

Threshold Pion Electro-Production at High Q^2

$$\gamma^* + p = \pi^0 + p'$$

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- Motivation: See recent talk by Braun at Jlab 12 GeV upgrade Workshop. Study PQCD
- Feasibility study for new measurements from $Q^2 = 0.50 - 10 \text{ GeV}^2$ at 12 GeV. Investigate kinematics and W resolution and missing mass resolution
 - Beam Energy 4 GeV - 12 GeV
 - Use HRS for Electron- Need good W resolution 1 MeV
 - Use MPS for Proton - large acceptance
- Several people and groups have shown interest In such a proposal recoil Including polarization measurements.

basic idea: decouple “hard” and “soft” momenta

$$M(Q_1^2, Q_2^2, \dots; k_1^2, k_2^2, \dots) = H(Q_1^2, Q_2^2, \dots; \mu^2) \otimes S(\mu^2; k_1^2, k_2^2, \dots)$$

↑ ↑
factorization scale

- H: hard function can be calculated in pQCD (if we are lucky)
- S: soft function can be simplified using LET (if we are lucky)

High Q^2

Low t

$$\gamma^* + p = \pi^0 + p'$$

Differential Cross Section

For unpolarized protons, the virtual photon cross section is

$$d\sigma_{\gamma^*} = \frac{\alpha_{em}}{8\pi} \frac{k_f}{W} \frac{d\Omega_\pi}{W^2 - m_N^2} |\mathcal{M}_{\gamma^*}|^2$$

with

$$|\mathcal{M}_{\gamma^*}|^2 = M_T + \epsilon M_L + \sqrt{2\epsilon(1+\epsilon)} M_{LT} \cos(\phi_\pi) + \epsilon M_{TT} \cos(2\phi_\pi) + \lambda \sqrt{2\epsilon(1-\epsilon)} M'_{LT} \sin(\phi_\pi)$$

$$f_\pi^2 M_T = \frac{4\bar{k}_i^2 Q^2}{m_N^2} |G_1^{\pi N}|^2 + \frac{c_\pi^2 g_A^2 \bar{k}_f^2}{(W^2 - m_N^2)^2} Q^2 m_N^2 G_M^2 + \cos\theta \frac{c_\pi g_A |\bar{k}_i| |\bar{k}_f|}{W^2 - m_N^2} 4Q^2 G_M \text{Re } G_1^{\pi N}$$

$$f_\pi^2 M_L = \bar{k}_i^2 |G_2^{\pi N}|^2 + \frac{4c_\pi^2 g_A^2 \bar{k}_f^2}{(W^2 - m_N^2)^2} m_N^4 G_E^2 - \cos\theta \frac{c_\pi g_A |\bar{k}_i| |\bar{k}_f|}{W^2 - m_N^2} 4m_N^2 G_E \text{Re } G_2^{\pi N}$$

$$f_\pi^2 M_{LT} = -\sin\theta \frac{c_\pi g_A |\bar{k}_i| |\bar{k}_f|}{W^2 - m_N^2} Q m_N \left[G_M \text{Re } G_2^{\pi N} + 4G_E \text{Re } G_1^{\pi N} \right]$$

$$f_\pi^2 M_{TT} = 0,$$

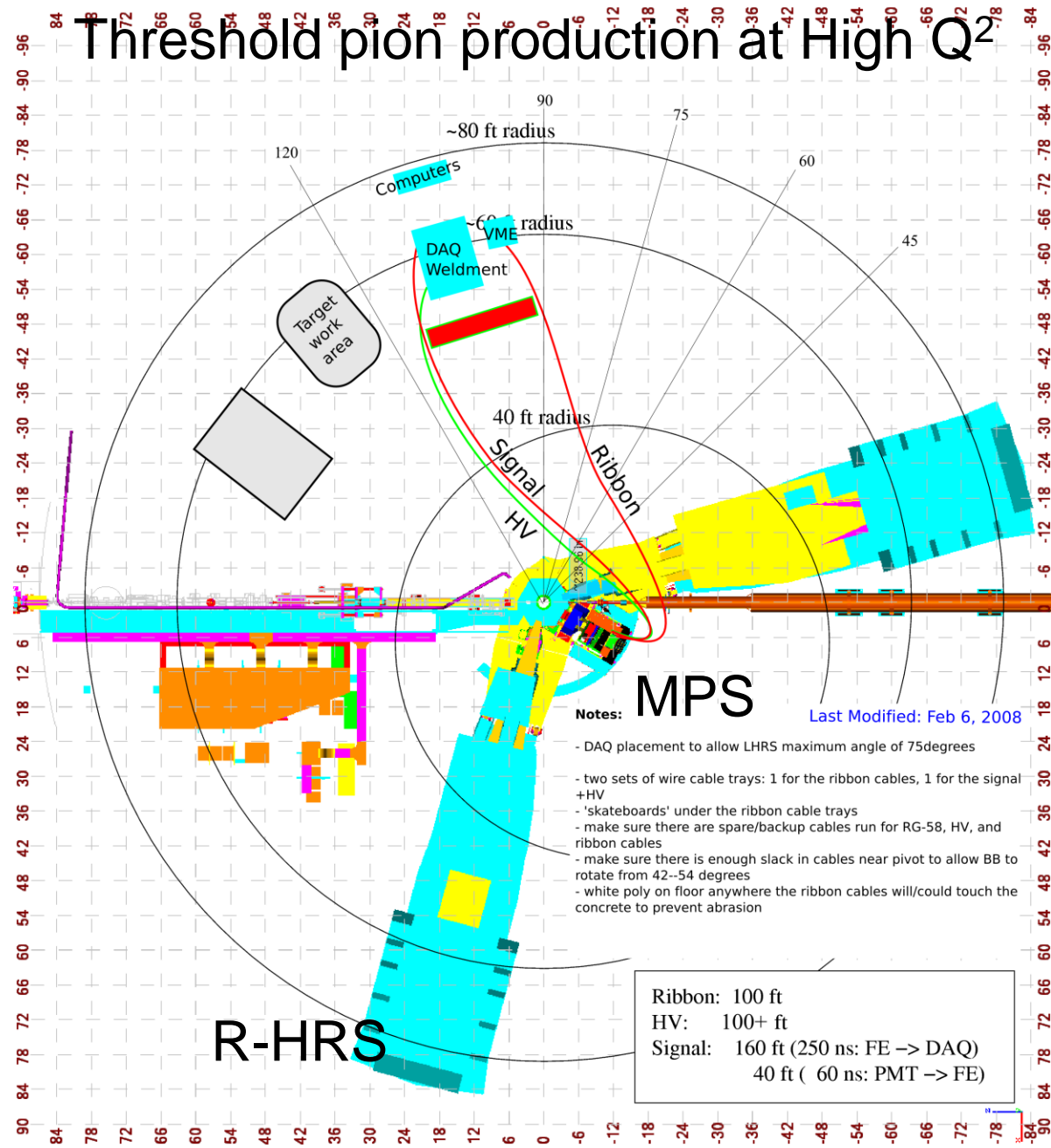
$$f_\pi^2 M'_{LT} = -\sin\theta \frac{c_\pi g_A |\bar{k}_i| |\bar{k}_f|}{W^2 - m_N^2} Q m_N \left[G_M \text{Im } G_2^{\pi N} - 4G_E \text{Im } G_1^{\pi N} \right]$$

$p(e,e'p)\pi^0$				
Facility	$Q^2(\text{GeV}^2)$	$W(\text{GeV})$	Ref	Comment
SLAC	7.1,9.4	1.10 -1.40	P.E. Bosted et al. PR D 49,3091(1994)	Poor W Resolution
Jlab Hall C	2.8,4.0	1.1 - 1.7	V.V Frolov et al. PRL 82,45(1999)	Delta & S11 No threshold data
Jlab Hall A	1.0	1.1 - 2.0	G. Laveissiere et al. PRC69,45203(2004)	Delta No Threshold Data
Jlab Hall A	1.0	1.23	J. Kelly et al. PRL 95,10200(2005)	Recoil Polarization No Threshold data
Jlab Hall C	6.4,7.7	1.1 - 1.7	E01-002	Delta & S11 No Threshold data
Jlab Hall B	1 - 4	Threshold	E01-113	DVCS Poor W Resolution
$p(e,e' \pi^+)n$				
Jlab Hall C	1 - 10	Above Resonances	T. Horn Proposal PAC 32	Test PQCD Concepts, Scaling
$p(p,p'p)\pi^0$				
IUCF E=400,325	Low	Threshold	H.O. Meyer et al. PRC 63,64002(2001)	Spin Observables

Threshold pion production at High Q²

Beam

L-HRS



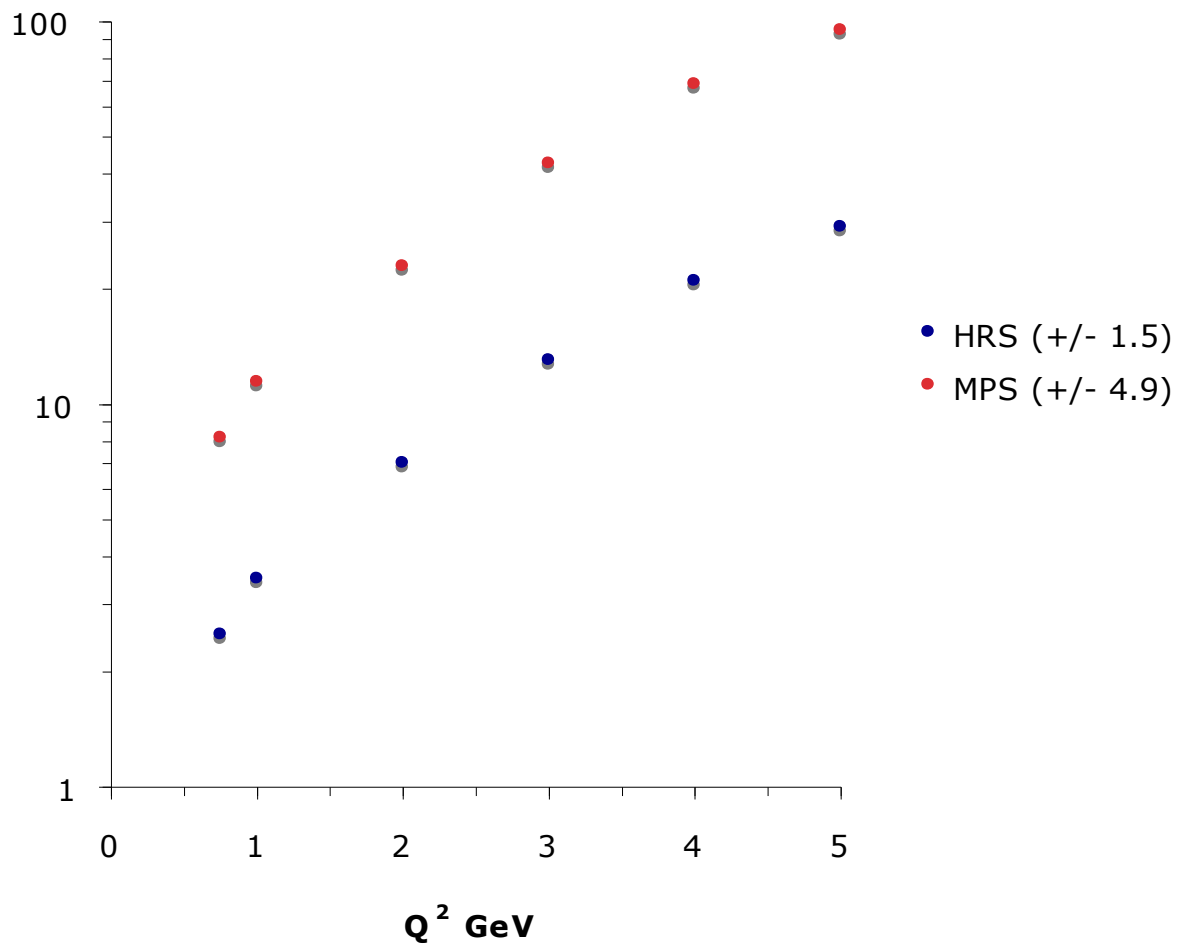
Notes: **MPS** Last Modified: Feb 6, 2008

- DAQ placement to allow L-HRS maximum angle of 75degrees
- two sets of wire cable trays: 1 for the ribbon cables, 1 for the signal +HV
- 'skateboards' under the ribbon cable trays
- make sure there are spare/backup cables run for RG-58, HV, and ribbon cables
- make sure there is enough slack in cables near pivot to allow BB to rotate from 42-54 degrees
- white poly on floor anywhere the ribbon cables will/could touch the concrete to prevent abrasion

Ribbon:	100 ft
HV:	100+ ft
Signal:	160 ft (250 ns: FE -> DAQ)
	40 ft (60 ns: PMT -> FE)

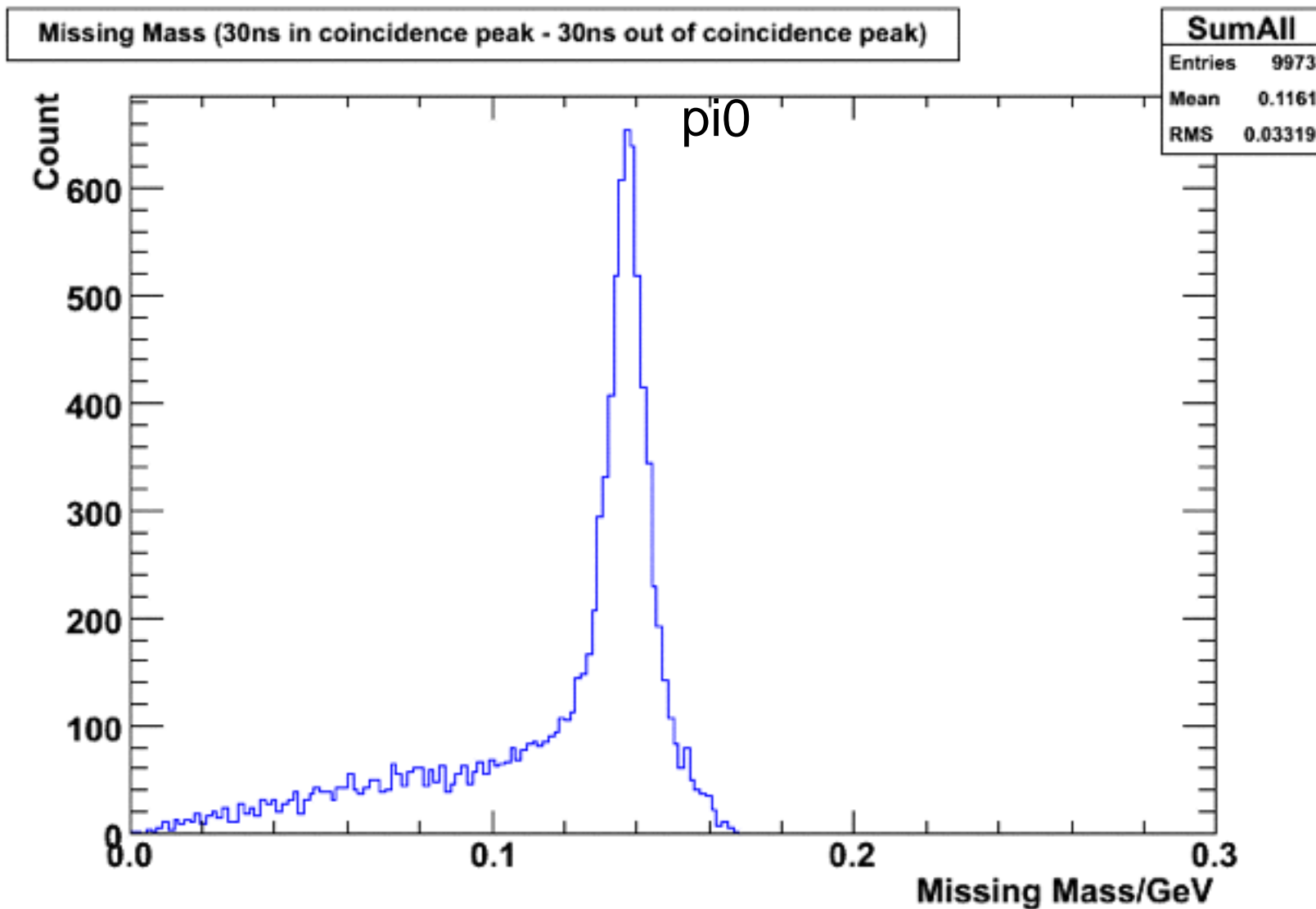
R-HRS

Range of ΔW that includes complete kinematic coverage for the proton detected in the HRS or MPS vs Q^2

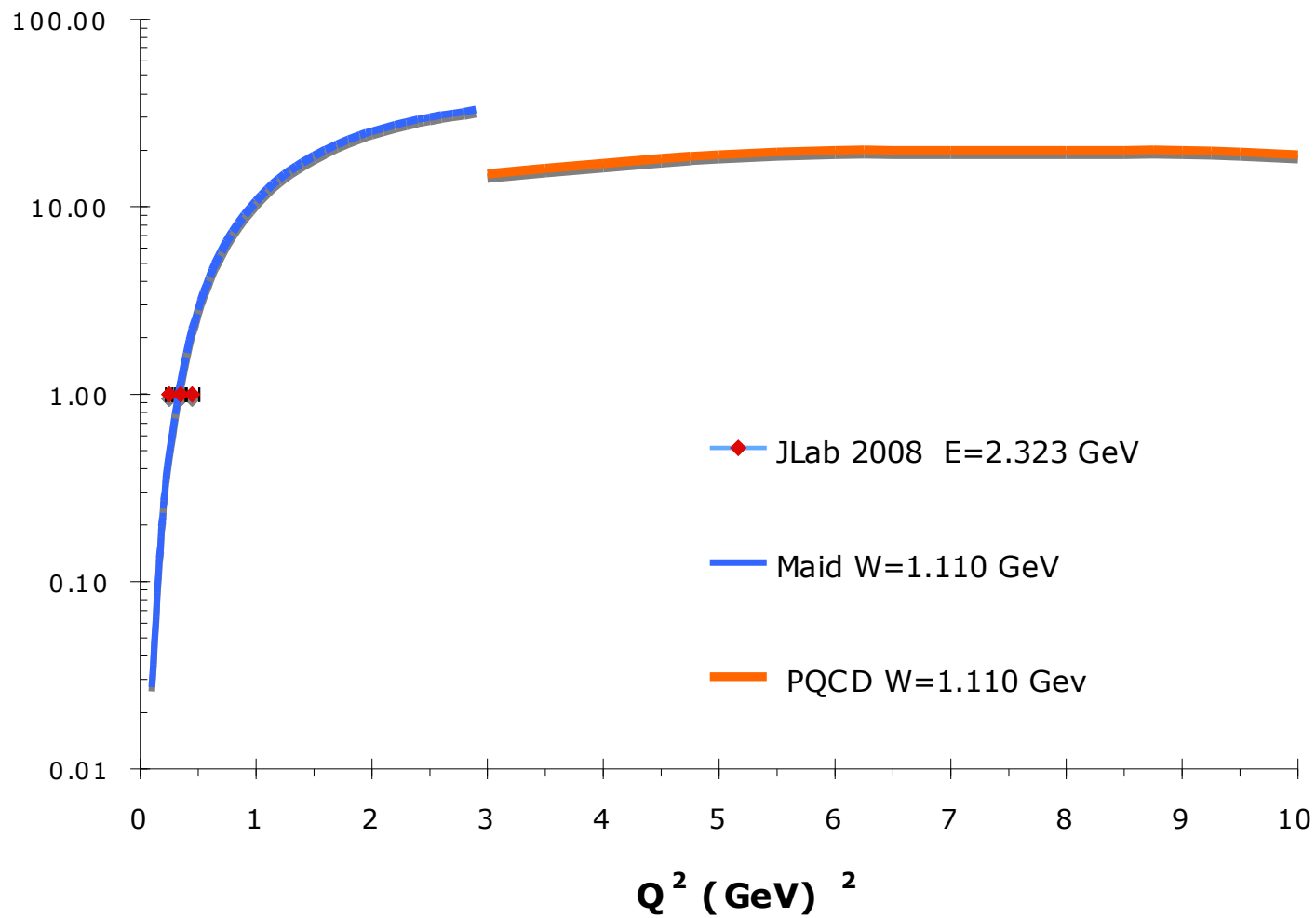


Data from E04-007

$$Q^2 = -0.45 \text{ GeV}^2$$

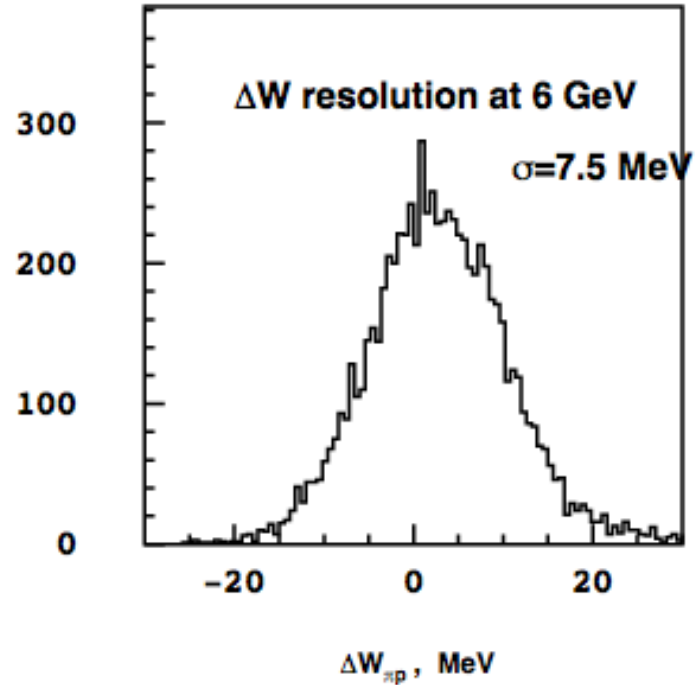
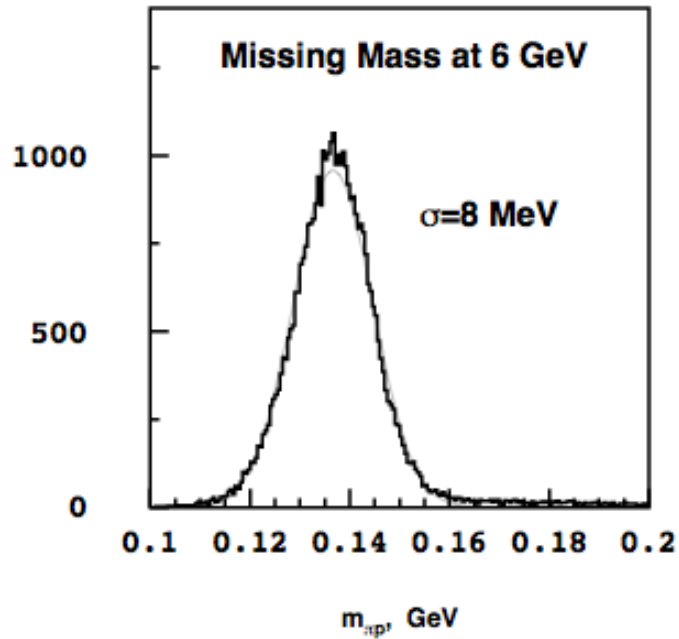


$\sigma Q^6 (\mu\text{b-GeV}^6) \text{ vs } Q^2$
 GeV^2



E Beam GeV	HRS Angle	Q ² GeV ²	Electron P (MeV)	Proton P(MeV)	Counts/day
			HRS	MPS	L= 10 ³⁸ Hz/cm ²
4	15.8	- 1	3320	1055	432,000
5	18.7	- 2	3788	1626	24,500
6	19.7	- 3	4254	2146	2390
8	17.0	- 4	5271	2646	849
10	15.2	- 5	7175	3136	386
11	19.0	- 10	8230	5520	12

E=6 GeV, LHRS(e), RHRS(p)

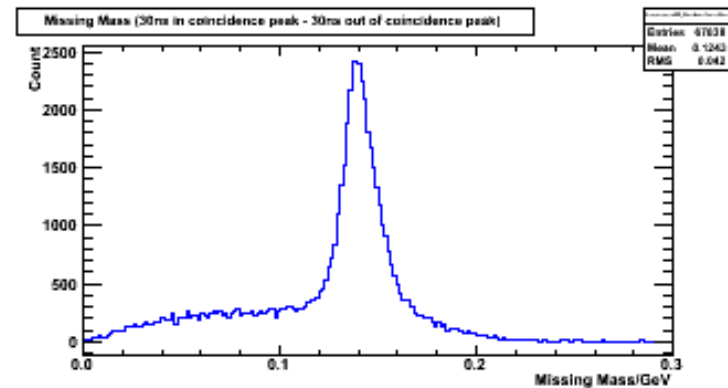
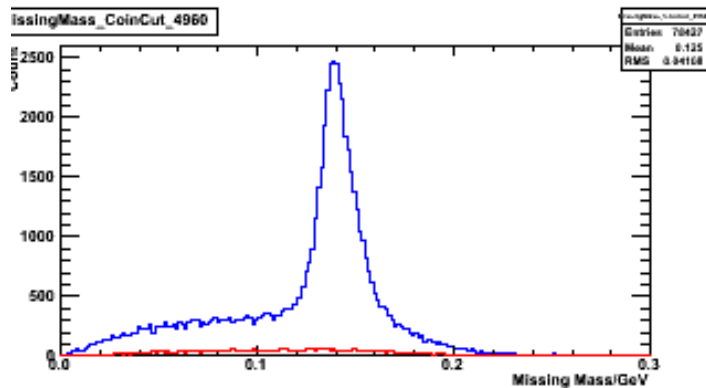
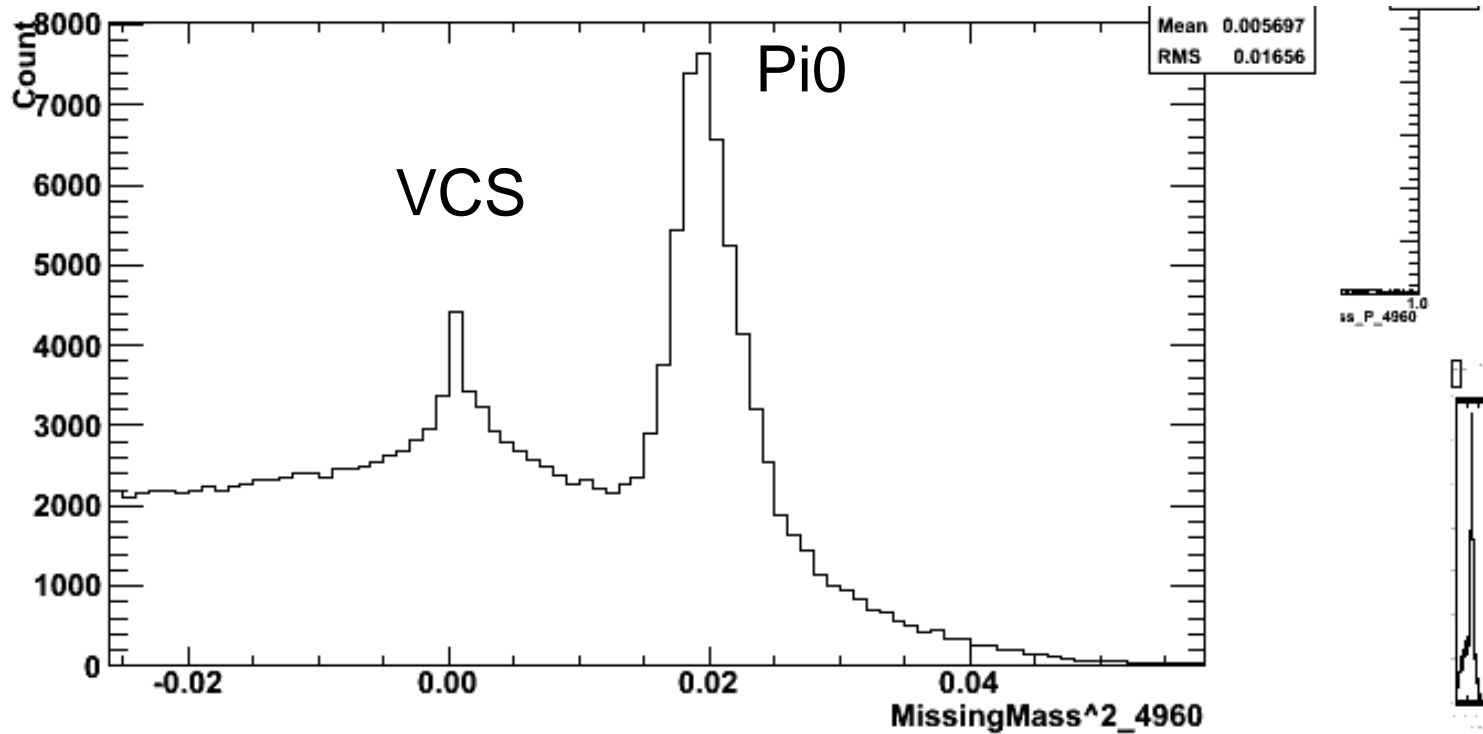


Questions

- What are the limits in resolution in W and how close to threshold do you have to be to expect these PQCD and ChPT predictions to be valid? How coarse or fine do you bin the data?
- How low in Q^2 can you go for PQCD to work?
- How do you get higher than $Q^2 = -3 \text{ GeV}^2$?
 - Electron energy too high
 - MPS resolution not good enough
 - For large angles rates are too low.

Data from E04-007

$$Q^2 = -0.45 \text{ GeV}^2$$



theory:

- detailed predictions possible near to threshold
- present LCSR accuracy probably 50%; can be improved to 20%(?)
- elastic πN rescattering can be taken into account
- can be incorporated in full PWA

- There exist two form factors related to multipoles
 - $S_{\{0+\}}$, and $L_{\{0+\}}$ and I expect that clean separation
 - between them is much easier if you have polarization
 - observables, similar to the situation with the proton
 - electric vs. magnetic form factors. Also, you will have
 - less problems with two-photon exchange contributions.
 - So, polarization is obviously a great plus.
-
- 3) I think you can say that the interesting issue is
 - how chiral symmetry works in hard reactions.

- For small $q^2 \sim m_\pi^2$ we have chiral perturbation theory and it works nicely (as it must). Question is
 - what happens when m_π is still treated as small parameter,
 - but q^2 is large, of order or larger than hadronic scale.
 - In this situation chiral perturbation theory is not applicable,
 - but people expect that the approach based on chiral algebra (soft pion theorems) still has to work up to a scale of
 - order $Q^2 \sim m_N^3/m_\pi$ whereafter QCD based approaches
 - are expected to become applicable. The task is to check
 - this scenario.
-
- As a bonus, if you find that soft-pion theorems
 - "work" for Q^2 of order 1 GeV or so, it will allow to extract
 - proton axial form factor from the data on π^+ production.

- Just a technical detail: chiral perturbation theory is good because
- it allows to calculate correction proportional to powers of the
- pion mass. When I say it is not applicable for $Q^2 \gg m_\pi^2$, this
- will imply that terms of order m_π^2 will remain out of control.
- This is what one loses. Chiral algebra, on the other hand,
- does need small q^2 , but remain selfconsistent to order m_π^1 .
- This means that for large q^2 you cannot expect theoretical accuracy
- better than smth like 5%, but this is not bad at all.

Hall C Data

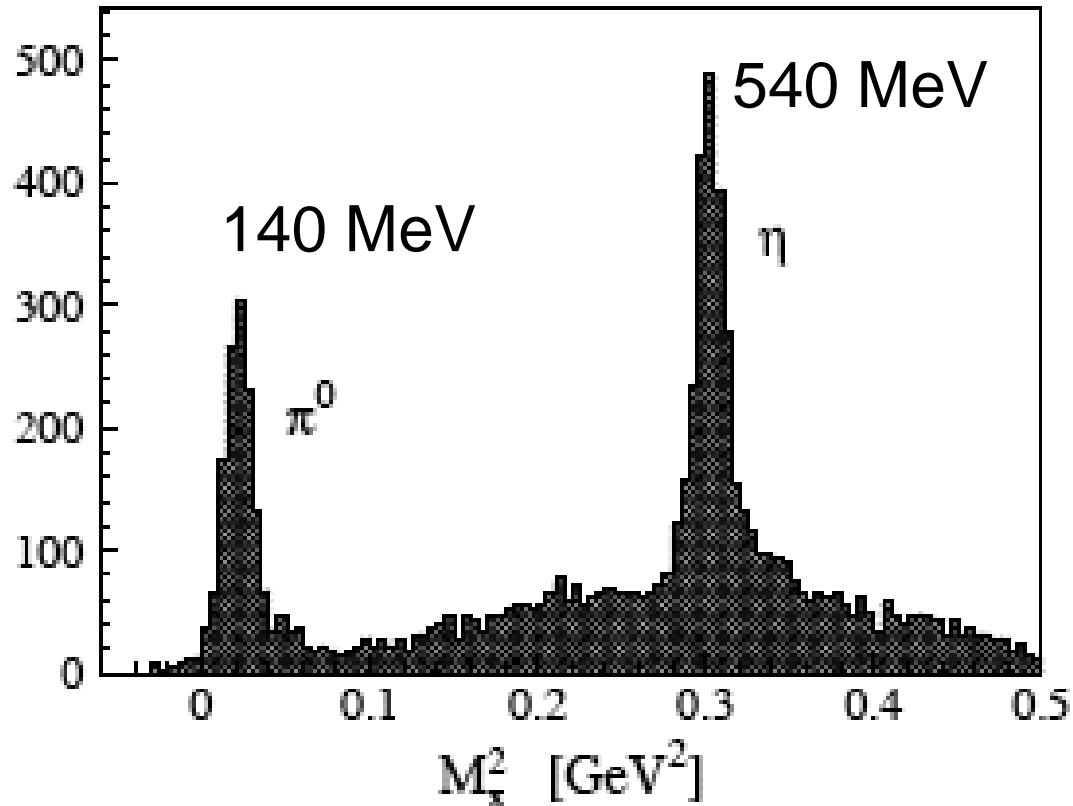


FIG. 1. Example of a missing mass distribution for the reaction $p(e, e'p)X$. The data are from one experimental setting at $Q^2 = 2.8 \text{ GeV}^2$ ($1.1 < W < 1.7 \text{ GeV}$).

V.V Frolov et al. PRL 82,45(1999)

$\theta_{central}$, degree	Ω , msr	D, meter	Hor. range, degree	Vert. range, degree
3.5	5	9.5	± 1.3	± 3.3
5.0	12	5.8	± 1.9	± 4.9
7.5	30	3.2	± 3	± 8
15	72	1.6	± 4.8	± 12.2
30	76	1.5	± 4.9	± 12.5

Table 1: The solid angle of MPS vs. central scattering angle.

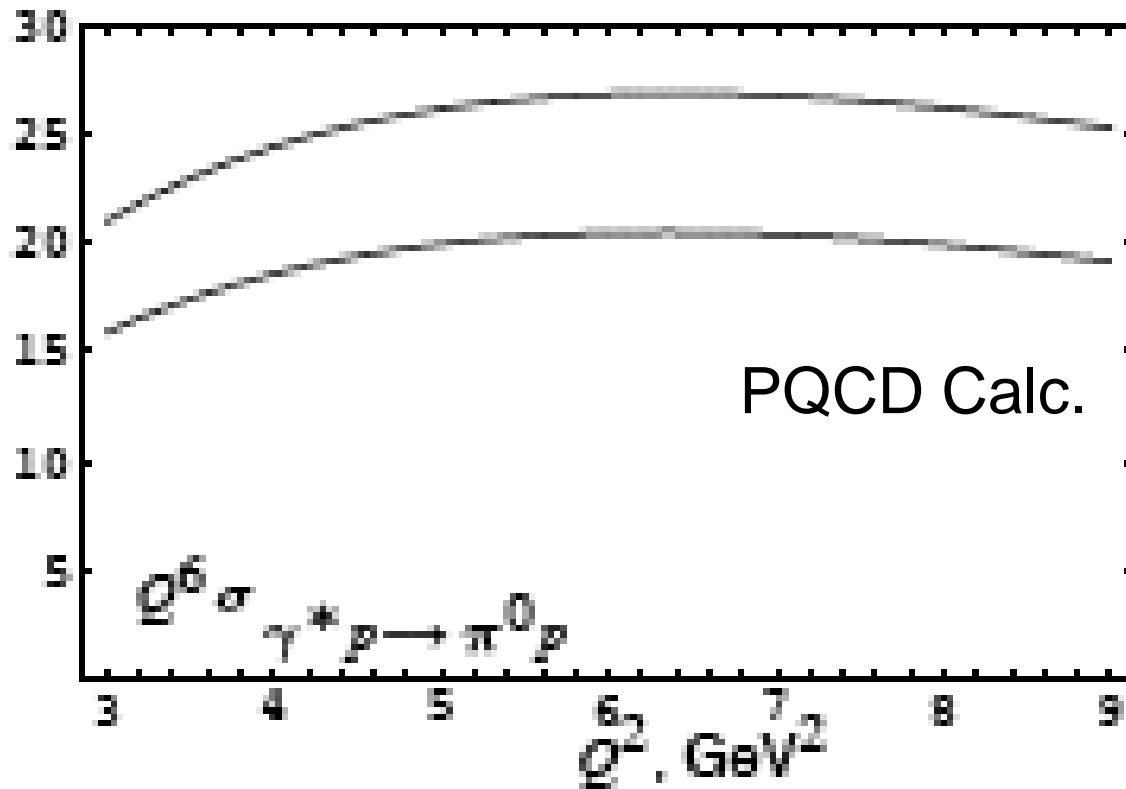


Figure: Integrated cross section $Q^6 \sigma_{\gamma^* p \rightarrow \pi^0 p}$ for $W = 1.11 \text{ GeV}$ (lower curve) and $W = 1.15 \text{ GeV}$ (upper curve)

$$W_{\text{threshold}} = 1.07326 \text{ GeV}$$

Threshold Pion Production in QCD

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based on

V.M. Braun, D. Ivanov, A. Lenz and A. Peters, Phys.Rev.D75:014021,2007

V.M. Braun, D. Ivanov and A. Peters, arXiv:0710.3265, PRD, to appear

JLAB, 02/19/08

$E=4$ GeV, LHRS(e), RHRS(p)

