Measurement of the Coulomb quadrupole amplitude in the $\gamma^* p \rightarrow \Delta(1232)$ reaction in the low momentum transfer region

David Anez Dalhousie University April 20, 2010

Introduction

- Take measurements of the $p(e,e'p)\pi^0$ (pion electroproduction) reaction
 - \square At energies in the area of the Δ resonance
 - With low momentum transfer (Q^2) between the electron and proton
 - To better understand the Coulomb quadrupole transition amplitude behavior in this region and how it affects nucleon deformation

Overview

- Motivation: Constituent Quark Model
- Kinematics and Transition Amplitudes
- Response Functions and Amplitude Extraction
- World Data and Models
- Experimental Setup
- Conclusions

- Three "heavy" quarks in nucleon
 - Each quark has mass $\frac{1}{3}$ of nucleon
 - Each quark has intrinsic spin angular momentum of $\frac{1}{2}$
 - Combines to give $S = \frac{1}{2}$ or $S = \frac{3}{2}$
 - If L = 0
 S = ¹/₂, J^π = ¹/₂⁺ → corresponds to N(939)
 S = ³/₂, J^π = ³/₂⁺ → corresponds to Δ(1232)
 If L = 2
 S = ¹/₂, J^π = ³/₂⁺ → corresponds to Δ(1232)
 S = ³/₂, J^π = ¹/₂⁺ → corresponds to N(939)

Wave functions created $|N(939)\rangle = a_{S} |(S = \frac{1}{2}, L = 0)J^{\pi} = \frac{1}{2}^{+}\rangle + a_{D} |(S = \frac{3}{2}, L = 2)J^{\pi} = \frac{1}{2}^{+}\rangle$ $|\Delta(1232)\rangle = b_{S} |(S = \frac{3}{2}, L = 0)J^{\pi} = \frac{3}{2}^{+}\rangle + b_{D} |(S = \frac{1}{2}, L = 2)J^{\pi} = \frac{3}{2}^{+}\rangle$



- Measure non-spherical components by measuring quadrupole moment
 - Cannot measure quadrupole moment directly
 - Measure quadrupole moment of $N \rightarrow \Delta$ transition
- Three electromagnetic transitions
 - M1 magnetic dipole
 - \Box E2 electric quadrupole
 - C2 Coulomb/scalar quadrupole

Magnetic Dipole

- Spin-flip
- Dominant



- Electric quadrupole
- Coulomb quadrupole
 - Only with virtual photons





One-body interactions

$$\square \quad \hat{Q}_{[1]} = \sqrt{\frac{16\pi}{5} \sum_{i=1}^{3} e_i r_i^2 Y_0^2(\vec{r}_i)} = \sum_i e_i \left(3z_i^2 - r_i^2\right)$$

Two-body interactions





Pion Cloud



Kinematics – Electronic Vertex

- Incoming electron
 - Energy E
 - Momentum \vec{k}_i
- Scattered electron
 - □ Energy *E*′
 - Momentum \bar{k}_f
 - Angle θ_e
- Virtual photon
 - Energy ω
 - Momentum \bar{q}
 - Angle θ_q



Kinematics – Electronic Vertex

• Momentum transfer, Q^2

$$Q^{2} = -q^{2} = -(\omega^{2} - \bar{q}^{2})$$

$$Q^{2} \approx 4EE' \sin^{2} \frac{\theta_{e}}{2}$$

$$e^{\frac{\vec{k}_{i}}{\sqrt{\theta_{e}}}}$$

$$e^{\frac{\vec{k}_{i}}{\sqrt{\theta_{e}}}}$$

$$e^{\frac{\vec{k}_{i}}{\sqrt{\theta_{e}}}}$$

∽v*

Kinematics – Hadronic Vertex

Recoil proton

- Energy E_p
- Momentum \bar{p}_p
- Angle θ_{pq}
- Recoil pion
 - Energy E_{π}
 - Momentum \bar{p}_{π}
 - Angle θ_{π}
 - Not detected



Kinematics – Planes

- Scattering plane $-\vec{k_i}$ and $\vec{k_f}$
- Recoil plane \vec{p}_p and \vec{p}_{π}
- Azimuthal angle ϕ_{pq}



Multipole Amplitudes

- General form of πN Multipoles: $X_{\ell \pm}^{I}$
 - $\Box X type of excitation (M, E, S)$
 - \Box *I* isospin of excited intermediate state

$$\Box \ \ell \pm - J_{\Delta} = \ell \pm \frac{1}{2}$$

- Magnetic dipole $M1 / M_{1+}^{3/2}$
- Electric quadrupole $E2 / E_{1+}^{3/2}$
- Coulomb quadrupole C2 / $S_{1+}^{3/2}$ • $\omega S_{1+} = |\vec{q}| L_{1+}$

Multipole Amplitudes

γN-Multipoles	Initial State		Excited State			Final State		πN -Multipoles	
С, Е, М	L_{γ}^{π}	$S \frac{\pi}{N}$	J_{R}^{π}	$N^* I_{2I2J} \Delta$		$S \frac{\pi}{N}$	I_{π}^{π}	$L_{\ell\pm}, E_{\ell\pm}, M_{\ell\pm}$	
<i>C</i> 0	0+	1/2+	1/2+	P ₁₁	P ₃₁	1/2+	1+	L ₁₋	
<i>C</i> 1, <i>E</i> 1	1-	1/2+	1/2-	S ₁₁	S ₃₁	1/2+	0-	L_{0+}, E_{0+}	
		1/2+	3/2-	D ₁₃	D ₃₃	1/2+	2-	L ₂₋ , E ₂₋	
<i>M</i> 1	1+	1/2+	1/2+	P ₁₁	P ₃₁	1/2+	1+	<i>M</i> ₁₋	
		1/2+	³ / ₂ +	P ₁₃	P ₃₃	1/2+	1+	<i>M</i> ₁₊	
C2, E2	2+	1/2+	3/2+	P ₁₃	P ₃₃	1/2+	1+	L_{1+}, E_{1+}	
		1/2+	5/2+	F ₁₅	F ₃₅	1/2+	3+	L ₃₋ , E ₃₋	
М2	2-	1/2+	3/2-	D ₁₃	D ₃₃	1/+	2-	M ₂₋	
		1/2+	5/2-	D ₁₅	D ₃₅	1/2+	2-	M ₂₊	

Multipole Amplitudes

- E_{1+} and S_{1+} at same magnitude as background amplitudes
- Measure ratio to dominant M_{1+}

• EMR =
$$R_{EM}^{3/2} = \operatorname{Re}\left(\frac{E_{1+}^{3/2}}{M_{1+}^{3/2}}\right) = \frac{\operatorname{Re}\left(E_{1+}^*M_{1+}\right)}{\left|M_{1+}\right|^2}$$

• CMR = $R_{CM}^{3/2} = \operatorname{Re}\left(\frac{S_{1+}^{3/2}}{M_{1+}^{3/2}}\right) = \frac{\operatorname{Re}\left(S_{1+}^*M_{1+}\right)}{\left|M_{1+}\right|^2}$

 Unpolarized cross section made up of four independent partial cross sections

$$\frac{d^5 \sigma}{dk_f d\Omega_e d\Omega^*} = \Gamma_{\gamma} \frac{k}{q_0} \left(\sigma_L + \sigma_T + \sigma_{LT} + \sigma_{TT} \right)$$

$$\Gamma_{\gamma} = \frac{\alpha}{2\pi^2} \frac{k_f}{k_i} \frac{k_{\gamma}}{Q^2} \frac{1}{1-\varepsilon} \qquad \varepsilon_s = \frac{Q^2}{q^2} \varepsilon \qquad k = \sqrt{\frac{W^2 + m_{\pi}^2 - m_p^2}{4W^2} - m_{\pi}^2}$$

$$k_{\gamma} = \frac{W^2 - m_p^2}{2m_p} \qquad \varepsilon = \left(1 + 2\frac{q^2}{Q^2} \tan^2 \frac{\theta_e}{2} \right)^{-1} \qquad q_0 = \frac{W^2 - m_p^2}{2W}$$

 Unpolarized cross section made up of four independent partial cross sections

$$\frac{d^{5}\sigma}{dk_{f}d\Omega_{e}d\Omega^{*}} = \Gamma_{\gamma}\frac{k}{q_{0}}(\sigma_{L} + \sigma_{T} + \sigma_{LT} + \sigma_{TT})$$

$$\sigma_L = \varepsilon_S R_L$$
 $\sigma_{LT} = \sqrt{2\varepsilon_S(1+\varepsilon)}R_{LT}\sin\theta_{pq}^*\cos\phi_{pq}^*$

$$\sigma_T = R_T \qquad \sigma_{TT} = \varepsilon R_{TT} \sin^2 \theta_{pq}^* \cos 2\phi_{pq}^*$$

$$R_{L}\sqrt{\frac{\omega_{cm}^{2}}{Q^{2}}} = |L_{0+}|^{2} + 4|L_{1+}|^{2} + |L_{1-}|^{2} - 4\operatorname{Re}\left\{L_{1+}^{*}L_{1-}\right\} + 2\cos\theta\operatorname{Re}\left\{L_{0+}^{*}\left(4L_{1+}+L_{1-}\right)\right\} + 12\cos^{2}\theta\left(|L_{1+}|^{2} + \operatorname{Re}\left\{L_{1+}^{*}L_{1-}\right\}\right)$$

$$R_{T} = |E_{0+}|^{2} + \frac{1}{2}|2M_{1+} + M_{1-}|^{2} + \frac{1}{2}|3E_{1+} - M_{1+} + M_{1-}|^{2} + 2\cos\theta\operatorname{Re}\left\{E_{0+}^{*}\left(3E_{1+} + M_{1+} - M_{1-}\right)\right\} + \cos^{2}\theta\left(|3E_{1+} + M_{1+} - M_{1-}|^{2} - \frac{1}{2}|2M_{1+} + M_{1-}|^{2} - \frac{1}{2}|3E_{1+} - M_{1+} - M_{1-}|^{2}\right)$$

$$R_{LT}\sqrt{\frac{\omega_{cm}^{2}}{Q^{2}}} = -\sin\theta\operatorname{Re}\left\{L_{0+}^{*}\left(3E_{1+} - M_{1+} + M_{1-}\right) - \left(2L_{1+}^{*} - L_{1-}^{*}\right)E_{0+} + 6\cos\theta\left(L_{1+}^{*}\left(E_{1+} - M_{1+} + M_{1-}\right) + L_{1-}^{*}E_{1+}\right)\right\}$$

$$R_{TT} = 3\sin^2\theta \left(\frac{3}{2}|E_{1+}|^2 - \frac{1}{2}|M_{1+}|^2 - \operatorname{Re}\left\{E_{1+}^*(M_{1+} - M_{1-}) + M_{1+}^*M_{1-}\right\}\right)$$

Truncated Multipole Expansion

$$R_{L} \approx 0$$

$$R_{T} \approx \frac{5}{2} |M_{1+}|^{2} + 2\cos\theta \operatorname{Re} \{ E_{0+}^{*} M_{1+} \} - \frac{3}{2}\cos^{2}\theta |M_{1+}|^{2}$$

$$R_{LT} \approx \sin\theta \operatorname{Re} \{ L_{0+}^{*} M_{1+} \} - 6\cos\theta (L_{1+}^{*} M_{1+})$$

$$R_{TT} \approx -3\sin^{2}\theta (\frac{1}{2} |M_{1+}|^{2} + \operatorname{Re} \{ E_{1+}^{*} M_{1+} + M_{1+}^{*} M_{1-} \}$$

- Model Dependent Extraction
 - Fit theoretical model to existing data
 - Insert model values for background amplitudes

- Models
 - MAID
 - SAID
 - DMT
 - □ Sato-Lee
 - Chiral EFT
 - Lattice QCD

- $p(e,e'p)\pi^0$ experiments
 - □ CEA 1969
 - □ DESY 1970-1972
 - □ NINA 1971
 - □ ELSA 1997
 - □ MIT-Bates 2000
 - □ MAMI 2001
 - □ CLAS 2002
 - □ MAMI 2005-2006













- Sato and Lee
 - Suggest separating nucleon into quark core and pion cloud
 - "bare" quark core links to lattice QCD
 - "full" nucleon links to experimental data



The Experiment

- Jefferson Lab, Hall A
- April 3rd April 8th, 2011
- 1115 MeV, 75µA *e*⁻ beam
- 6 cm LH_2 target
- Two high resolution spectrometers
 - □ HRSe and HRSh



High Resolution Spectrometers

- Vertical drift chambers
 - Particle tracking
- Scintillators
 - **D** Timing information
 - Triggering DAQ
- Cerenkov detectors
 - Aerogel and gas
 - Particle identification
- Lead glass showers
 - Particle identification



Settings

$Q^2 (\text{GeV}/c)^2$	W(MeV)	$\boldsymbol{\theta}_{\scriptscriptstyle pq}^{*}{}^{\boldsymbol{\circ}}$	$ heta_{_{\!$	$P'_{e}(\text{MeV/}c)$	$\boldsymbol{\theta}_{p}^{~\boldsymbol{\circ}}$	$P_p'(\text{MeV/}c)$	Time (hrs)
0.040	1221	0	12.52	767.99	24.50	547.54	1.5
0.040	1221	30	12.52	767.99	12.52	528.12	2
0.040	1221	30	12.52	767.99	36.48	528.12	3.5
0.040	1260	0	12.96	716.42	21.08	614.44	1.5
0.090	1230	0	19.14	729.96	29.37	627.91	1.5
0.090	1230	40	19.14	729.96	14.99	589.08	3
0.090	1230	40	19.14	729.96	43.74	589.08	4.5
0.125	1232	0	22.94	708.69	30.86	672.56	3.5
0.125	1232	30	22.94	708.69	20.68	649.23	7
0.125	1232	30	22.94	708.69	41.03	649.23	7
0.125	1232	55	22.94	708.69	12.52	596.43	3.5
0.125	1232	55	22.94	708.69	49.19	596.43	3.5
0.125	1170	0	21.74	788.05	37.31	575.57	3
0.125	1200	0	22.29	750.16	34.06	622.63	2
Configuration changes							17
Calibrations							8
Total:							72

Conclusion

- Important step forward in understanding nucleon's internal structure
- Help bridge and validate experimental world data
- Help theoretical models better understand
 role of pion cloud in nucleon deformation
 role of QCD in low momentum transfer region