

# SBS Optics and Spin Transport

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SBS Weekly Meeting

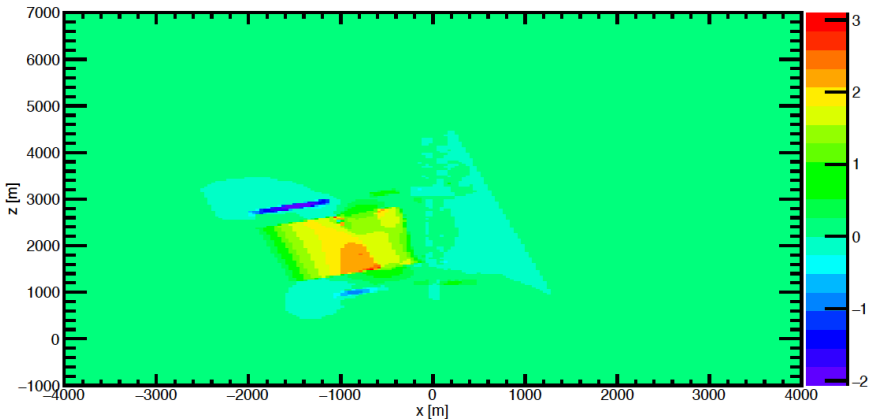
Jan. 27, 2016

# Outline

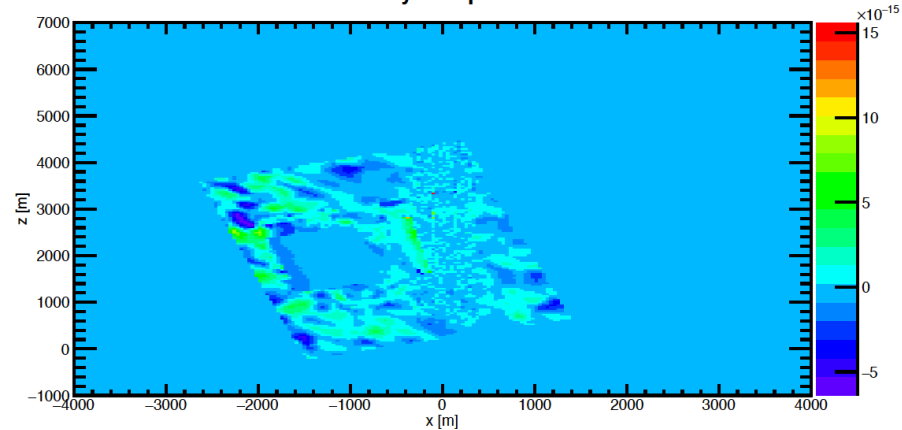
- SBS/BigBite Optics studies from GEANT4
- Overview of previous studies in GEANT4
- Fitting SBS optics matrix—”reconstruction” coefficients
- SBS Spin Transport in GEANT4
  - Formalism—BMT equation
  - Ideal dipole approximation
  - Approach for fitting—added technical challenges compared to optics
- Current status and plans

# SBS magnetic field layout from TOSCA

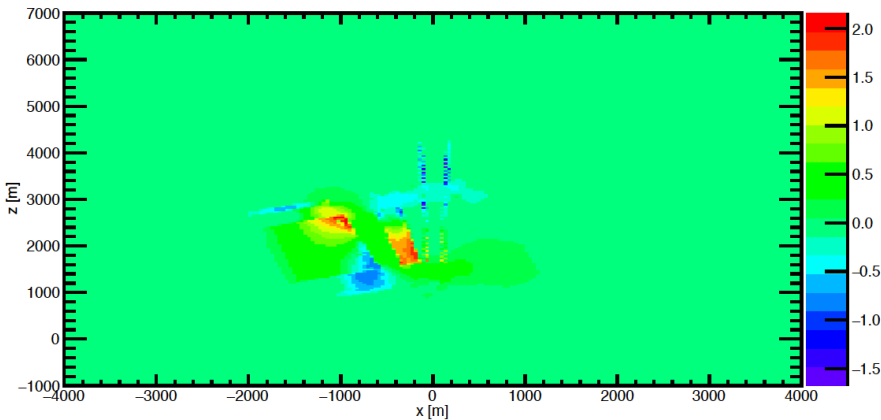
Field x component



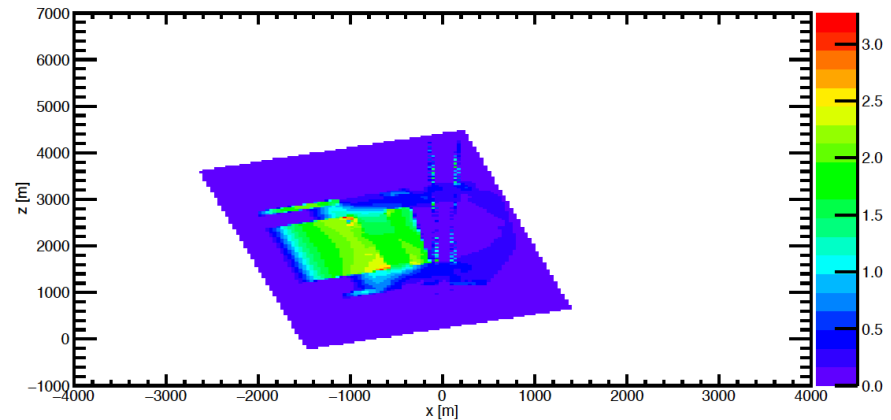
Field y component



Field z component



Field total magnitude



- SBS field components calculated by TOSCA, projected on the xz plane, in Tesla.
  - y component in the xz plane (at  $y=0$ ) is zero by symmetry
  - Includes field clamps and beamline magnetic shielding
- Field is predominantly horizontal and perpendicular to SBS axis in dipole gap
- Good approximation to a pure dipole field

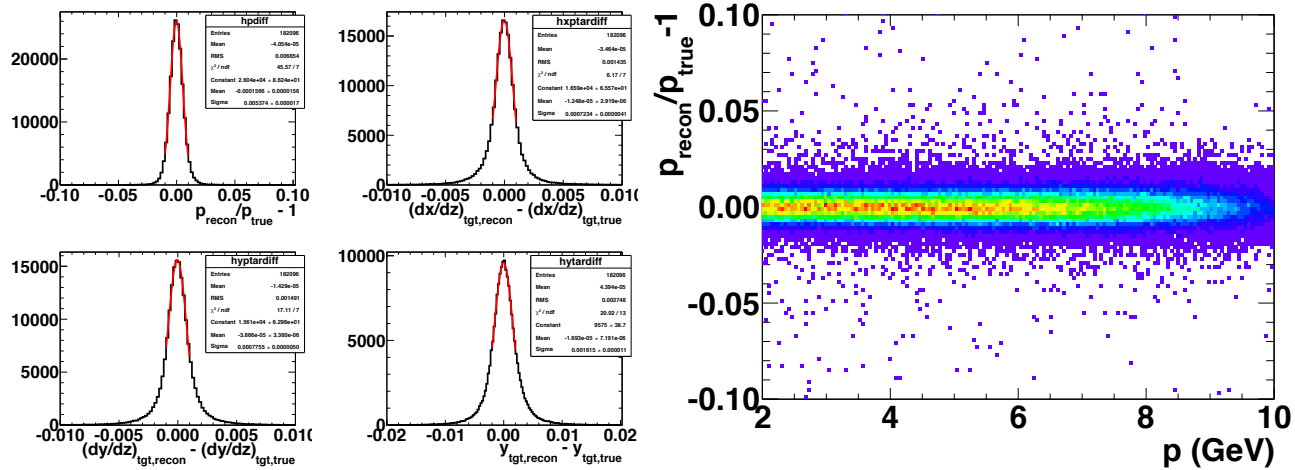
# BigBite and SBS optics from GEANT4

- Charged particles are traced through magnetic field layout of BigBite/SBS in GEANT4 using classical 4<sup>th</sup>-order Runge-Kutta numerical integration of the equation of motion.
- Charged-particles deposit energy in GEMs, making “hits”—GEM “hits” are then smeared by a Gaussian with a  $\sigma$  of 70  $\mu\text{m}$ , representing the coordinate resolution of GEMs, and a straight-line track is fitted.
- In MC, we know the track parameters at the target and at the “focal plane” (GEM location).
- We expand the reverse transport matrix in a power series in the measured track parameters:

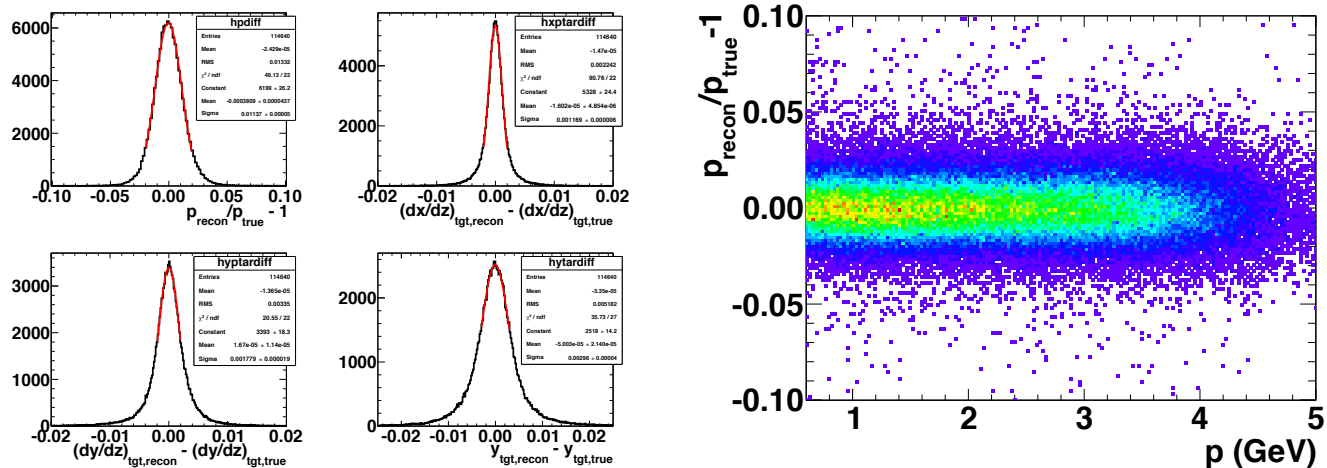
$$(x'_{tgt}, y'_{tgt}, y_{tgt}, 1/p) = \sum_{i+j+k+l+m \leq 6} C_{x',y',y,1/p}^{ijklm} x_{fp}^i y_{fp}^j x'_{fp}{}^k y'_{fp}{}^l x_{tgt}^m$$

- Track parameters at the target are *linear* functions of the expansion coefficients
  - Use standard linear algebra libraries, e.g., Singular Value Decomposition (SVD), to fit the coefficients, avoid pitfalls of nonlinear fitting/numerical minimization/Minuit
- **Why fit 1/p instead of p or the traditional  $\delta = 100 \times (p/p_0 - 1)$  in the expansion?**
  - **SBS/BigBite are large-acceptance spectrometers—range of “delta” can equal or exceed  $\pm 100\%$ , it is no longer a good expansion variable.**
  - **SBS/BigBite are non-focusing, dipole spectrometers  $\rightarrow$  an expansion in 1/p converges very quickly— $x_{fp}, x'_{fp}$  are almost linear in 1/p for dipole magnets**

# SBS and BigBite Optics/Resolution from g4sbs—Old results

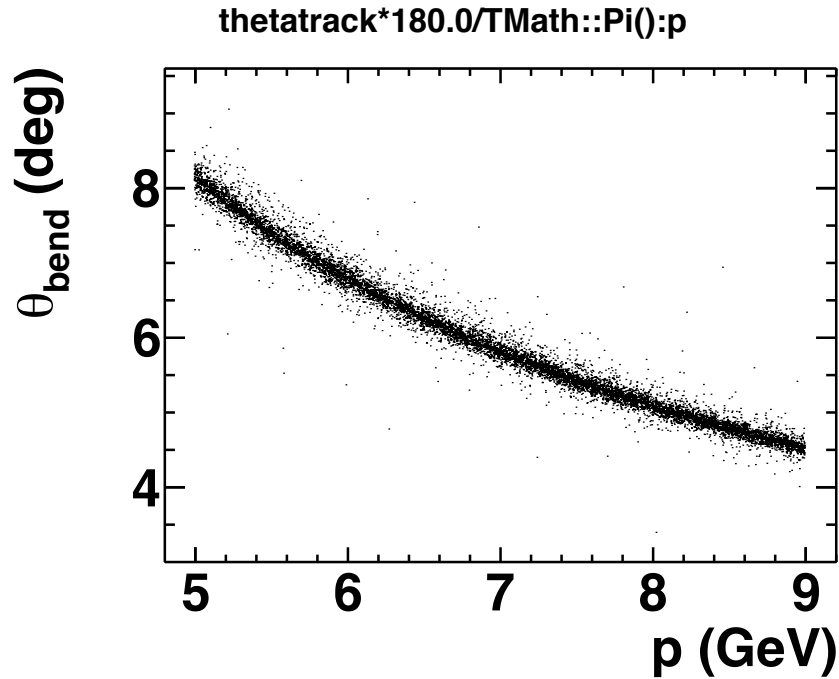


SBS angle, vertex and momentum resolution for 1.4-Tesla uniform field,  $\sigma_p/p \sim 0.5\%$   
(average for 2-10 GeV pions)

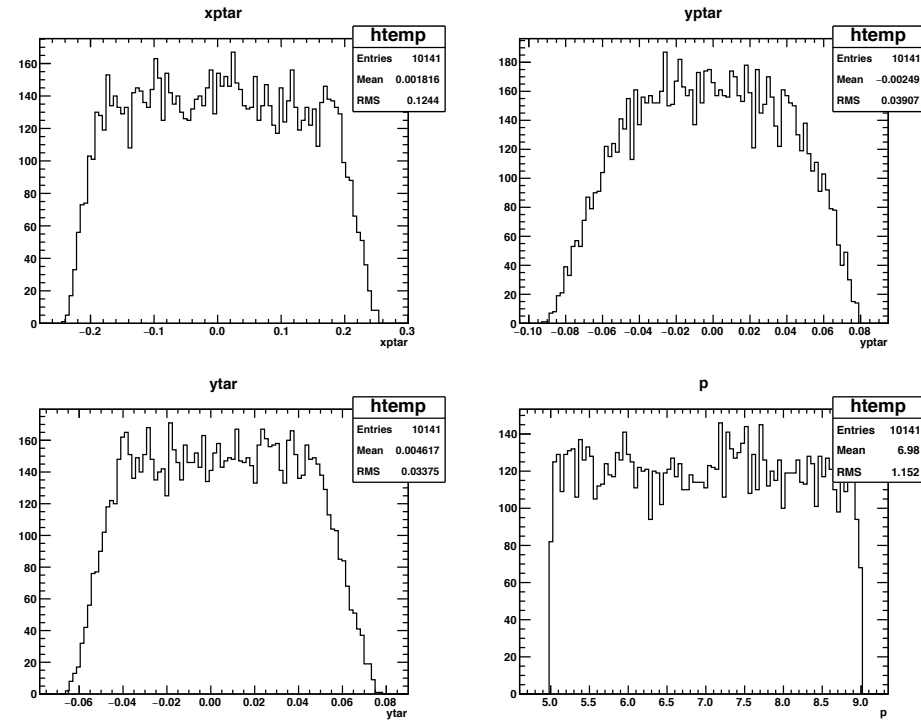


BigBite angle, vertex and momentum resolution for "map\_696A.dat",  $\sigma_p/p \sim 1.1\%$

# Optics and Spin Transport Studies for GEP, $Q^2 = 12 \text{ GeV}^2$



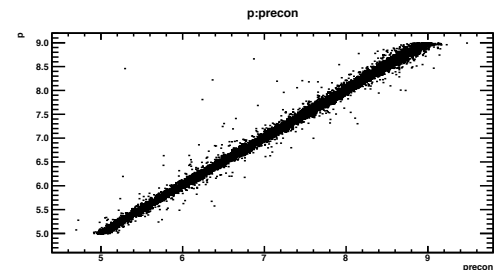
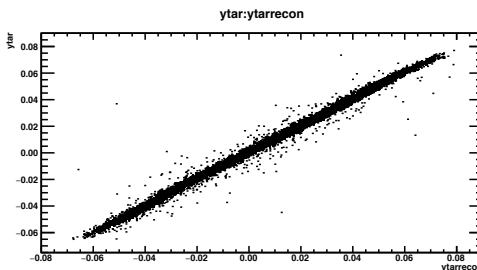
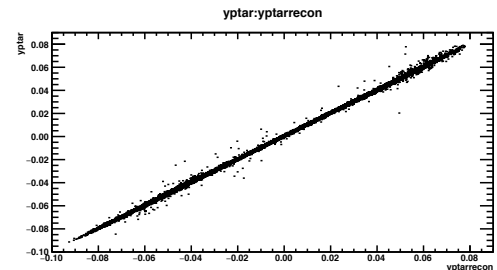
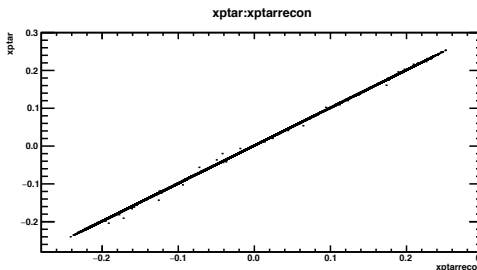
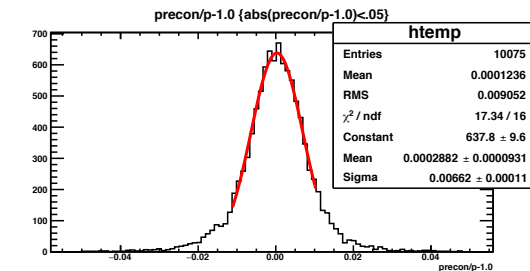
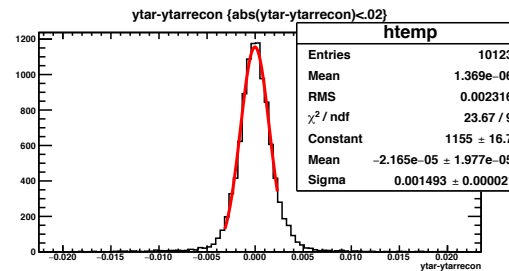
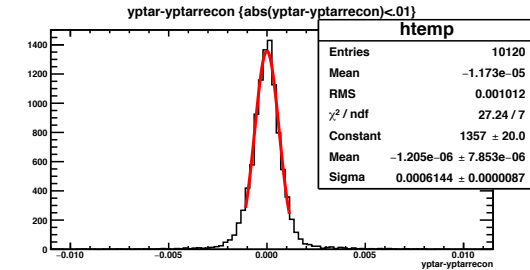
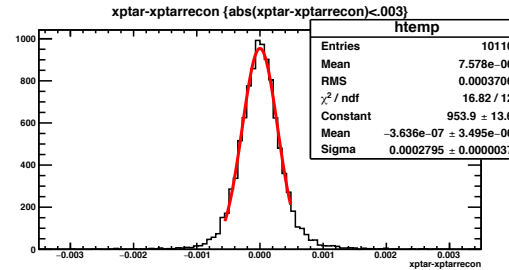
Track vertical bend angle vs. momentum



SBS angular and vertex acceptance for GEP  
highest  $Q^2$

- GEANT4 simulation for optics and spin transport:
  - Use “particle gun” generator with limits chosen wide enough to populate full acceptance of SBS (use 40 cm target)
  - Proton momenta generated in the range of 5-9 GeV (corresponding to highest  $Q^2$  of GEP)
  - Generate 10,000 protons in three different starting spin orientations in the fixed TRANSPORT coordinate system:
    - Pure “X” (vertically down)
    - Pure “Y” (horizontal, toward small angle)
    - Pure “Z” (along SBS central ray)
  - Fit reconstruction coefficients and spin transport matrix elements

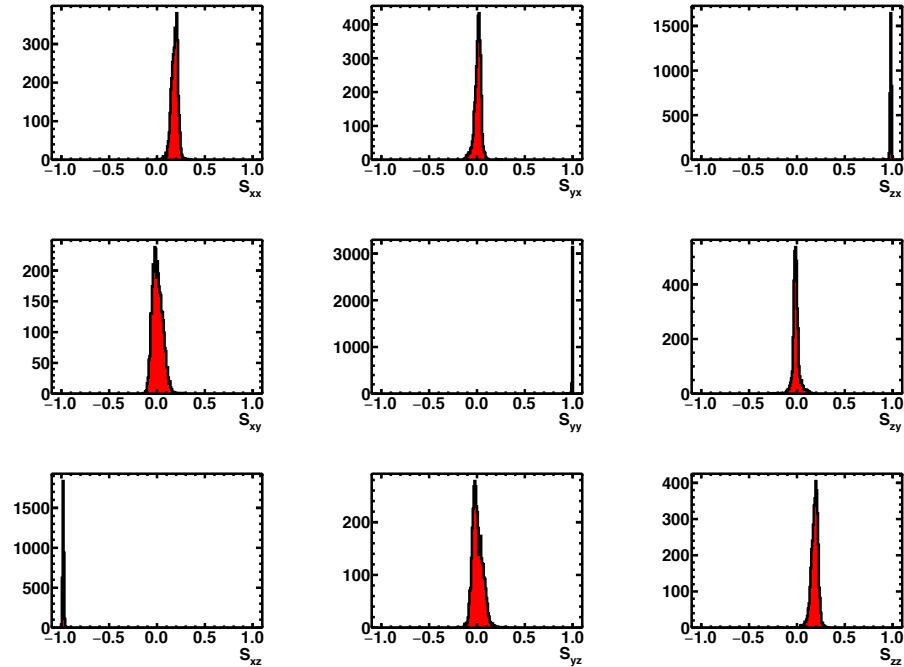
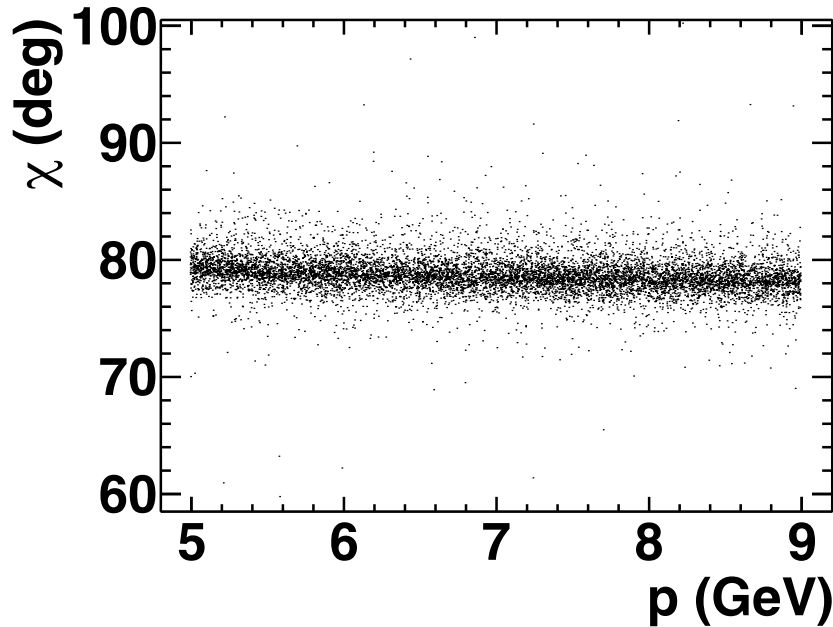
# SBS optics fitting from GEANT4



- SBS angle, vertex and momentum resolution for 5-9 GeV protons
- $\sigma(x_{\text{ptar}}) \sim 0.3$  mrad
- $\sigma(y_{\text{ptar}}) \sim 0.6$  mrad
- $\sigma(y_{\text{tar}}) \sim 1.5$  mm
- $\sigma(p)/p \sim 0.66\%$
- Improvement of the fit not significant beyond about 4<sup>th</sup>-order expansion of reconstruction coefficients

# Spin transport properties of SBS in GEP

$\chi \cdot 180.0 / \text{TMATH}::\text{PI}():p$



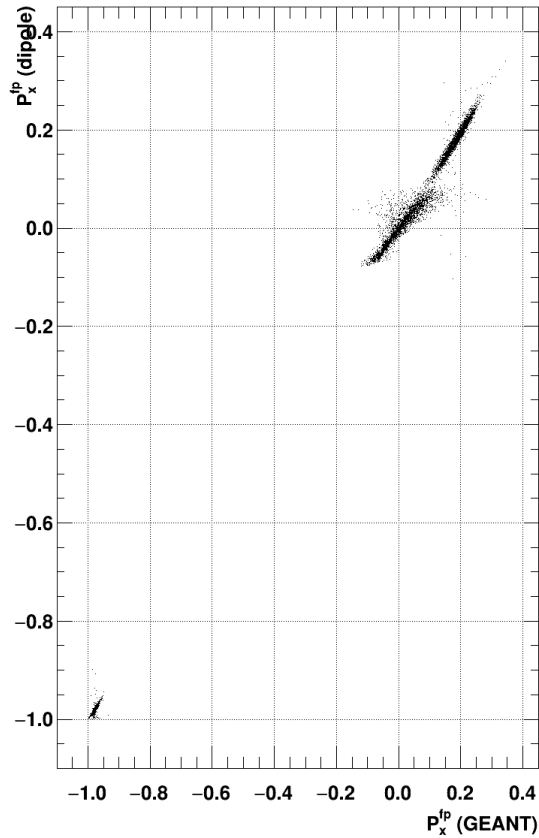
$$\frac{d\mathbf{S}}{dt} = \frac{e}{m\gamma} \mathbf{S} \times \left[ \frac{g}{2} \mathbf{B}_{\parallel} + \left( 1 + \gamma \left( \frac{g}{2} - 1 \right) \right) \mathbf{B}_{\perp} \right] \quad \text{BMT equation for protons}$$

- Spin precession in a magnetic field for is governed by the BMT equation
- For an almost pure dipole field, as in SBS, the proton spin precesses relative to its trajectory by an angle:  $\chi = \gamma \kappa_p \theta_{bend}$
- Precession angle is almost constant within useful acceptance of SBS for elastic ep events (cancellation between momentum dependence of gamma and thetabend)

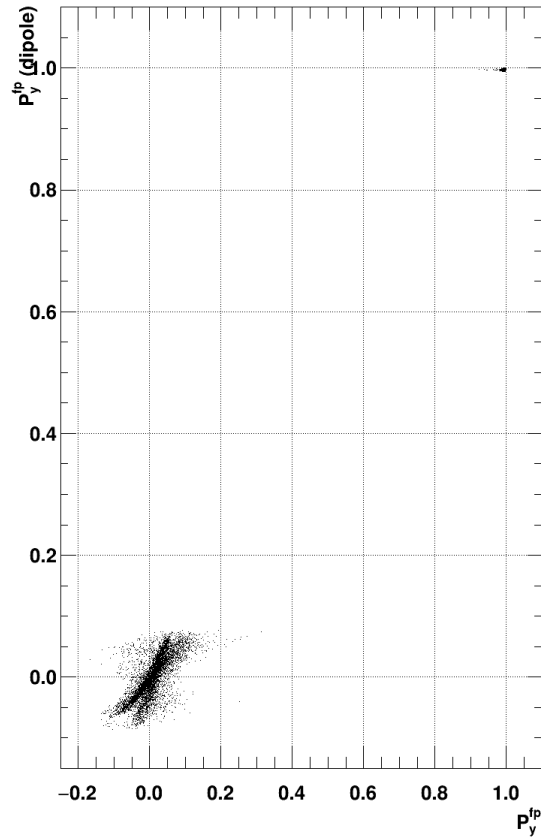


# Differences between dipole and full calculation

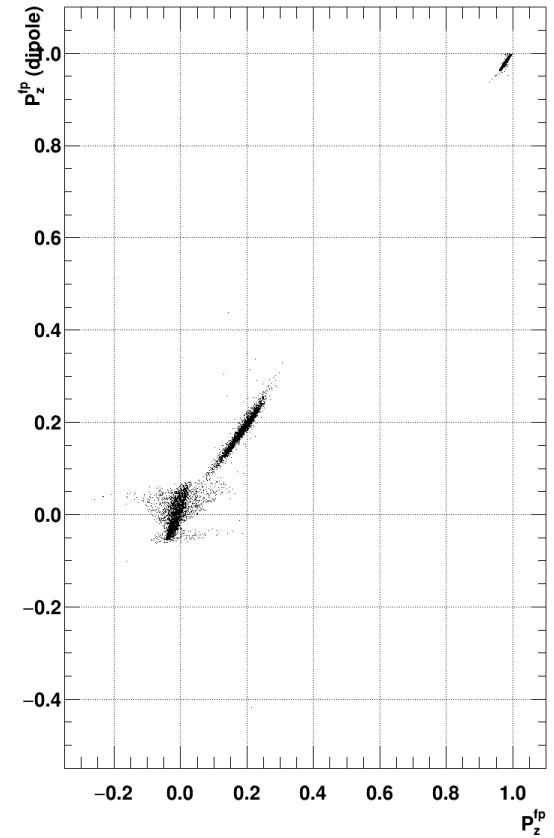
Px<sup>fp</sup>dipole:Px<sup>fp</sup>



Py<sup>fp</sup>dipole:Py<sup>fp</sup>

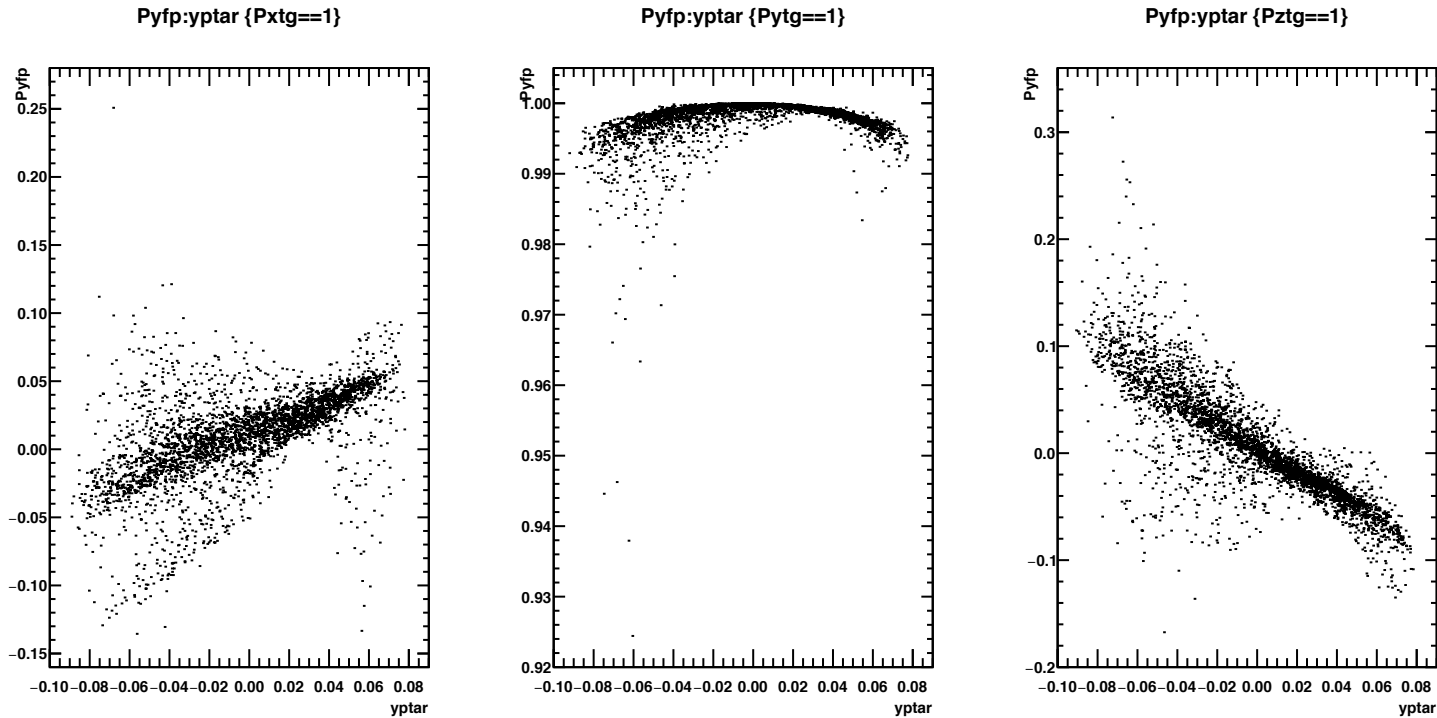


Pz<sup>fp</sup>dipole:Pz<sup>fp</sup>



- In dipole approximation, the spin component parallel to the field ( $P_y$ ) does not precess;
- Since SBS is non-focusing, the trajectory bend angle in the non-dispersive plane is close to zero for most trajectories
- Nevertheless, a small precession in the non-dispersive plane occurs for non-central trajectories
- This precession mixes  $P_T$  and  $P_L$  in extraction of  $GE_p$ —important systematic uncertainty

# Non-dispersive precession

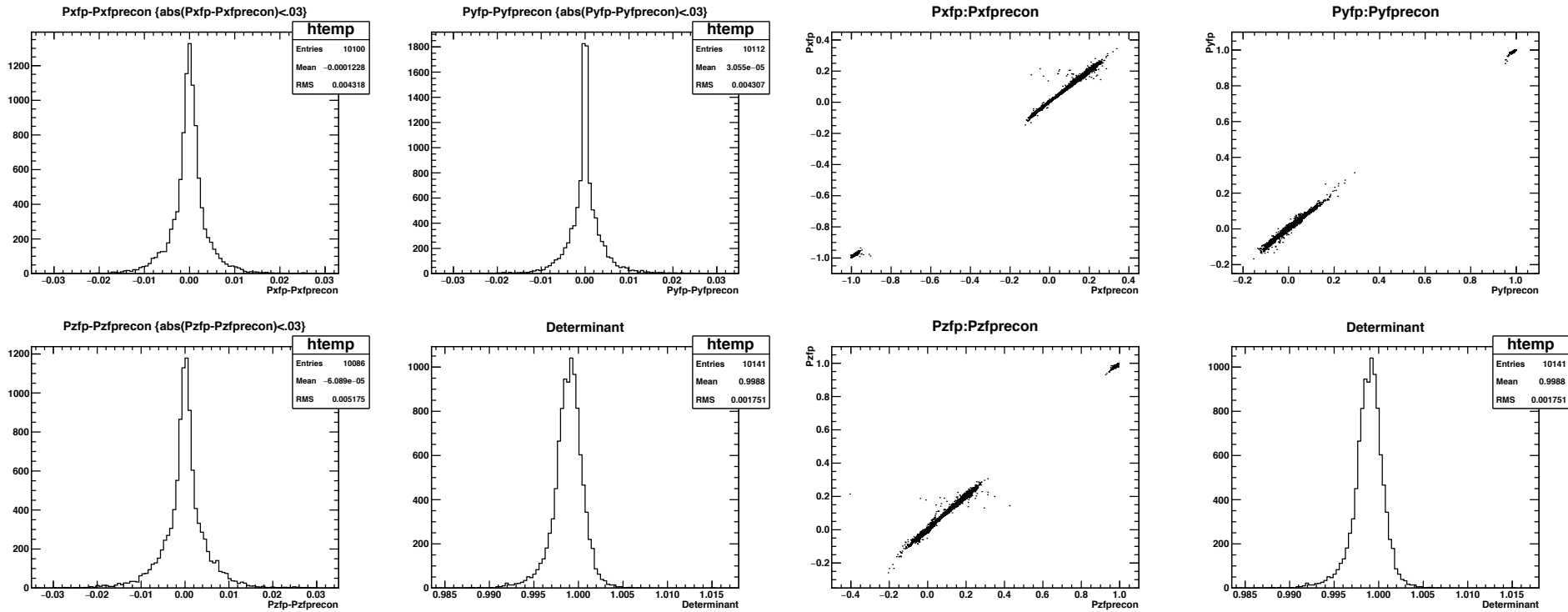


- Non-dispersive-plane precession is non-zero mainly for rays with “yptar” =  $dy/dz \neq 0$  (“phitarget” in usual Hall A notation)
- $S_{yx}$  has a weak positive correlation with yptar
- $S_{yz}$  (which mixes  $P_T$  and  $P_L$ ), has a stronger negative correlation with yptar
- The slope of this correlation sets the scale for how accurately yptar needs to be determined to achieve a given accuracy on GEp/GMp (no, we haven’t done the calculation yet)

# Formalism for fitting spin transport matrix elements

- In the usual Maximum-Likelihood analysis, the forward spin transport matrix elements are used. These in turn have to be computed from the reconstructed proton kinematics at the target.
- We expand the “small” deviations from the ideal dipole approximation as a power series in the proton trajectory parameters at the target:
  - $x_{tar}$ ,  $y_{tar}$ ,  $x_{ptar}$ ,  $y_{ptar}$ ,  $1/p$
- We fit using the SVD as in the case of the optics
- Additional technical challenge:
  - There is no guarantee that if we fit the individual matrix elements without enforcing any constraints, that the 3x3 matrix computed from the resulting expansion coefficients will be a proper rotation for any given event
  - We can try to fit the Euler angle and/or the angle-axis decomposition of the total rotation for a given event, but the problem is that such a decomposition is not unique, and it is difficult to define “good” expansion parameters.
  - Would like to come up with some kind of constrained optimization procedure to guarantee that the fitted Spin transport coefficients are guaranteed to give a proper rotation in each event

# Spin fit results (5<sup>th</sup>-order)



- Fit deviations from the ideal dipole approximation up to 5<sup>th</sup>-order (still some room for improvement)
- Fit matrix elements directly
- Don't enforce any constraints
- Determinant results close to 1 in any case.

# Status and plans

- Figure out how to enforce orthogonality in every event
- Add sieve slit to MC, develop optics and spin transport calibration/optimization procedures
- Quantify systematics