

Precision Measurement of the Neutron d_2 : The Color Field Response to the Polarized Nucleon

Graduate Students:

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Outline

- 1 The Spin Structure of the Nucleon
 - Current Status of the Spin Structure
- 2 d_2^n : Higher Twist Effects
 - Quark-Gluon Correlations
 - Expressions of d_2^n
- 3 The Experiment
 - The Measurement of d_2^n
 - The Experimental Setup
- 4 Preliminary Analysis
 - Left High-Resolution Spectrometer
 - BigBite Spectrometer
 - Compton Polarimeter
- 5 Projected Results
- 6 Summary

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Current Status of the Spin Structure

- For the nucleon, the total spin is broken down (in the light-cone gauge) into its constituents:
 - $\frac{1}{2} = \frac{1}{2}\Delta q + \Delta G + L_q + L_G$
 Δq = valence and sea quark spin
 ΔG = gluon spin
 $L_{q,G}$ = orbital angular momenta of quarks and gluons
 - Measurements show that quark contribution to the nucleon spin is $\sim 30\%$ (CERN, DESY, SLAC)
 - Measurements of the gluon contribution are negligible (BNL)
- All of this work corresponds to the investigation of g_1 under the interpretation of the Feynman Parton Model and pQCD
- The g_2 structure function is not as well known. It contains **quark-gluon correlations** inside the initial nucleon
 - Subsequently, g_2 does not have a simple interpretation in the Parton Model

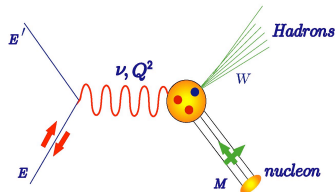
Polarized DIS

- Allows access to the **spin structure functions**:

$$\frac{d^2\sigma(\downarrow\uparrow-\uparrow\uparrow)}{dE'd\Omega} = \frac{4\alpha^2}{MQ^2} \frac{E'}{\nu E} \left[(E + E' \cos\theta) g_1(x, Q^2) - \frac{Q^2}{\nu} g_2(x, Q^2) \right]$$

$$\frac{d^2\sigma(\downarrow\Rightarrow-\uparrow\Rightarrow)}{dE'd\Omega} = \frac{4\alpha^2 \sin\theta}{MQ^2} \frac{E'^2}{\nu^2 E} [\nu g_1(x, Q^2) + 2E g_2(x, Q^2)]$$

- We can write g_1 and g_2 in terms of these **measurable** asymmetries and unpolarized cross sections



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Quark-Gluon Correlations (1)

- The g_2 structure function provides a direct probe into **quark-gluon interactions**, and is given by a spin-flip Compton amplitude
 - Seen in the imaginary part of virtual Compton scattering:

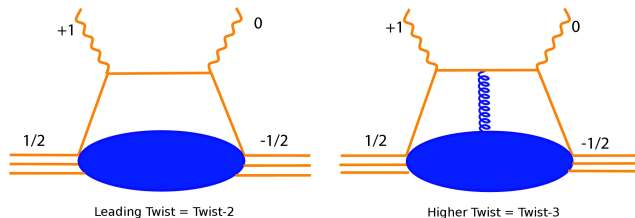


Figure: Higher twist contributions to virtual Compton scattering

Quark-Gluon Correlations (2)

The g_2 Structure Function

- g_2 can be broken into two parts:

$$g_2(x, Q^2) = g_2^{WW}(x, Q^2) + \overline{g_2}(x, Q^2)$$

where:

$$g_2^{WW}(x, Q^2) = -g_1(x, Q^2) + \int_x^1 \frac{g_1(y, Q^2)}{y} dy$$

$$\overline{g_2}(x, Q^2) = - \int_x^1 \frac{1}{y} \frac{\partial}{\partial y} \left[\frac{m_q}{M} h_T(y, Q^2) + \xi(y, Q^2) \right] dy$$

- The transverse polarization density (h_T) is suppressed by the ratio of the quark and target masses
- ξ is a twist-3 term arising from quark-gluon correlations

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Expressions of d_2^n (1)

d_2^n From the Structure Functions

- d_2^n is expressed as the second moment of a linear combination of g_1 and g_2 :

$$\begin{aligned} d_2^n(Q^2) &= \int_0^1 x^2 [2g_1(x, Q^2) + 3g_2(x, Q^2)] dx \\ &= 6 \int_0^1 x^2 \overline{g_2}(x, Q^2) dx \end{aligned}$$

- d_2^n is a **direct measure** of twist-3 effects in the neutron

Expressions of d_2^n (2)

d_2^n as a Lorentz Color Force

- The expression for the transverse (color) force on the active quark right after it is struck by the virtual photon in the interaction reads:

$$\begin{aligned}
 F^y(0) &\equiv -\frac{\sqrt{2}}{2P^+} \langle P, S | \bar{\psi}_q(0) G^{+y}(0) \gamma^+ \psi_q(0) | P, S \rangle \\
 &= -\frac{1}{2} M^2 d_2^n
 \end{aligned}$$

- d_2^n is a **measure** of this transverse Lorentz color force (M. Burkardt)

Expressions of d_2^n (3)

The Color Field Polarizabilities

- Considering OPE in the rest frame of the nucleon and introducing color singlet operators $O_{B,E}$,

$$O_B = \psi^\dagger g \vec{B} \psi \quad O_E = \psi^\dagger \vec{\alpha} \times g \vec{E} \psi$$

a relation containing the gluon **color field polarizabilities** χ is obtained:

$$\langle P, S | O_{B,E} | P, S \rangle = 2\chi_{B,E} M^2 \vec{S}$$

d_2^n can be written as a linear combination of $\chi_{B,E}$

$$d_2^n = \frac{1}{8} (\chi_E + 2\chi_B)$$

- d_2^n can be seen as a measure of the **response** of the gluon color fields to the polarization of the nucleon (X. Ji)

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The Measurement of d_2^n

- Writing g_1, g_2 in terms of $\sigma_0, A_{\parallel}, A_{\perp}$ we obtain the explicit form of d_2^n to be evaluated:

$$g_1 = \frac{MQ^2}{4\alpha^2} \frac{2y}{(1-y)(2-y)} \sigma_0 [A_{\parallel} + \tan(\theta/2) A_{\perp}]$$

$$g_2 = \frac{MQ^2}{4\alpha^2} \frac{y^2}{(1-y)(2-y)} \sigma_0 \left[-A_{\parallel} + \frac{1 + (1-y) \cos \theta}{(1-y) \sin \theta} A_{\perp} \right]$$

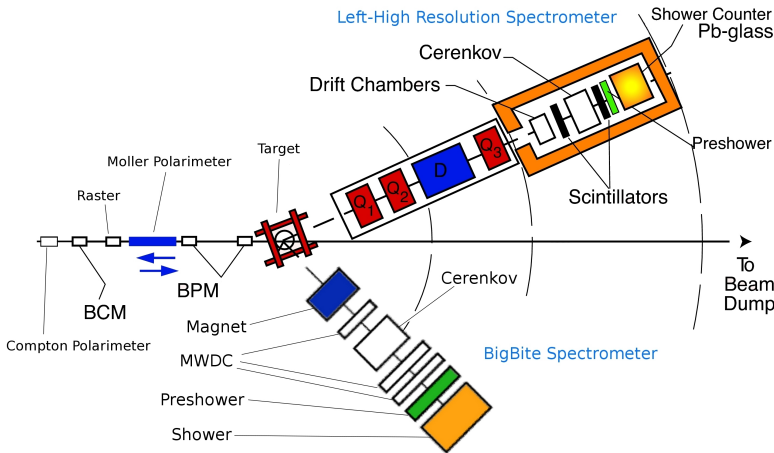
$$d_2^n = \int_0^1 \frac{MQ^2}{4\alpha^2} \frac{x^2 y^2}{(1-y)(2-y)} \sigma_0 \times \left[\left(3 \frac{1 + (1-y) \cos \theta}{(1-y) \sin \theta} + \frac{4}{y} \tan(\theta/2) \right) A_{\perp} + \left(\frac{4}{y} - 3 \right) A_{\parallel} \right] dx$$

$$A_{\parallel} = \frac{\sigma^{\downarrow\uparrow} - \sigma^{\uparrow\uparrow}}{2\sigma_0} \quad A_{\perp} = \frac{\sigma^{\downarrow\Rightarrow} - \sigma^{\uparrow\Rightarrow}}{2\sigma_0}$$

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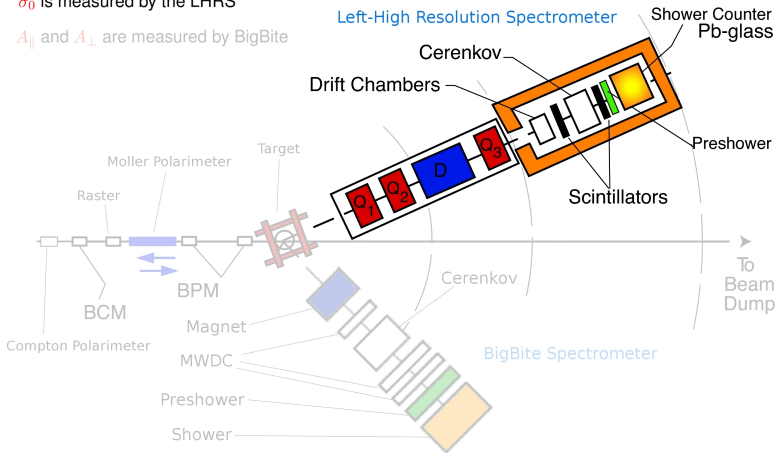
The Experimental Setup (1)



The Experimental Setup (1)

σ_0 is measured by the LHRS

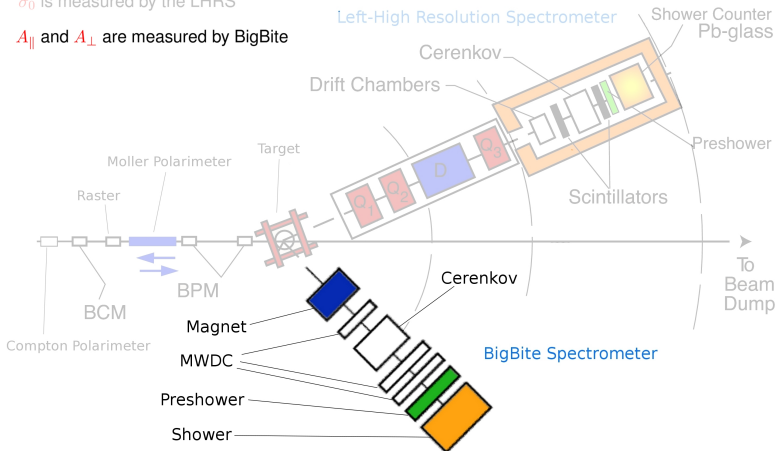
A_{\parallel} and A_{\perp} are measured by BigBite



The Experimental Setup (1)

σ_0 is measured by the LHRS

A_{\parallel} and A_{\perp} are measured by BigBite



The Experimental Setup (2)

LHRS Detectors

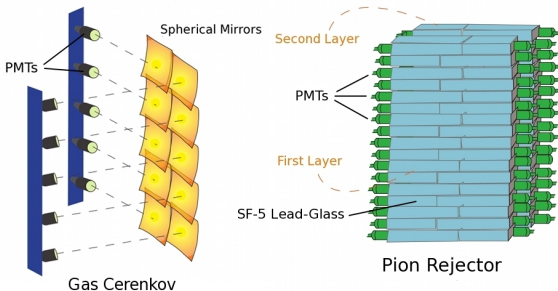


Figure: Drawings of the Gas Cerenkov and Pion Rejector in the Left High-Resolution Spectrometer

The Experimental Setup (3)

BigBite Gas Cerenkov

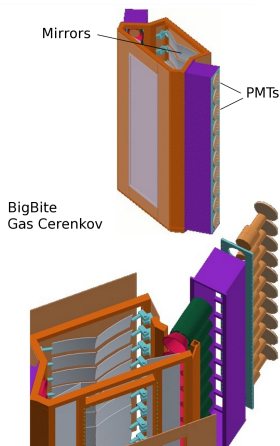
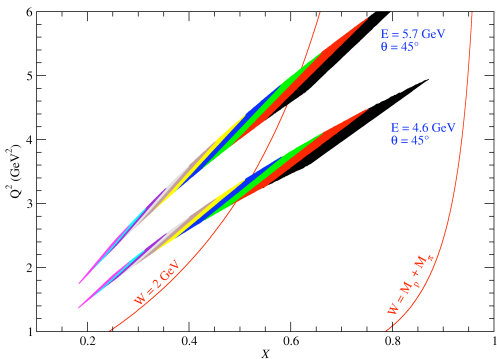


Figure: Drawings of the Gas Cerenkov in the BigBite Spectrometer

Kinematic Range



- The two bands represent the angular acceptance of the BigBite Spectrometer
- The ten colored stripes represent the different momentum settings in the LHRS for each beam energy

Outline

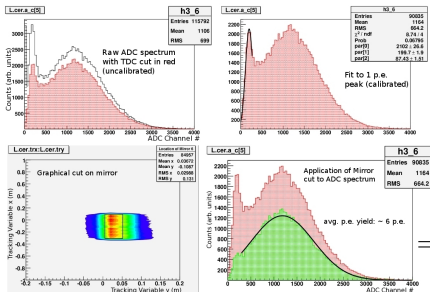
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Left High-Resolution Spectrometer (1)

Gas Cerenkov Calibration

- Gain-match
1 photoelectron (p.e.)
peaks
- Determine avg. # of p.e.
for each PMT

Mirror	1 p.e.	$n_{p.e.}$	# p.e.
1	196.8	967.3	4.45
2	198.6	715.3	3.48
3	198.4	1335	6.73
4	197.8	1344	6.75
5	198.5	1154	5.83
6	199.7	1184	5.93
7	199.5	1212	6.08
8	198.6	1225	6.17
9	194.9	1072	5.50
10	196.7	1010	5.13

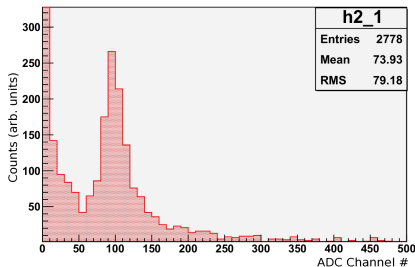


Left High-Resolution Spectrometer (2)

Pion Rejector Calibration

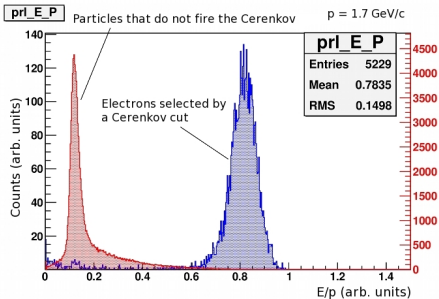
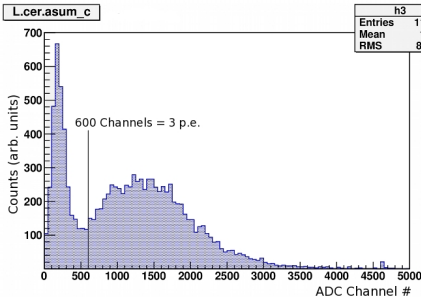
- We want each block in the pion rejector to have the same response for any given p
- Gain-match each PMT
 - Do this for the **pions** in the ADC spectra, since pions deposit the same amount of energy in the calorimeter for any p for the kinematic range of our experiment

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Left High-Resolution Spectrometer (3)

Pion Rejector Calibration



- A cut on the Cerenkov ADC spectrum (sum) is applied to the shower to select good electrons
- **Clean** separation of e^- , π

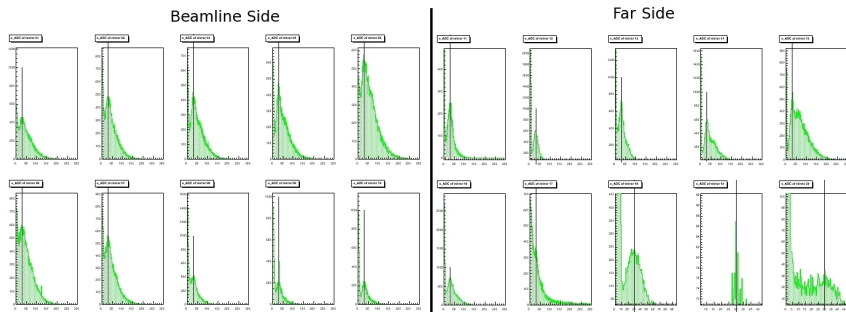
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BigBite Spectrometer (1)

Gas Cerenkov Calibration

- 1 p.e. gain-matching using LED runs:



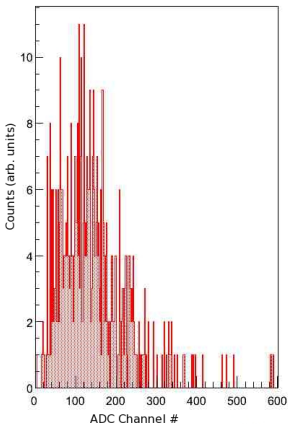
Plots provided by M. Posik

BigBite Spectrometer (2)

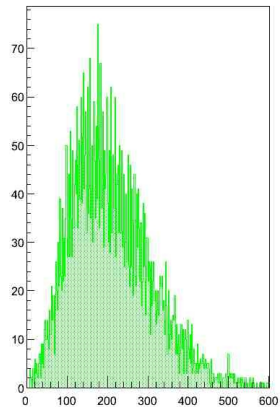
Gas Cerenkov Calibration

- Applying to production runs:

Background ADC of mirror 18



Good Electrons in ADC of mirror 18



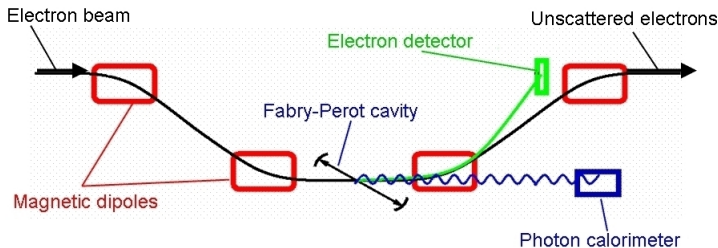
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Compton Polarimeter (1)

Schematic Diagram

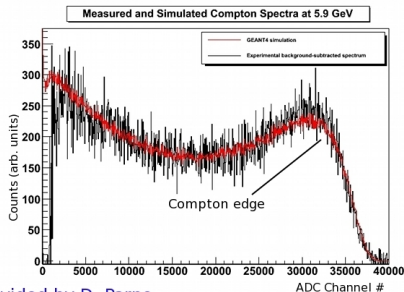
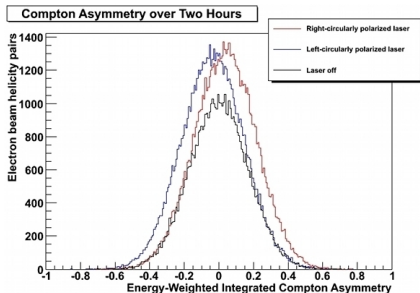


Drawing provided by D. Parno

- Polarization of the beam is measured using the **Compton Polarimeter**
- Relies on the **asymmetry** of the Compton cross section, due to the relative orientations of the e^- , γ polarizations

Compton Polarimeter (2)

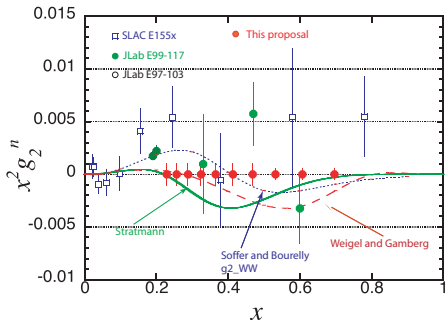
Asymmetries and Cross Sections



Plots provided by D. Parno

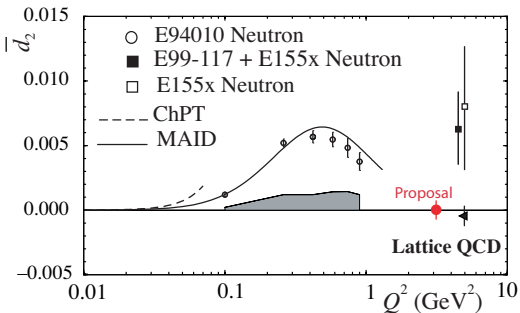
- The final Compton asymmetry is the weighted average of the mean asymmetries for the two polarization states (L,R)
- Compton spectrum shows a measurement of the scattered photon energy, proportional to the Compton cross section. The plot shows the **Compton edge** – the maximum amount of energy a photon can acquire

Projected Error on $x^2 g_2^n(x, Q^2)$



- The experiment is designed to minimize the **moment** of g_2 , so as to improve the accuracy of d_2^n

Projected Measurement of d_2^n



- This experiment is expected to have a statistical uncertainty of $\Delta d_2^n = 5 \times 10^{-4}$
 - **Four** times better than the current world average
 - Provides a **benchmark** test for Lattice QCD

Summary

- Interested in **quark-gluon correlations**
- Investigate this by exploring transverse **spin interactions** through the g_2 structure function leading to higher twist effects seen in the matrix element d_2^n
- Sheds light upon the **Lorentz color force** inside the nucleon
- This measurement provides a **benchmark** test on Lattice QCD

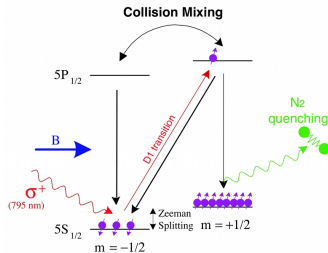
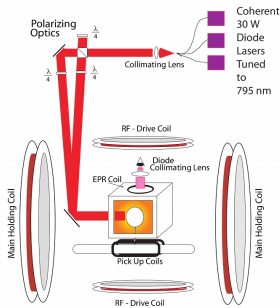
Future Work

- Short term goals:
 - Continue work on calibrations in LHRS, BB, and Compton
 - Work towards extracting preliminary $\sigma_0, A_{\parallel}, A_{\perp}$

Acknowledgements

- I would like to thank the spokespeople X. Jiang, S. Choi, B. Sawatzky, and Z.-E. Meziani
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^3He Target



- Vaporized Rb is optically pumped using circularly polarized light to polarize its electrons
- Through **collision mixing** the Rb electrons transfer their spin to the ^3He nuclei