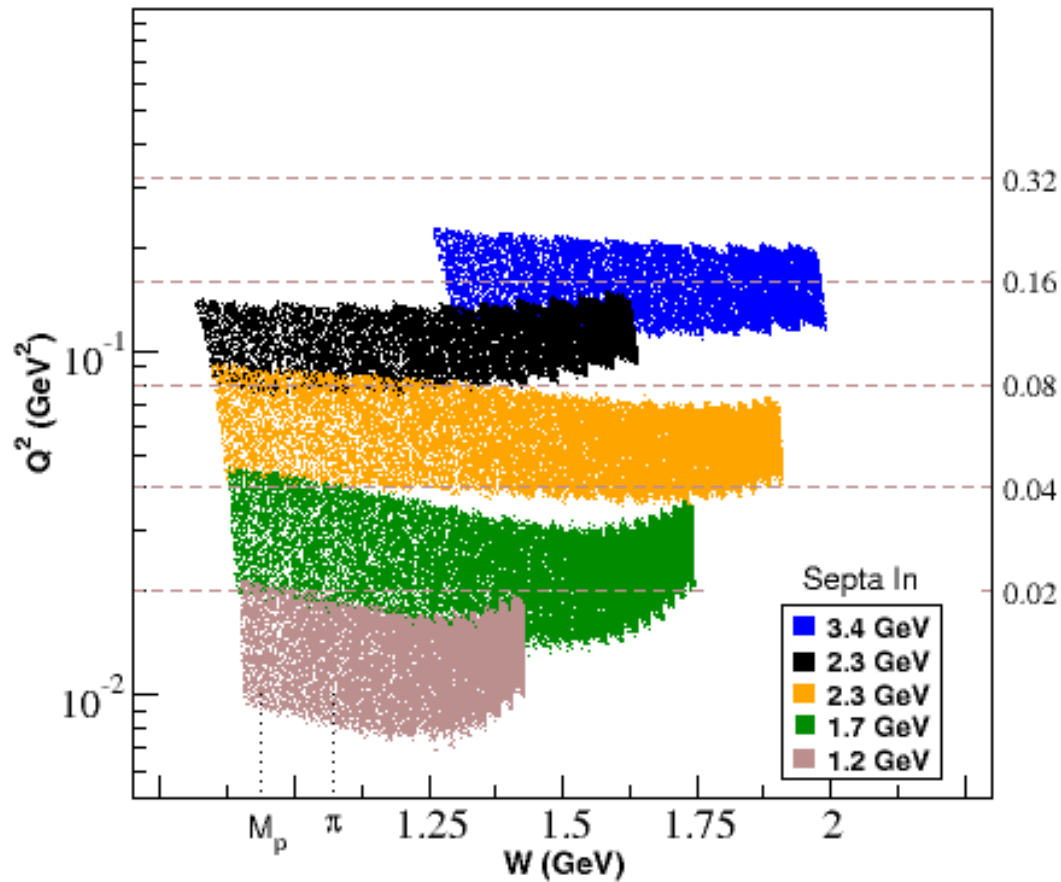


Large scale installation in Hall A

- DNP NH_3 target with 2.5/5 T magnetic field (longitudinal and transverse configurations)
- New beamline diagnostics for low current (<100 nA) running
- Chicane and septum magnets
- Local dump

Kinematic Coverage

First data on g_2 for proton at low Q^2



$W < 2 \text{ GeV}$
 $0.02 < Q^2 < 0.2 \text{ GeV}^2$

Beam Energy (GeV)	Target Field (T)
2.2	2.5
1.7	2.5
1.1	2.5
2.2	5.0
3.3	5.0

Status of Analysis

Completed

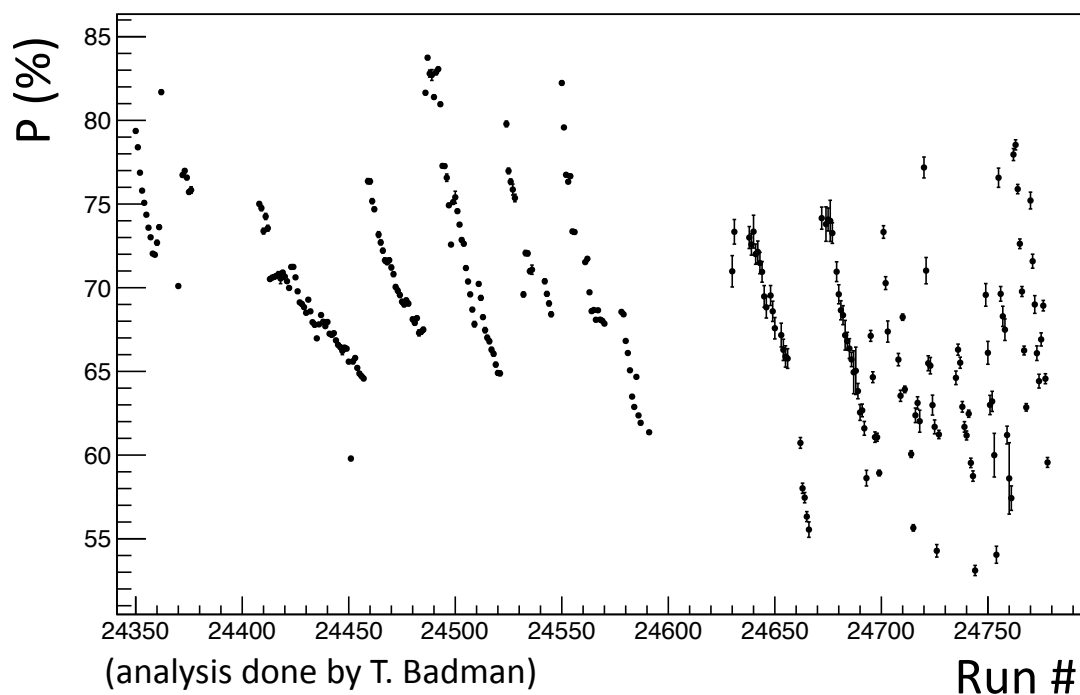
- Run Database
- HRS Optics
 - Field measurement analysis
 - VDC t_0 calibration
 - Simulation Package
 - Optics with target field (LHRS)
- Detector Calibrations/Efficiency Studies
 - Gas Cherenkov
 - Lead Glass Calorimeters
 - Scintillator trigger efficiencies
 - Multi-track analysis
- Scalers
 - BCM calibration
 - Helicity decoding
 - Dead time calculations
- Target Polarization Analysis
- BPM Calibrations

In Progress

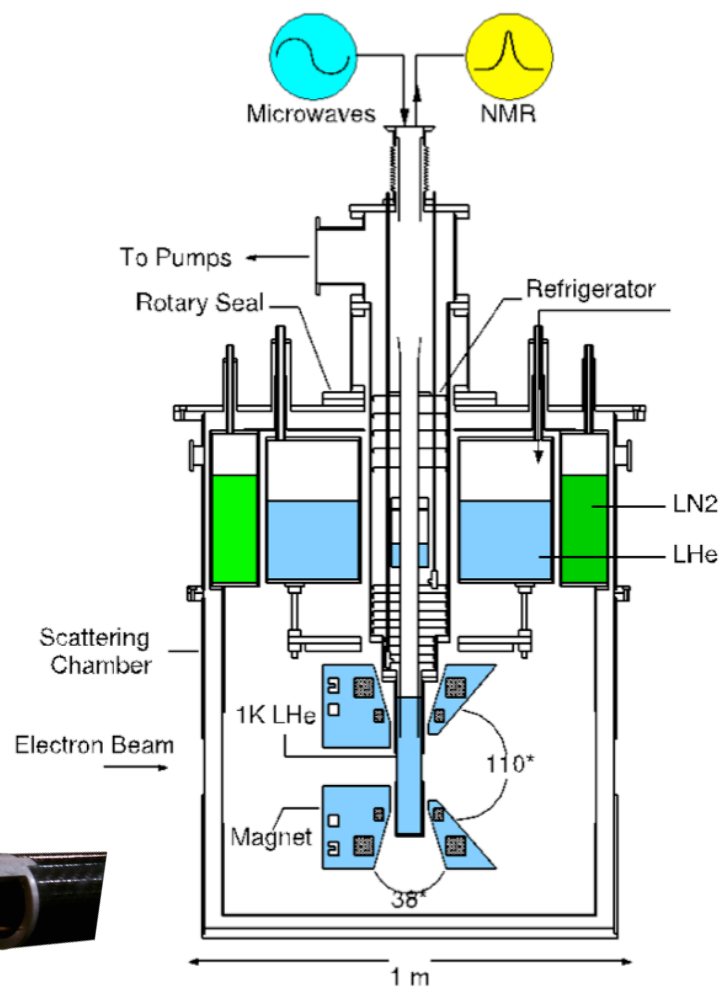
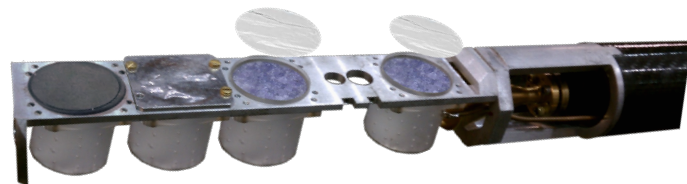
- Raster Size Calibrations
- Packing Fraction/Dilution Analysis
- Elastic Analysis
- Yields/Radiative Corrections

Polarized NH₃ Target

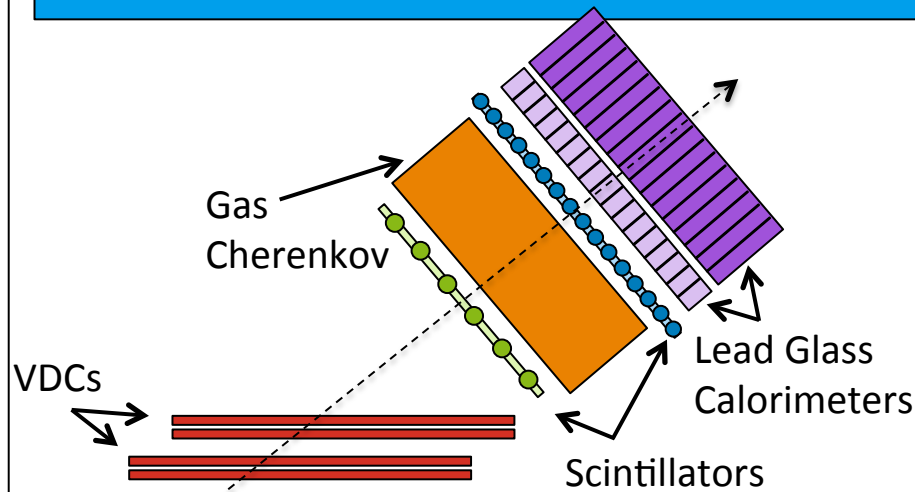
Target Polarization Results for 5T Field Setting



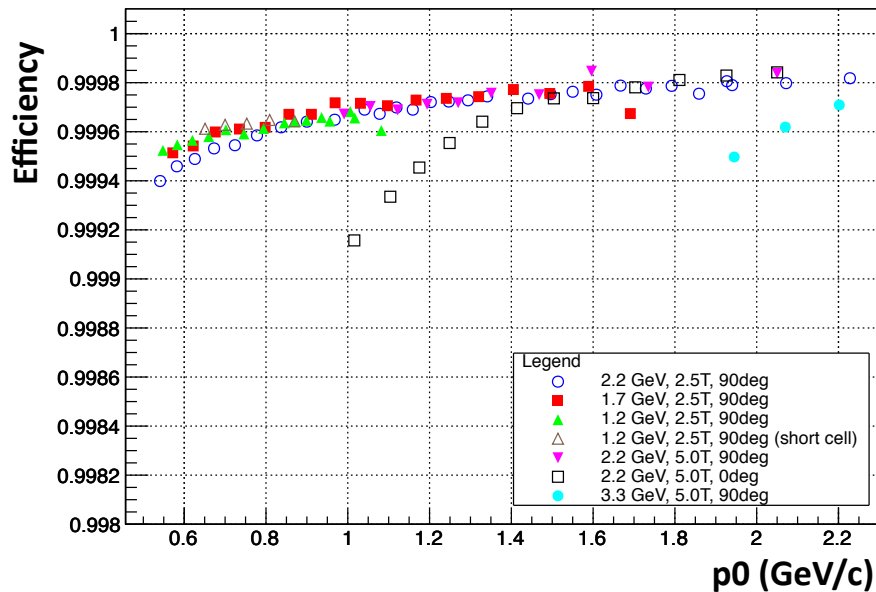
Average Polarization:
5T: ~70%
2.5T: ~15%



Detector Stack



Gas Cherenkov Detector Efficiency

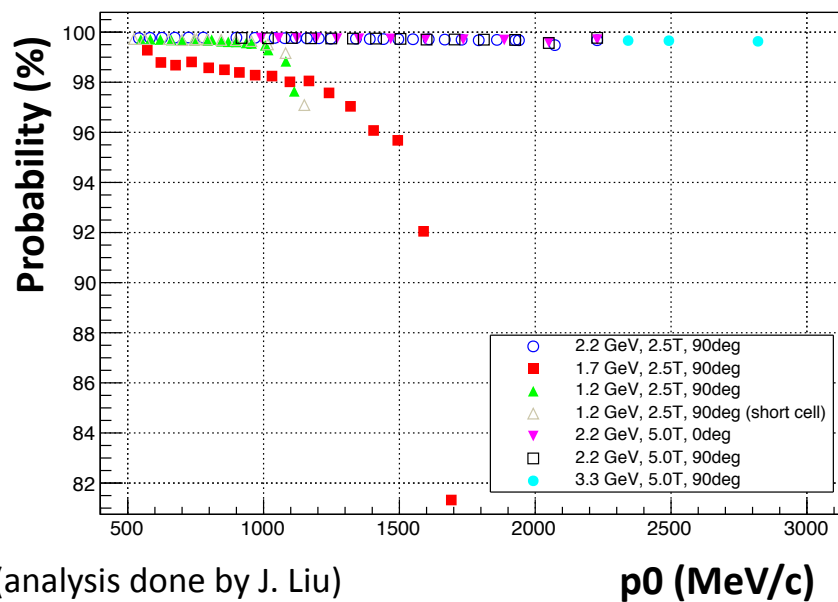


- Gas Cherenkov and lead glass electromagnetic calorimeters used for particle identification
- High Efficiency ($>99\%$) for gas Cherenkov and lead glass calorimeters
- Two planes of scintillators used to create singles trigger
- High Trigger efficiency ($>98\%$) for production runs

VDC Multi-track Analysis

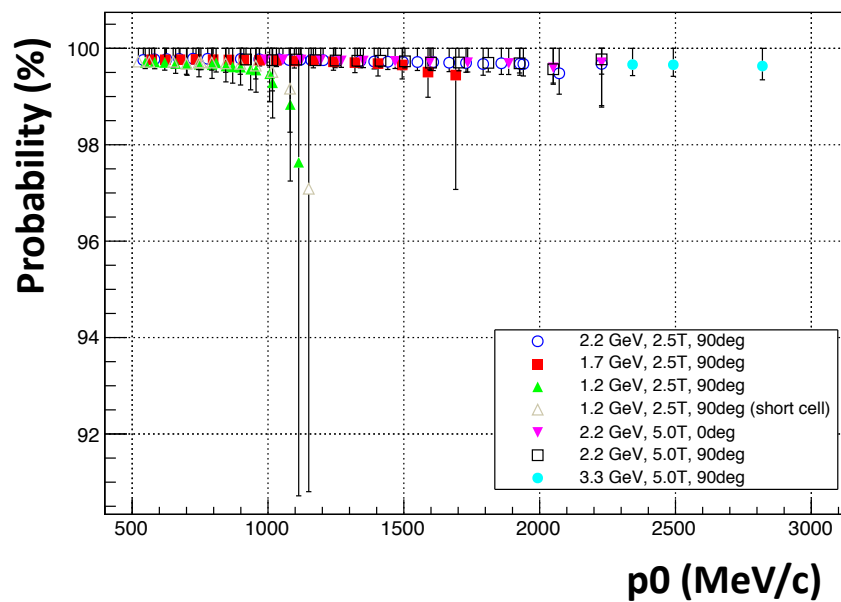
- VDC efficiency is the fraction of events with a successful track reconstruction
- Number of multi-track events can reach $\sim 20\%$ for some kinematic settings
- Cluster energy of tracks is examined to determine inefficiency

VDC Single-Track Probability



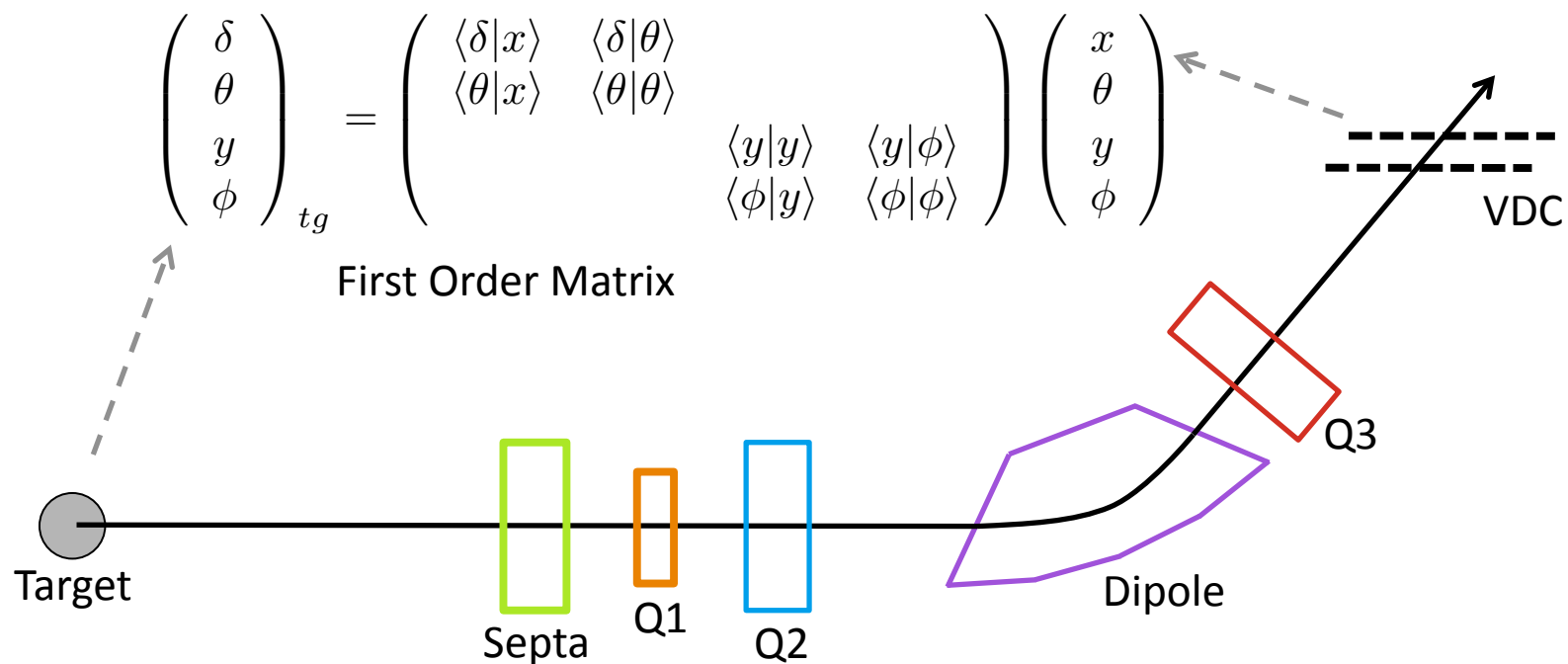
(analysis done by J. Liu)

VDC Total Good Track Probability



HRS Optics

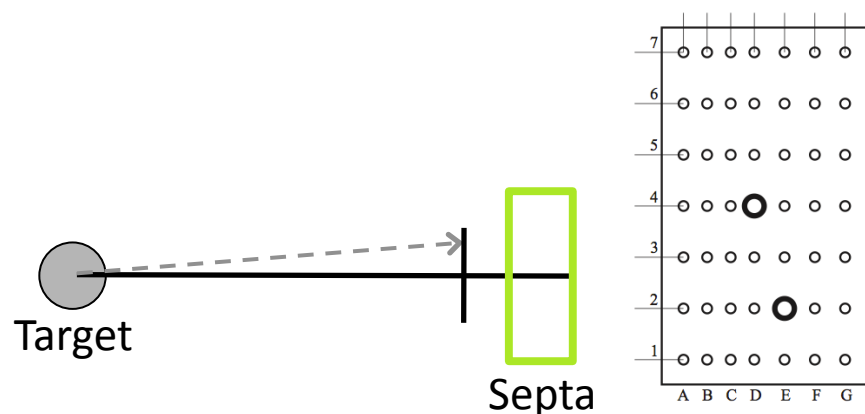
- HRS has a series 3 quadrupoles and 1 dipole for focusing and dispersion
- g_2^p experiment also used an additional septum magnet
- Optics study provides a matrix to transform VDC readouts to kinematic variables



HRS Optics

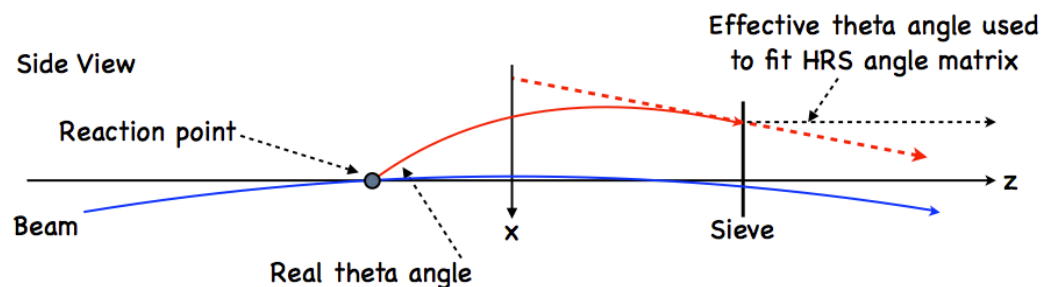
Without Target Field

- Calibrations done using a foil target and point beam
- Can use sieve slit to get the real scattering angle from geometry



With Target Field

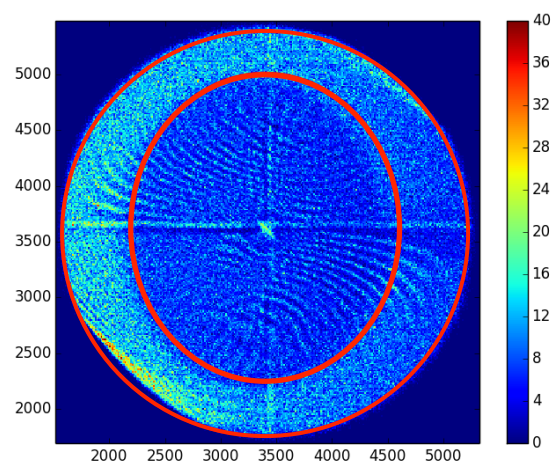
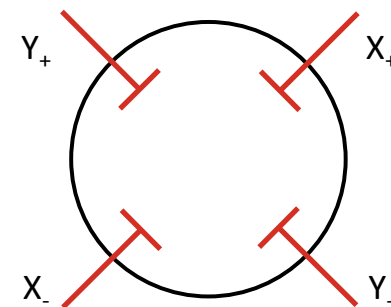
- Must split reconstruction into 2 parts:
 - Can use HRS transform matrix for reconstruction from VDC to sieve slit
 - Then use a target field map to do a ray trace of the scattered particle from the sieve slit to the target



Beam Position Calibrations

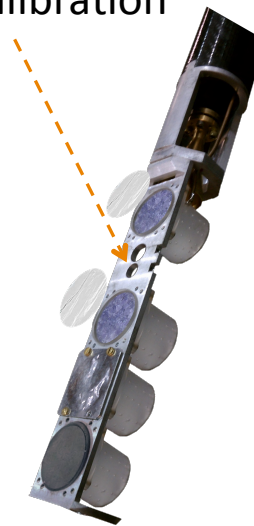
- 4 antennas used to determine the beam position
- Fast and Slow raster system used to reduce radiation damage and minimize depolarization of target
- BPMs can only give average beam position, need raster information to get position event-by-event

$$X = X_{bpm} + X_{raster_{fast}} + X_{raster_{slow}}$$



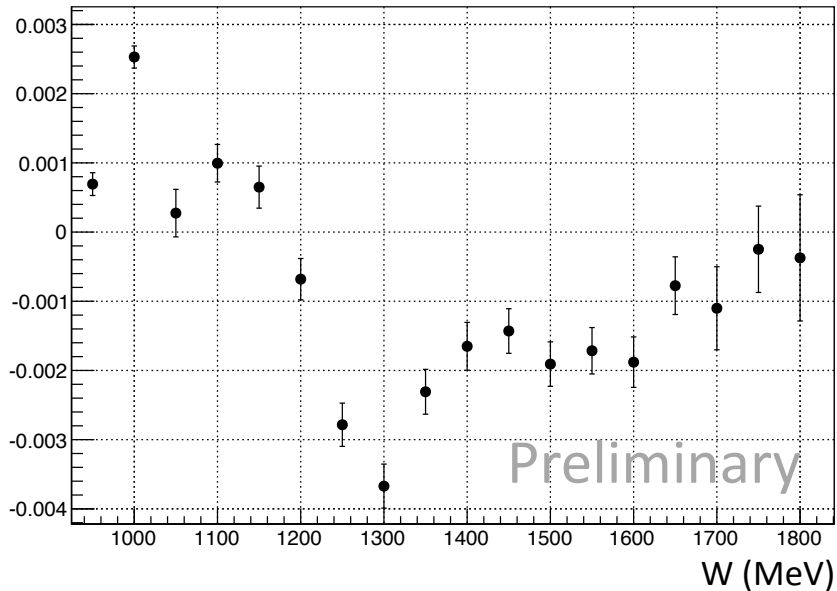
(analysis done by P. Zhu)

hole used for calibration

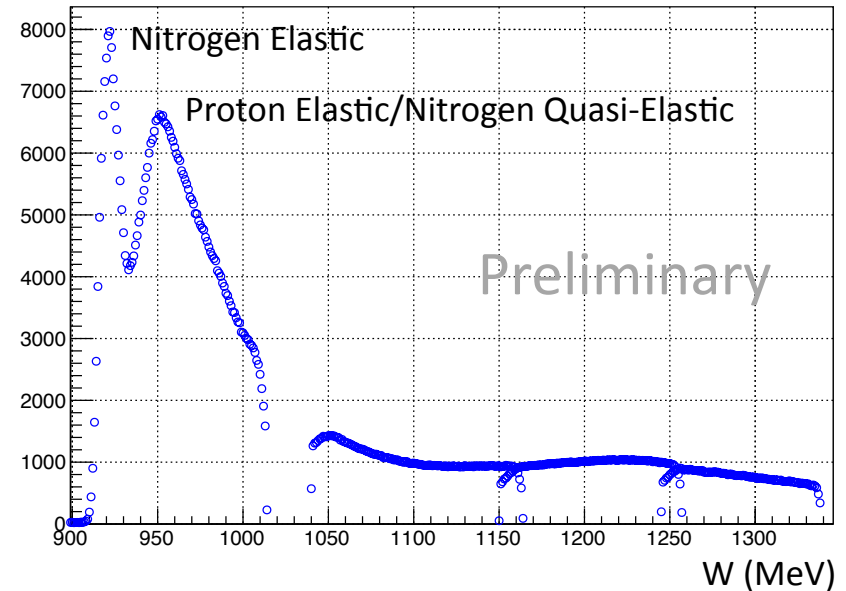


Preliminary Results

Asymmetry



Yield



$$A_{\perp} = \left(\frac{1}{P_b P_t} \right) \frac{Y_+ - Y_-}{Y_+ + Y_-}$$

$$Y_{\pm} = \frac{N_{\pm}}{Q_{\pm} L T_{\pm}}$$

$$\Delta\sigma_{\perp} = \sigma_{total} A_{\perp}$$

Summary of g_2^p

g_2^p experiment will provide first precision measurement for proton at low Q^2
 $0.02 < Q^2 < 0.2 \text{ GeV}^2$

- Will provide insight on several outstanding physics puzzles
 - BC Sum Rule: Violation suggested for proton at large Q^2 (SLAC E155x)
 - Longitudinal-transverse spin polarizability: benchmark test of χ PT, discrepancy seen for neutron data
 - Hydrogen hyperfine splitting: correction for proton structure contributes to uncertainty
 - Proton charge radius: contributions to uncertainty include proton polarizability