## E04-007: π° Threshold Electroproduction Status Report

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Hall A Collaboration Meeting December 10 2012

# Since last collaboration meeting

- Khem Chirapatimol received PhD and returned to Thailand.
- Richard Lindgren presented first public results at JLAB Chiral Dynamics Workshop August 2012.
- U. Meisner agreed to extend HBChPT electroproduction calculation after seeing our data.
- A. Bernstein and Nikos Sparveris want to use E04-007 data to investigate VCS.
- Continue to work with Khem on systematic error analysis in preparation for publication.

# Motivation

- Chiral Perturbation Theory (ChPT): Low energy effective theory based on spontaneous breaking of chiral symmetry in QCD.
- Long range degrees of freedom: Goldstone boson (mesons) and baryons and their mutual interactions (tree + loops).
- Short distance physics: Low Energy Constants (LEC) whose number can be constrained by suitable choice of kinematics. Crucial question: At what energy (W) and distance (Q<sup>2</sup>) scale does chiral dynamics fail (need baryon/ vector meson resonance physics)?
- Precision measurement of (Q<sup>2</sup>,W) evolution of p(e,e'p)π<sup>o</sup> reaction near threshold can test validity of low energy expansion *once LECs are fixed*.
- Total cross section, polarization observables, EM multipoles.

# **Theory: Current Status**

### V. Bernard et al.: Heavy Baryon HBChPT

- Q<sup>2</sup>=0: O(p<sup>4</sup>) s-wave (2 LECs) O(p<sup>4</sup>) p-wave (3 LECs)
- Q<sup>2</sup>>0: O(p<sup>5</sup>) s-wave (5 LECs) O(p<sup>3</sup>) p-wave (1 LEC)
- M. Hilt: Relativistic BChPT
  - All  $Q^2$ :  $O(p^4)$  with five LECs, but better convergence.
- Gasparyan and Lutz: Unitary chiral EFT beyond threshold
  - Chiral Lagrangian truncated at  $O(p^3)$ .
  - Resonances enter thru infinite summation of counter terms.
- A. Bernstein: Unitary fit
  - No LECs. Fits s-wave cusp using unitarity *ansatz*.
- Phenomenological Fits to Global Data: DMT, MAID, SAID
  - Used to constrain ChPT fits.

### Photoproduction Data: CB-TAPS@MAMI-C $\gamma + p \rightarrow \pi^{0} + p$





Fourth order theory required to describe most recent MAMI photoproduction data.

Surprising result: relativistic BChPT departs from data at lower photon energies compared to HBChPT (see above plot).

Data: D. Hornidge *et al.* arXiv:1211.5495v1 Theory: Marius Hilt, Bosen 2011

 $\mathcal{O}(q^3)$ ,  $\mathcal{O}(q^3)$ +VM,  $\mathcal{O}(q^4)$ 

### **Electroproduction Data: MAMI**

### 2002

2009

Q<sup>2</sup> = 0.10: Distler et al. PRL 80, 2294 (1998) Q<sup>2</sup> = 0.05: Merkel et al. PRL 88, 1230 (2002)

#### Q<sup>2</sup> = 0.05 - 0.15: Merkel et al. arXiv:1109.5075



----- HBChPT (1996) ---- MAID Theory: BChPT, Marius Hilt, Bosen 2011

 $\mathcal{O}(q^3), \mathcal{O}(q^3) + VM, \mathcal{O}(q^4)$ 

Conflicting measurements require more extensive data set

## JLAB Experiment E04-007: Hall A

Experiment was designed to provide finer  $Q^2$  granularity and extend the W range up to 20 MeV above  $\pi^{o}$  threshold.

### E04-007: Coincidence Kinematics





Low mass and thickness minimizes straggling for low energy protons



### Data and Simulation $\Delta W$ =19-21 MeV

# **Experimental Challenges**

HRS Electron rate from rad. tail and  $(\pi^0, p)$ , at 1200 MeV and 12.5 Deg. 1.E+03 HRS Rate (Hz/MeV) 1.E+02 1.E+01 1.E+00 Radiation tail pi0,p 1.E-01 1070 1075 1080 1095 1085 1090 1100 W (MeV)







LHRS: Large singles rate from radiative tail. Limits beam current to < 5 µA.

**BigBite:** Very low momentum protons. Helium bags and special inset flange to permit thin Ti window w/o sacrificing vertical acceptance.

## Quality of Data Above Threshold Region



$$+\frac{p_{\pi}^{*}}{k_{\gamma}^{*}}\sqrt{2\varepsilon_{L}(1+\varepsilon)}(A_{0}^{LT}+A_{1}^{LT}P_{1}(x))(1-x^{2})^{1/2}\cos\phi_{\pi}^{*} \qquad \text{LT}$$
$$+\frac{p_{\pi}^{*}}{k_{\gamma}^{*}}\varepsilon A_{0}^{TT}(1-x^{2})\cos 2\phi_{\pi}^{*} \qquad \text{TT}$$

# Quality of Data at Threshold





E=1.192 GeV

E=0.880 GeV

# Quality of Data at Threshold $A_0^{T+L}$

![](_page_10_Figure_1.jpeg)

JLAB 2012 E=1.192 GeV ε=0.943 MAMI 2011 E=0.880 GeV ε=0.882

# **Dominant Systematic Errors**

### W Calibration and Resolution

![](_page_11_Figure_2.jpeg)

Knowledge of absolute W calibration and simulation of W resolution crucial to obtaining correct cross sections at threshold.

### Target Window Backgound Subtraction

![](_page_11_Figure_5.jpeg)

Full target yield used (no VZ cuts). Data with W <  $\pi^{\circ}$  threshold used to estimate contribution from windows.

# Effect on Total Cross Section of Shift in Energy Calibration

![](_page_12_Figure_1.jpeg)

![](_page_12_Picture_2.jpeg)

JLAB 2012 220 keV shift in W
JLAB 2012 Nominal W calibration

# Carbon Elastic Peak vs. Run No.

![](_page_13_Figure_1.jpeg)

## Example of Shift in LHRS W Calibration

![](_page_14_Figure_1.jpeg)

## Possible Sources of C Elastic Peak Shifts

- LHRS mis-pointing (shift in recoil factor R)
  - $dR/d\theta = 4-6 \times 10^{-4}/deg$  0.1° error = 40-60 keV
  - Exp't variation in L.tr.vz +/- 4 mm or +/- 0.03°
- LHRS magnetic field
  - Dipole stable. Q1 Hall probe stability ~ 200 ppm.
- Beam energy shifts
  - Kinematics uses Tiefenbach energy.
- Beam Y-position shifts
  - Fluctuations of 2-3 mm common in early runs.
- Software (Hall A Analyzer misconfigured).

### Sensitivity to Y-Position

The standard HRS transport matrix is listed at hallaweb.jlab.org/news/minutes/fo matrix.html:  $\begin{pmatrix} x \\ \theta \end{pmatrix}_{fp} = \begin{pmatrix} -2.18 & -0.198 & 11.9 \\ -0.10 & -0.469 & 1.967 \end{pmatrix} \begin{pmatrix} x_0 \\ \theta_0 \\ \theta_0 \\ z \end{pmatrix}$  where distances are in meters, angles in radians. In my notation, the solution for the ray's vertical angle and momentum at the target is  $\begin{pmatrix} \theta_0 \\ d \end{pmatrix} = \begin{pmatrix} -0.198 & 11.9 \\ -0.469 & 1.967 \end{pmatrix}^{-1} \begin{pmatrix} x \\ \theta \end{pmatrix}_{c_1} - \begin{pmatrix} -0.198 & 11.9 \\ -0.469 & 1.967 \end{pmatrix}^{-1} \begin{pmatrix} -2.18 \\ -0.10 \end{pmatrix} x_0$ or  $\begin{pmatrix} \theta_0 \\ d \end{pmatrix} = \begin{pmatrix} 0.379 & -2.229 \\ 0.090 & -0.038 \end{pmatrix} \begin{pmatrix} x \\ \theta \end{pmatrix}_{\ell_0} - \begin{pmatrix} -1.055 \\ -0.201 \end{pmatrix} x_0$ To first order, the momentum is thus  $P-P_{central} = P_{central} d = P_{central} \{0.090 \text{ x} - 0.038 \theta - 0.201 \text{ x}_0\}$ (G. Franklin, Aug. 18, 2010)  $\times 10^{-2}$ 0.34 0.32  $\chi^2/ndf$ 96.03 96 0.3 Constant 322.0 URB.Y 350 0.2669E-02 Mean Siama 0.2305E-03 URB.Y 300 250 200 0.24 150 0.22 100  $\sigma = 230 \,\mu$ 50 0.2 -12 -110.26 0.28 200.PROY 0.22 0.24 0.3 0.32 0.34

x 10

Since raster was not used raster-y corrections were turned off.

This means our W calibration was vulnerable to shifts in vertical beam position > 0.5 mm.

RUN 3829 - CARBON - PO=1193.98 MEV

![](_page_16_Figure_5.jpeg)

## **Carbon Elastic Cross Sections**

![](_page_17_Figure_1.jpeg)

Cross section is sensitive to pointing errors, but not so much to energy calibration.

Fourier-Bessel fit to NIKHEF-K data taken from Offermann et al. and recalculated for E=1.192 GeV using DWBA phase-shift code from J. Heisenberg (R. Lindgren).

## Carbon Elastic – Ratio with FB Fit

![](_page_18_Figure_1.jpeg)

From Offermann et al.

Significant discrepancies only for 16.5° data near diffraction minimum. Similar to enhancements seen in previous measurements at lower beam energies.

# Summary and Still to Do

- Large data set with unprecedented W, Q<sup>2</sup>, C.M. coverage and many available chiral theories to test.
- Estimates of systematic errors due to W calibration and target background subtraction.
- Compare pi0 yields for similar kinematics having large C elastic peak shifts to estimate systematics.
- Start writing paper in early 2013. Chiral Dynamics Workshop writeup due in Jan 2013.

# Carbon Runs: Q1 and VZ stability

![](_page_20_Figure_1.jpeg)

![](_page_21_Figure_0.jpeg)

## $Q^2$ Dependence of $A_0^{TT}$

LEGENDRE COEFFICIENT  $A_{o}^{TT}(\mu b)$ 

![](_page_22_Figure_2.jpeg)

Both Q<sup>2</sup> and W dependence of  $\sigma_{TT}$  in strong disagreement with  $O(q^3)$  ChPT. Changing  $b_p$  LEC to compensate may destroy agreement with other p-wave observables.

![](_page_22_Picture_4.jpeg)

![](_page_22_Picture_5.jpeg)

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