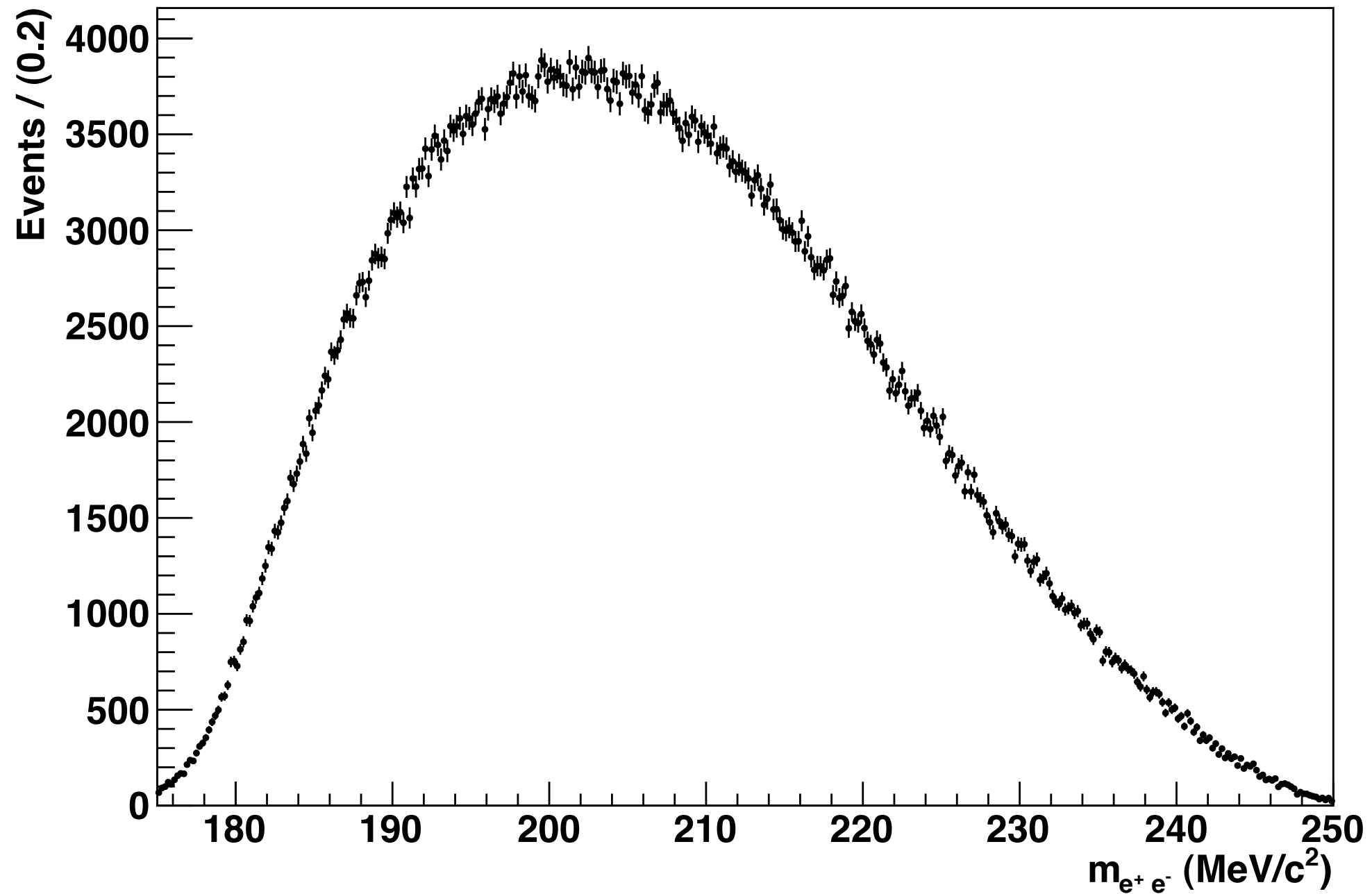


APEX Resonance Search

APEX Test Run Data, ~770K events



James Beacham
New York University



Outline

1) APEX resonance search

- Objectives of resonance search
- Procedure
- Optimization for binned analysis
- Characteristics of limit plots
- Test run unblinding and final results
- Plans for full run analysis

2) In progress: Update of reach calculation for full run using new acceptance

The Objectives

- Identify the small, narrow A' resonance on a smooth, high-statistics background spectrum
- Determine an upper limit on the number of signal events, S , consistent with background
- Upper limit on $S \rightarrow$ upper limit on coupling

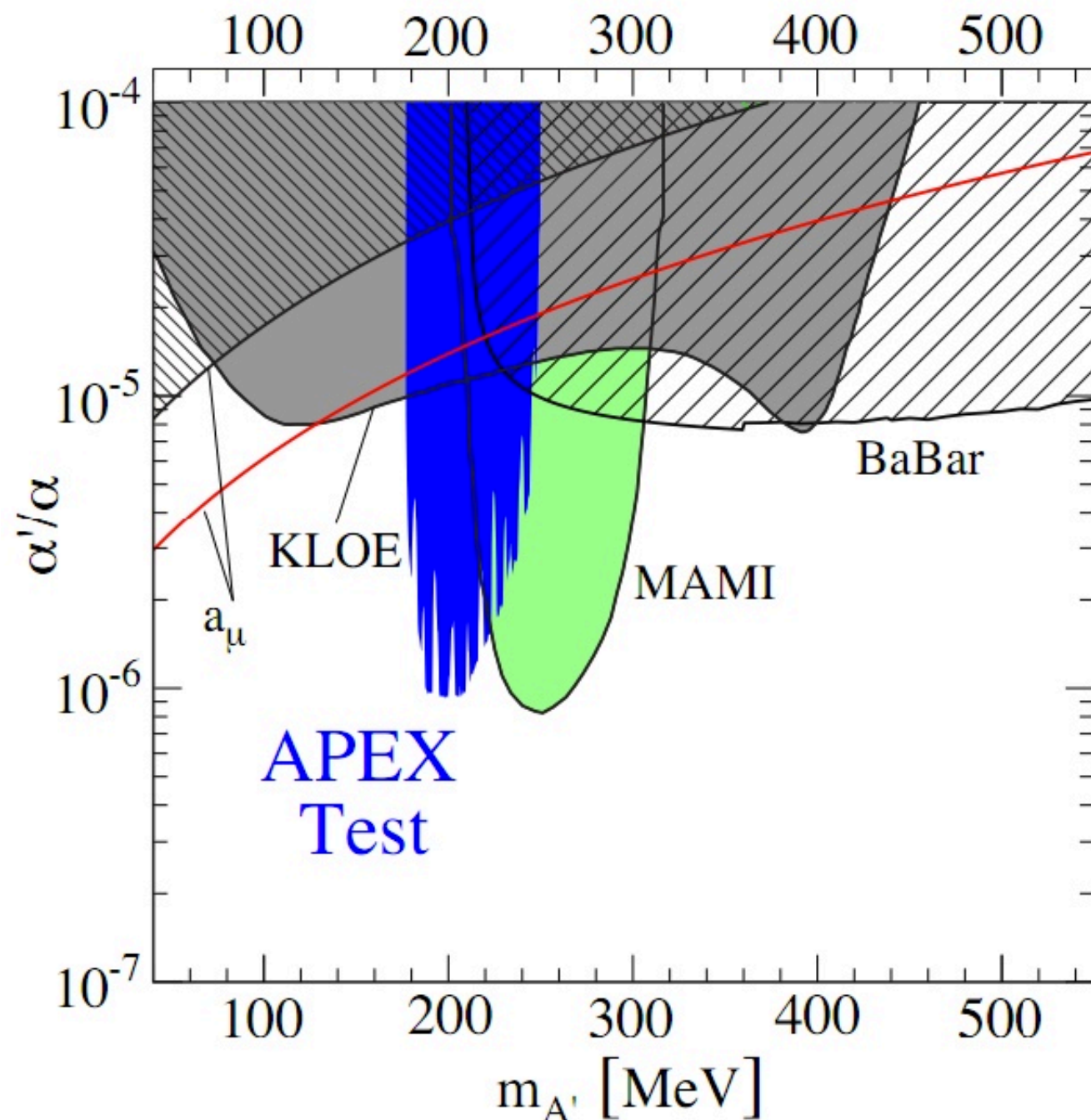
$$\frac{d\sigma(A')}{d\sigma(\gamma^*)} = \left(\frac{3\pi\epsilon^2}{2N_{\text{eff}}\alpha} \right) \frac{m_{A'}}{\delta m} = \frac{S_{\delta m}}{B_{\delta m}^{\gamma^*}}$$

(See APEX proposal)

Normalize all backgrounds to γ^* background

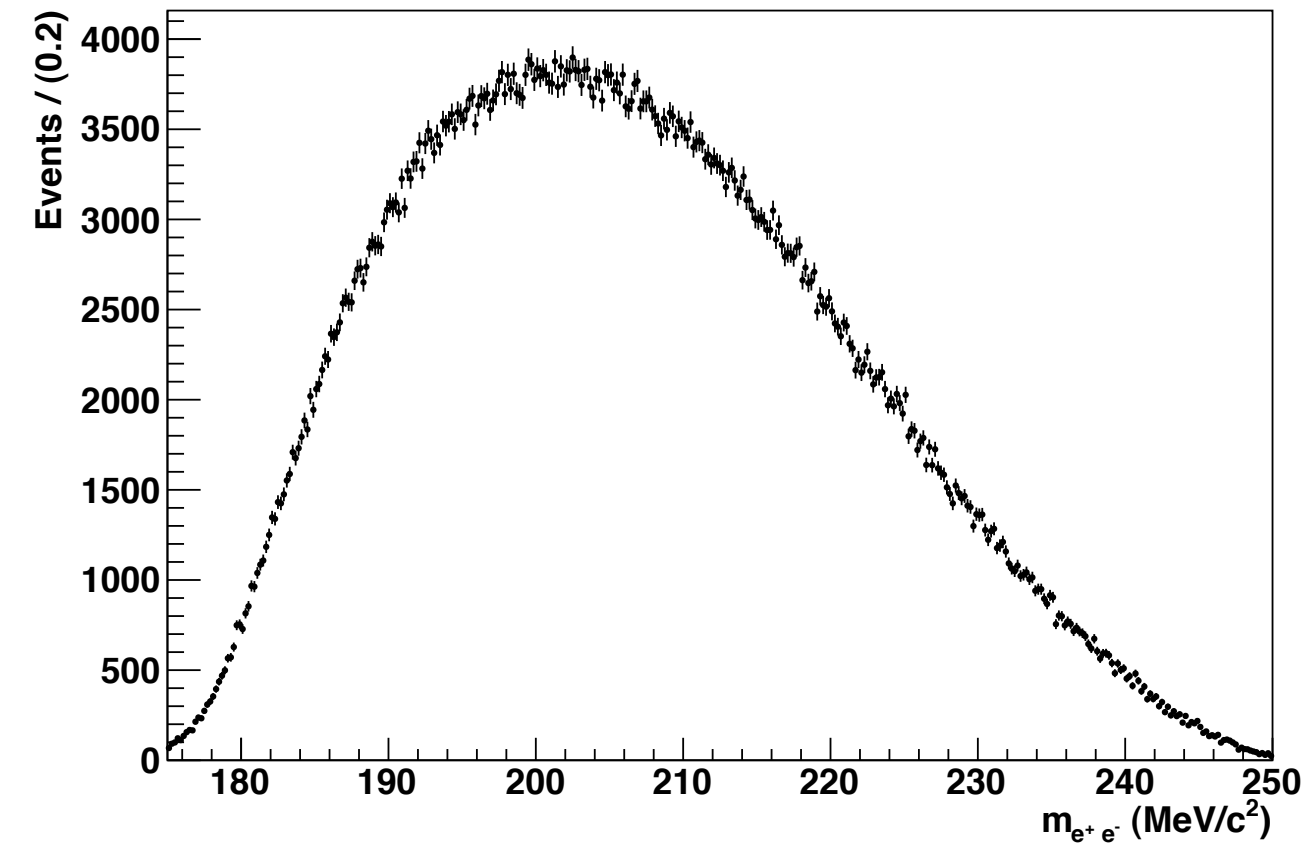
- Ratio f of radiative-only cross section to full trident cross section determined via Monte Carlo to vary linearly from 0.21 to 0.25 across APEX mass range

$$\left(\frac{\alpha'}{\alpha} \right)_{\text{max}} = \left(\frac{S_{\text{max}}/m_{A'}}{f \cdot \Delta B / \Delta m} \right) \times \left(\frac{2N_{\text{eff}}\alpha}{3\pi} \right)$$





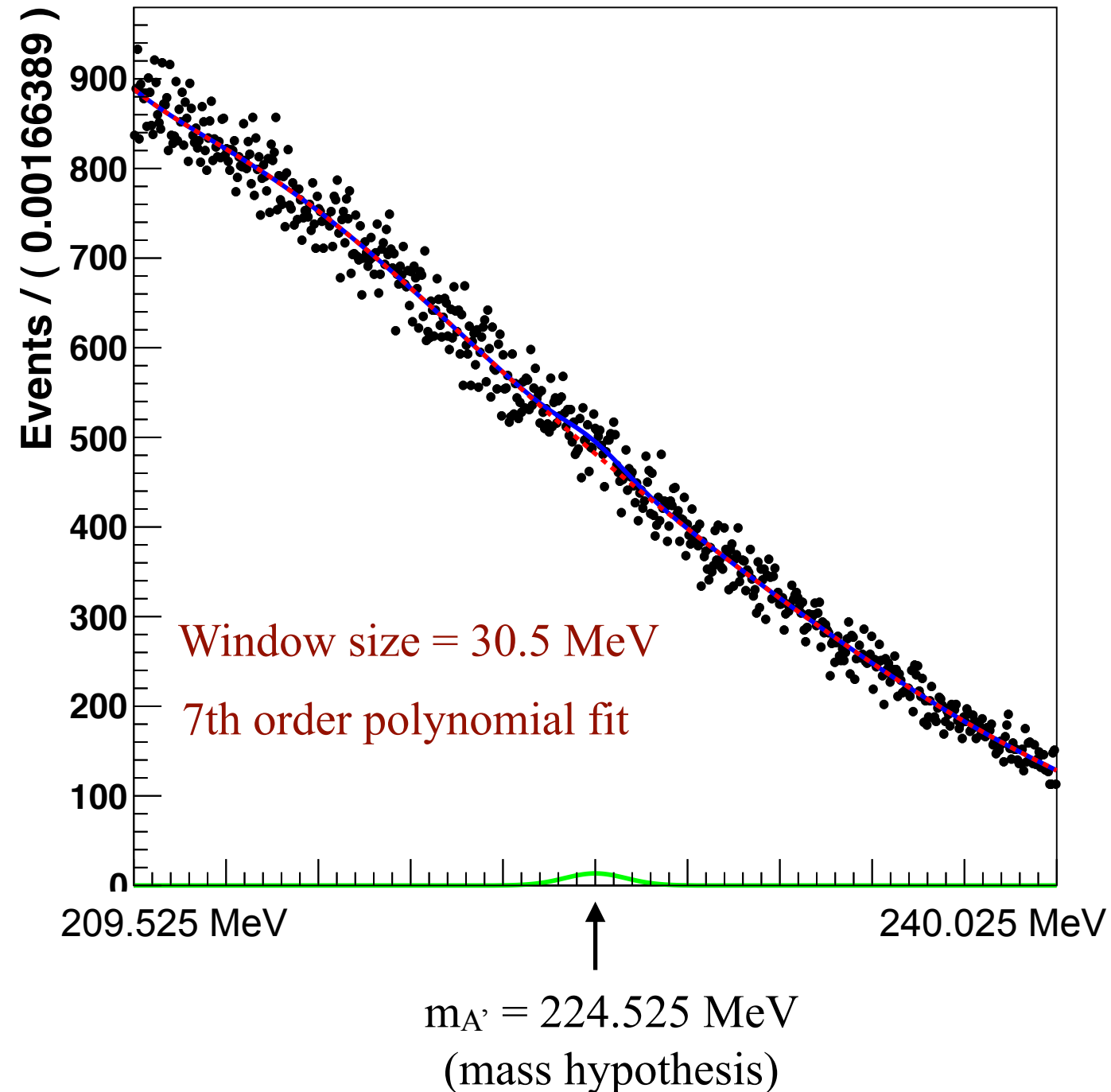
The Procedure



Procedure: Scan spectrum; look for significant excesses; determine upper limit on number of signal events, S , at each mass hypothesis.

Challenge: Very high statistics ($\sim 770,000$ events); regular raster scan methods (linear or quadratic background models) may be too crude.

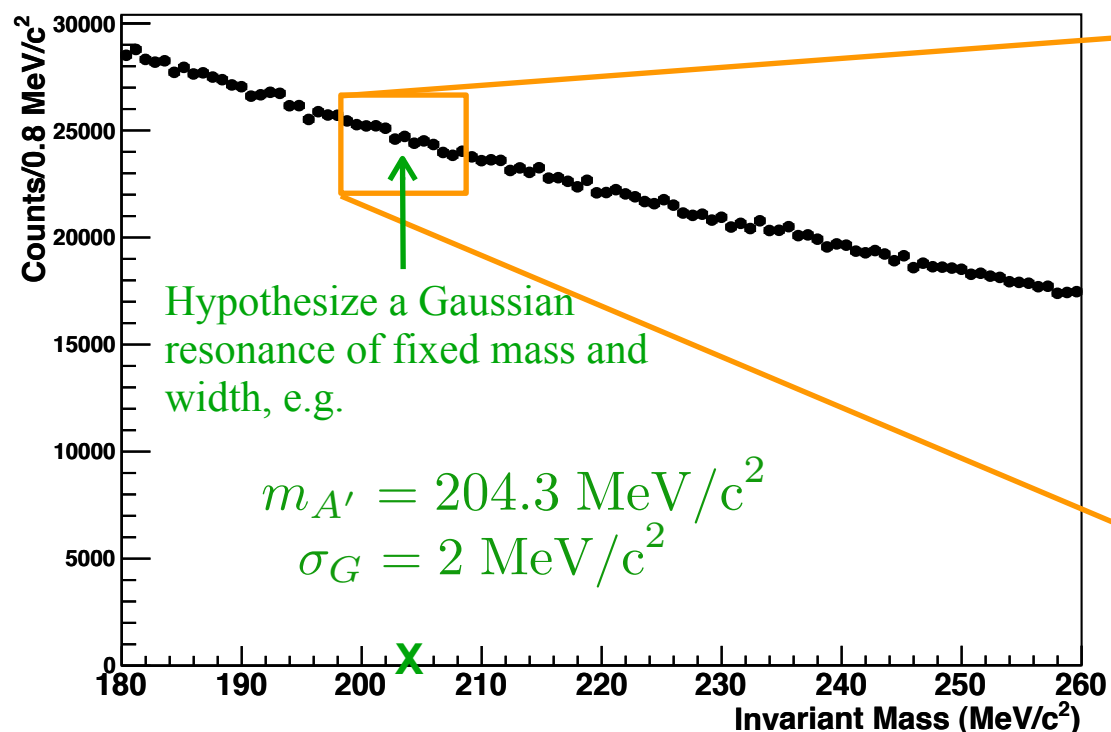
Final choice for test run: 30.5 MeV wide window and 7th order polynomial background fit.



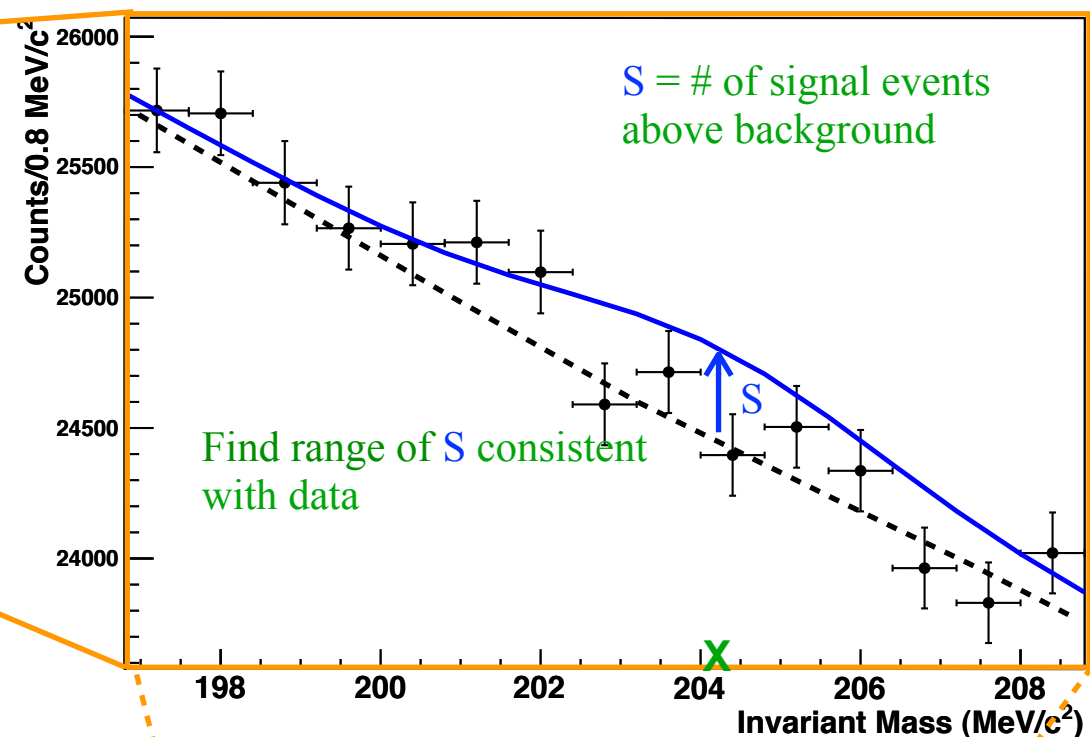


START: $S = 0$

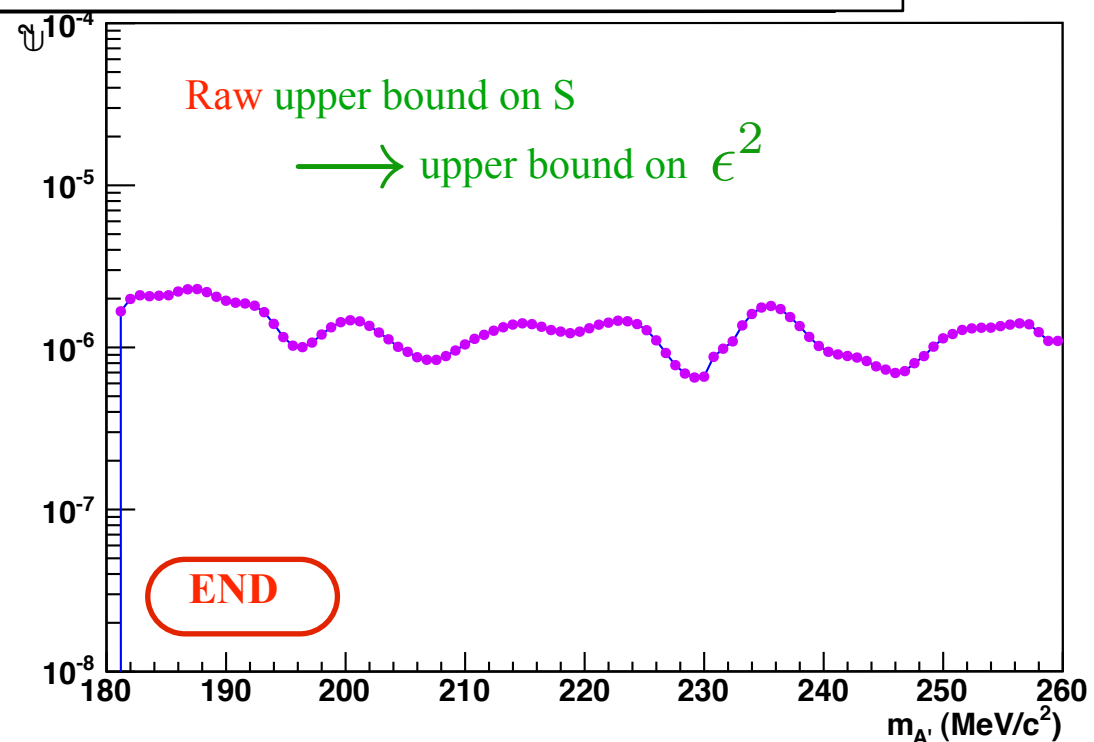
Toy model: Coincident e^+e^- pairs



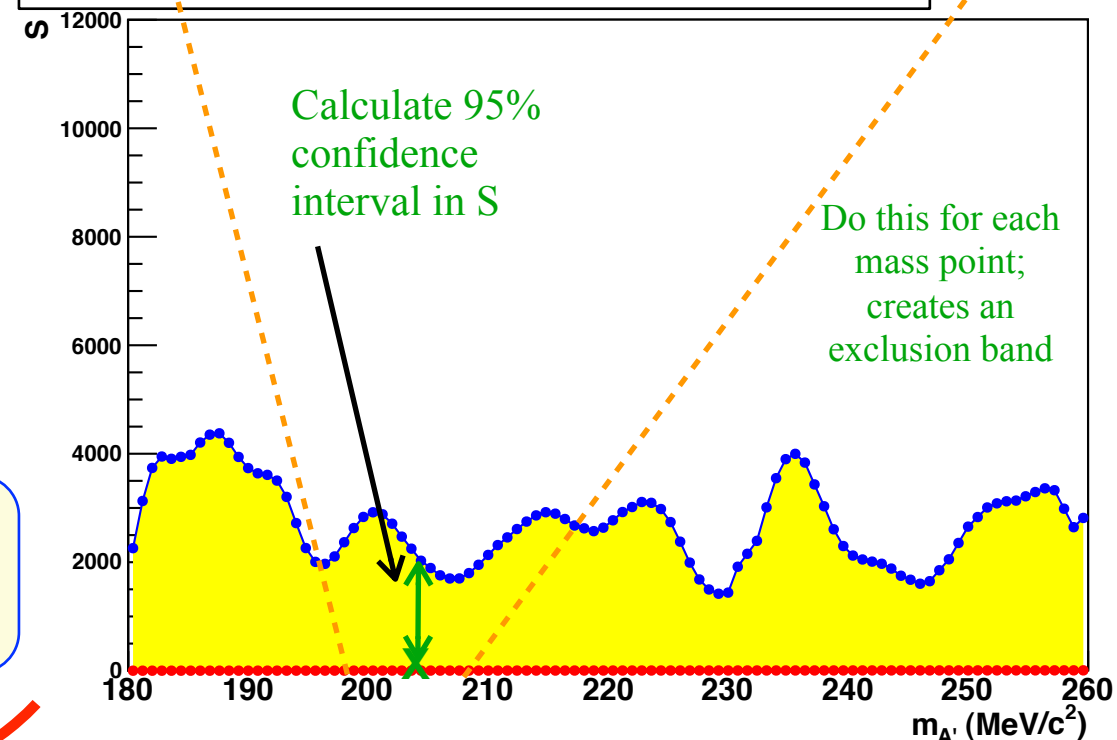
Toy model: Coincident e^+e^- pairs



Toy model: Upper bound on ϵ^2



Toy model: confidence intervals



One more step...



Test statistic based on **Profile Likelihood Ratio**

Probability model:

$$P(m_{e^+e^-} \mid m_{A'}, \sigma, S, B, a_i) = \frac{S \cdot N(m_{e^+e^-} \mid m_{A'}, \sigma) + B \cdot \text{Polynomial}(m_{e^+e^-}, a_i)}{S + B}$$

This becomes a likelihood function, L , as a function of the model parameters.

▸ Separate parameters into

- Parameters of interest (physics parameters: mass, # signal events, etc.)
- Nuisance parameters (systematics, i.e., everything else: background shape, etc.)

A test statistic that incorporates nuisance parameters is the **profile likelihood ratio**:

The PLR is the maximum likelihood under the null hypothesis as a fraction of its largest possible value; large values of λ indicate that the null is reasonable

Parameters of interest (POI),
set to values for null
hypothesis ($S = 0$)

$$\lambda(S) = \frac{L(S, \hat{\hat{B}}, \hat{\hat{a}}_i)}{L(\hat{S}, \hat{B}, \hat{a}_i)}$$

Conditional MLE of
nuisance parameters
(i.e., for $S = 0$)

MLE of POI

MLE of nuisance
parameters

Test statistic is then $-2 \log \lambda$

Wilks's Theorem: $-2 \log \lambda$ is distributed as a χ^2
with # d.o.f. = # POI

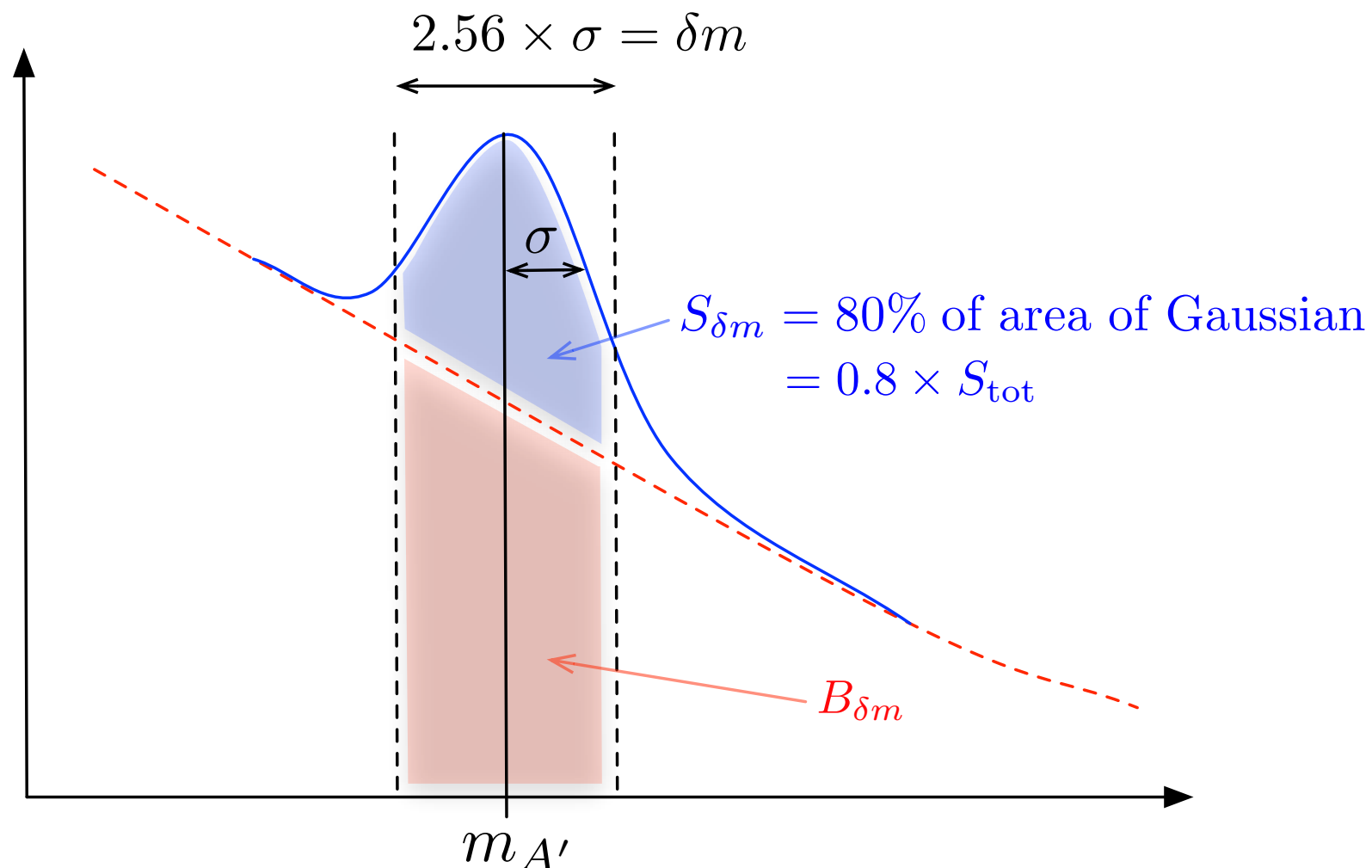
Upper limit on S yields an upper limit on ϵ^2

$$\epsilon^2(m_{A'}) = \frac{S_{\delta m}}{B_{\delta m}} \frac{F(m_{A'})}{m_{A'}} \delta m \left(\frac{2N_{\text{eff}}\alpha}{3\pi} \right)$$

δm = mass window = $2.56 \times \sigma$ σ = width of Gaussian ~ 1 MeV (mass-dependent)

Background factor: $F \sim 5$ (mass-dependent)

$N_{\text{eff}} = 1$

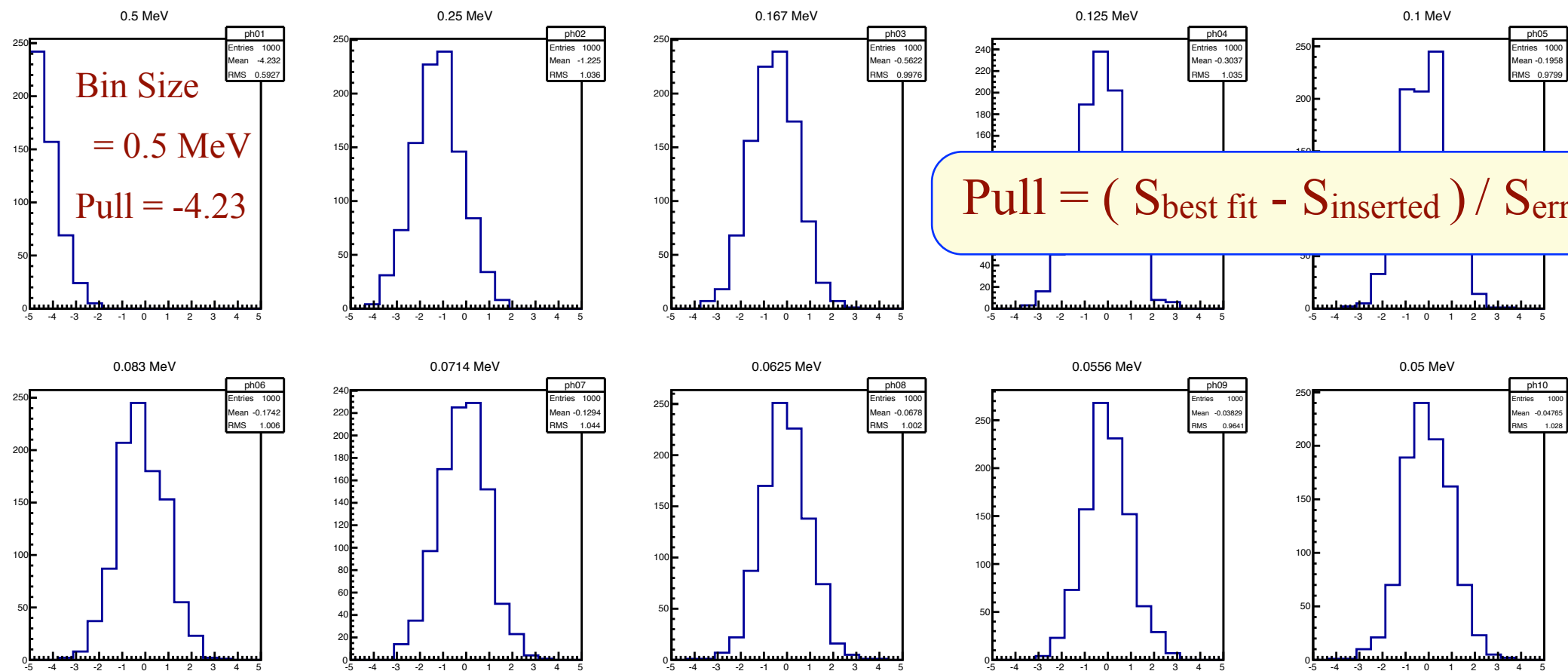


Setting limit on S: Binned analysis

We would ideally do an unbinned analysis; such a large number of events makes this intractably time consuming. A binned analysis is manageable and any pulls due to choice of statistical tool (ROOT, Mathematica, etc.) can be made negligible by choosing small enough bins.

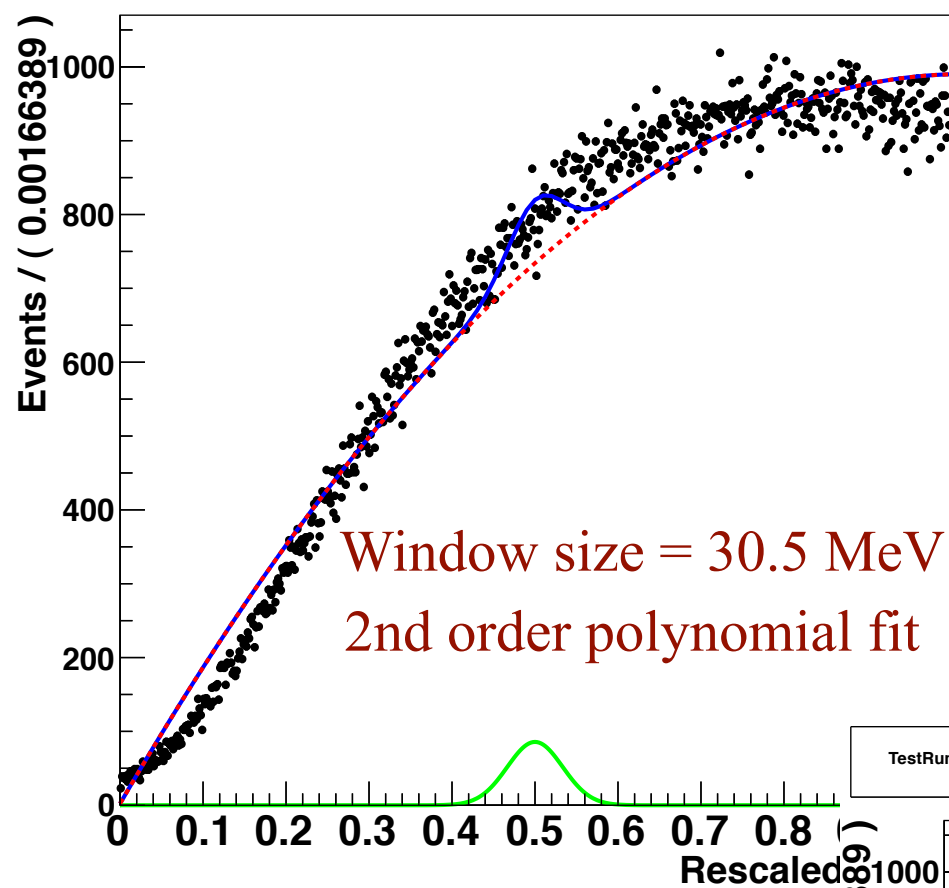
Example: With ROOT, do ensemble studies at a representative mass point; pseudodata sets based on spectrum of 10% of test run data, scaled up

Here, $S_{\text{inserted}} = 0$; large pulls for 0.5 MeV bins; pulls negligible at 0.06/0.05 MeV bins



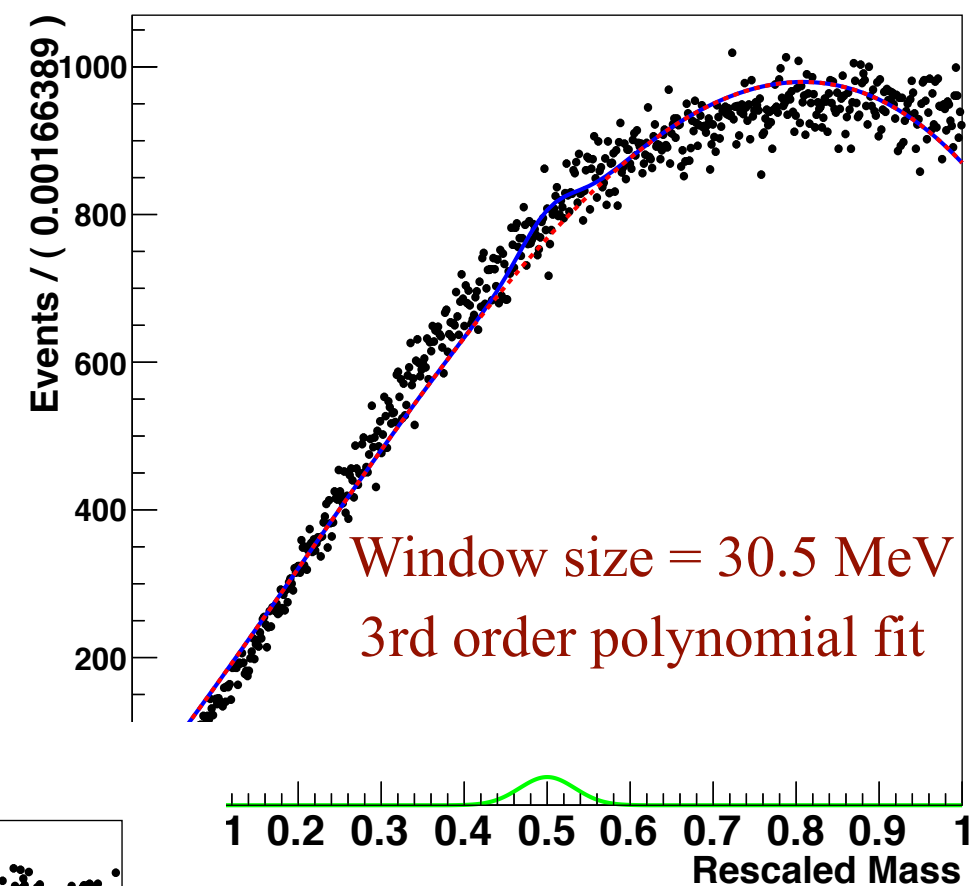


TestRunData_6_30_10_27, Wind 30, Cheb O(2), $m_A = 191.025$, $edm = 0.000001$, CQ=3

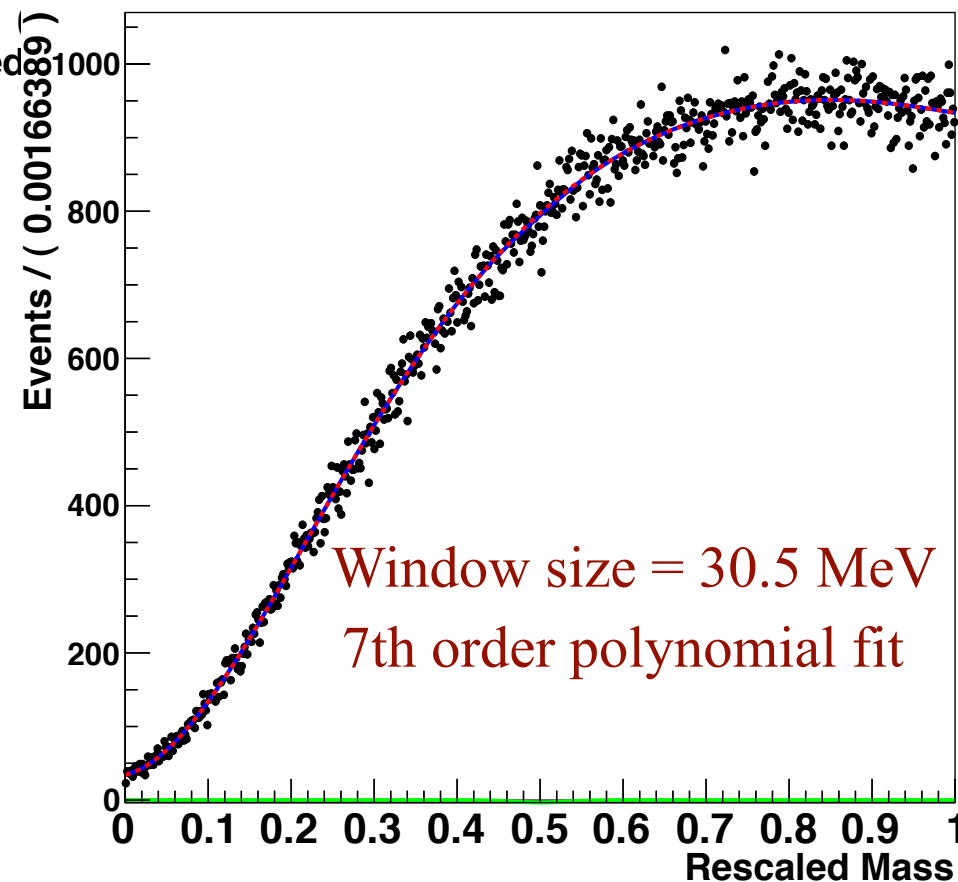


Lower order
polynomials fail to
capture the background
shape and lead to
spurious peaks

TestRunData_6_30_10_27, Wind 30, Cheb O(3), $m_A = 191.025$, $edm = 0.000173$, CQ=3



TestRunData_6_30_10_27, Wind 30, Cheb O(7), $m_A = 191.025$, $edm = 0.000299$, CQ=3



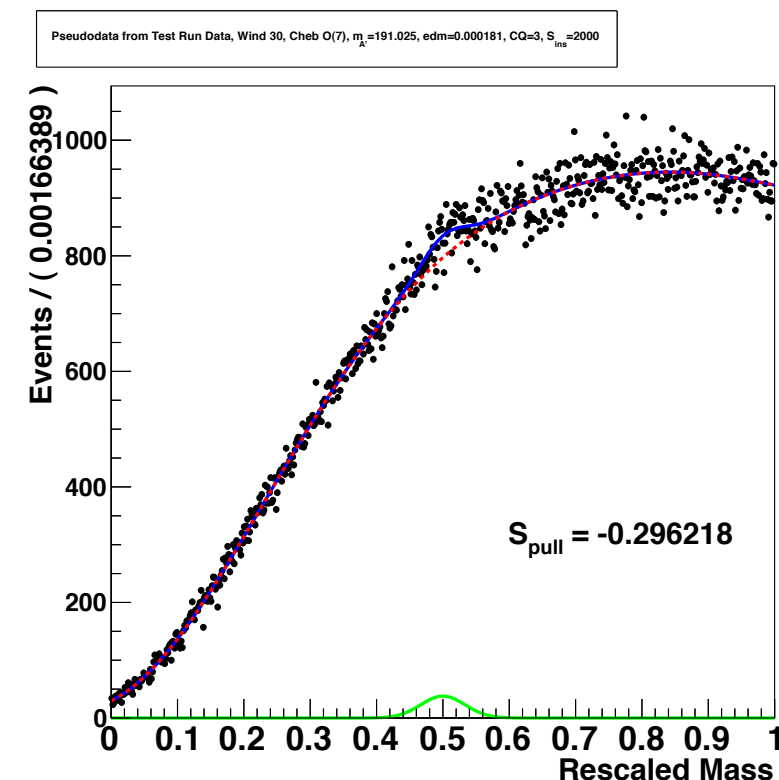
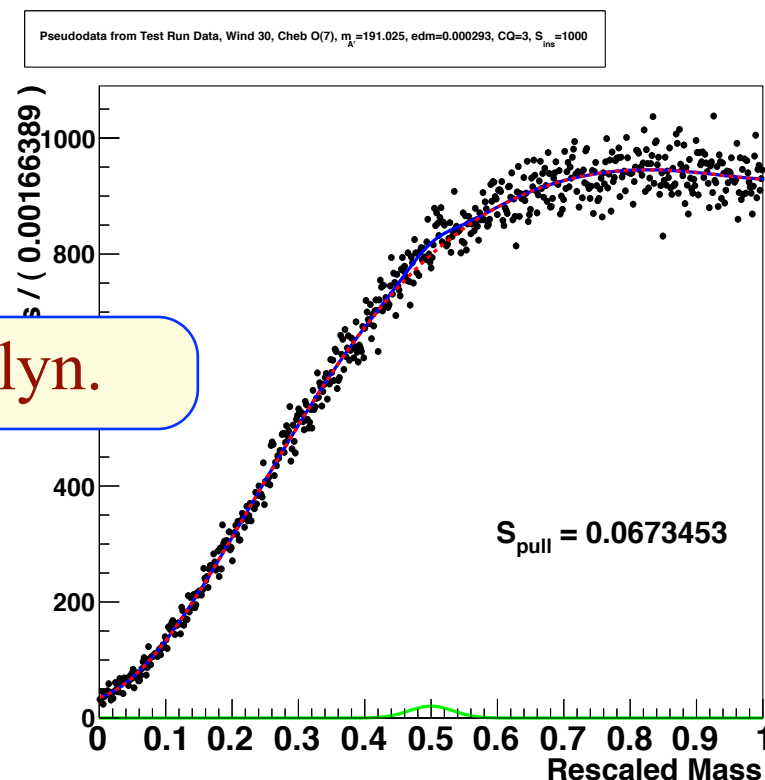
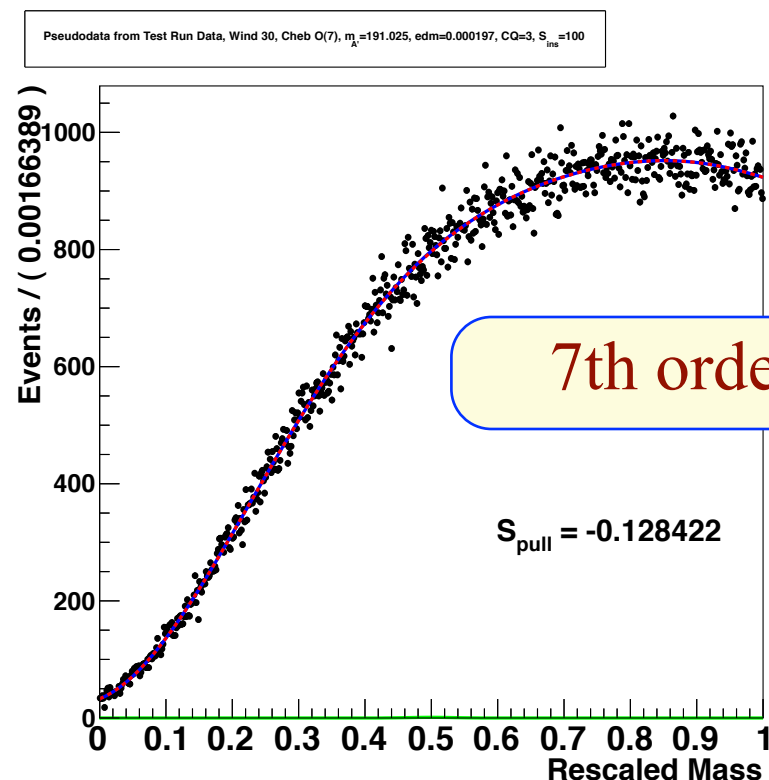
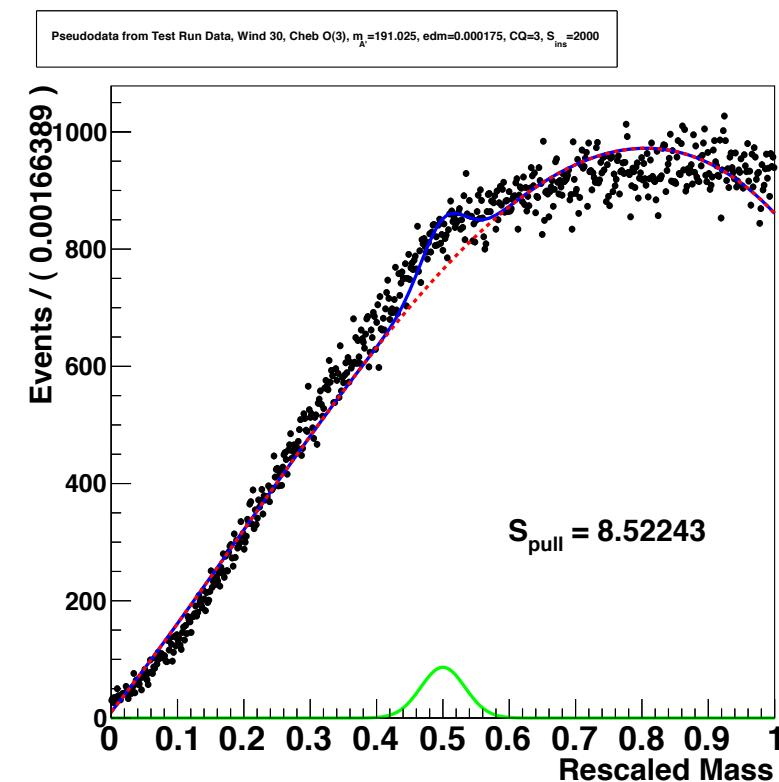
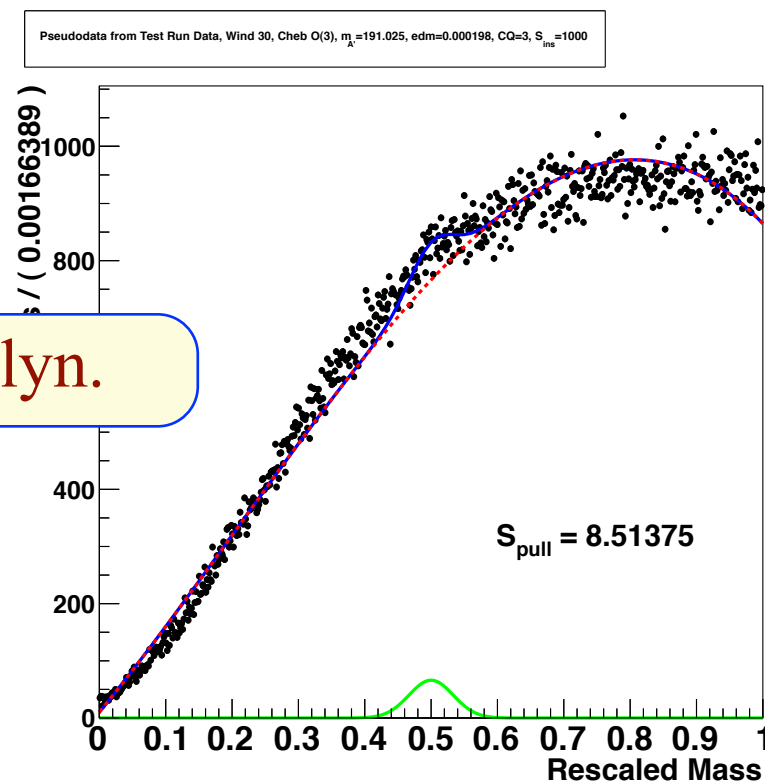
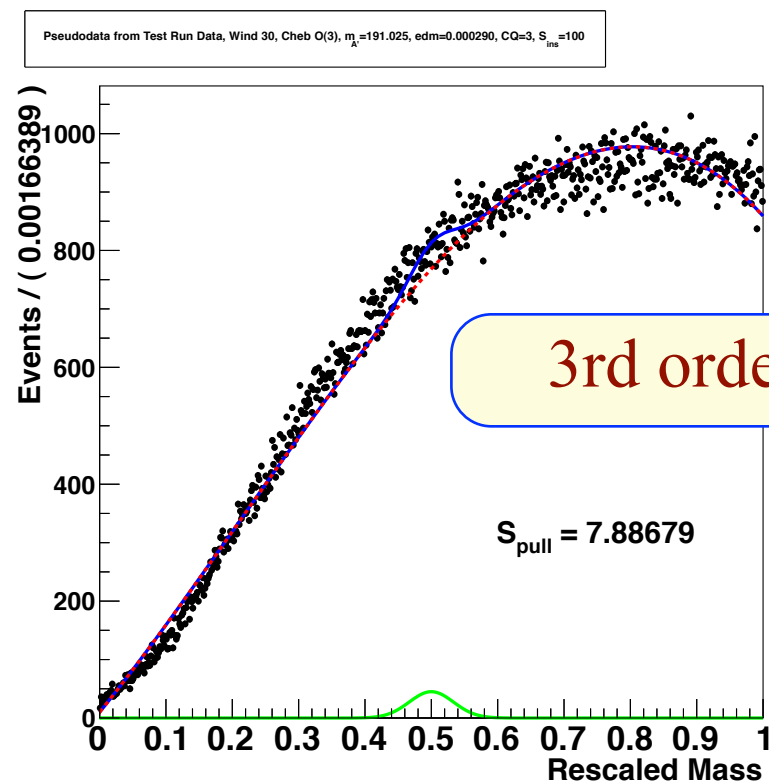
$S_{\text{inserted}} = 0$



$$S_{\text{inserted}} = 100$$

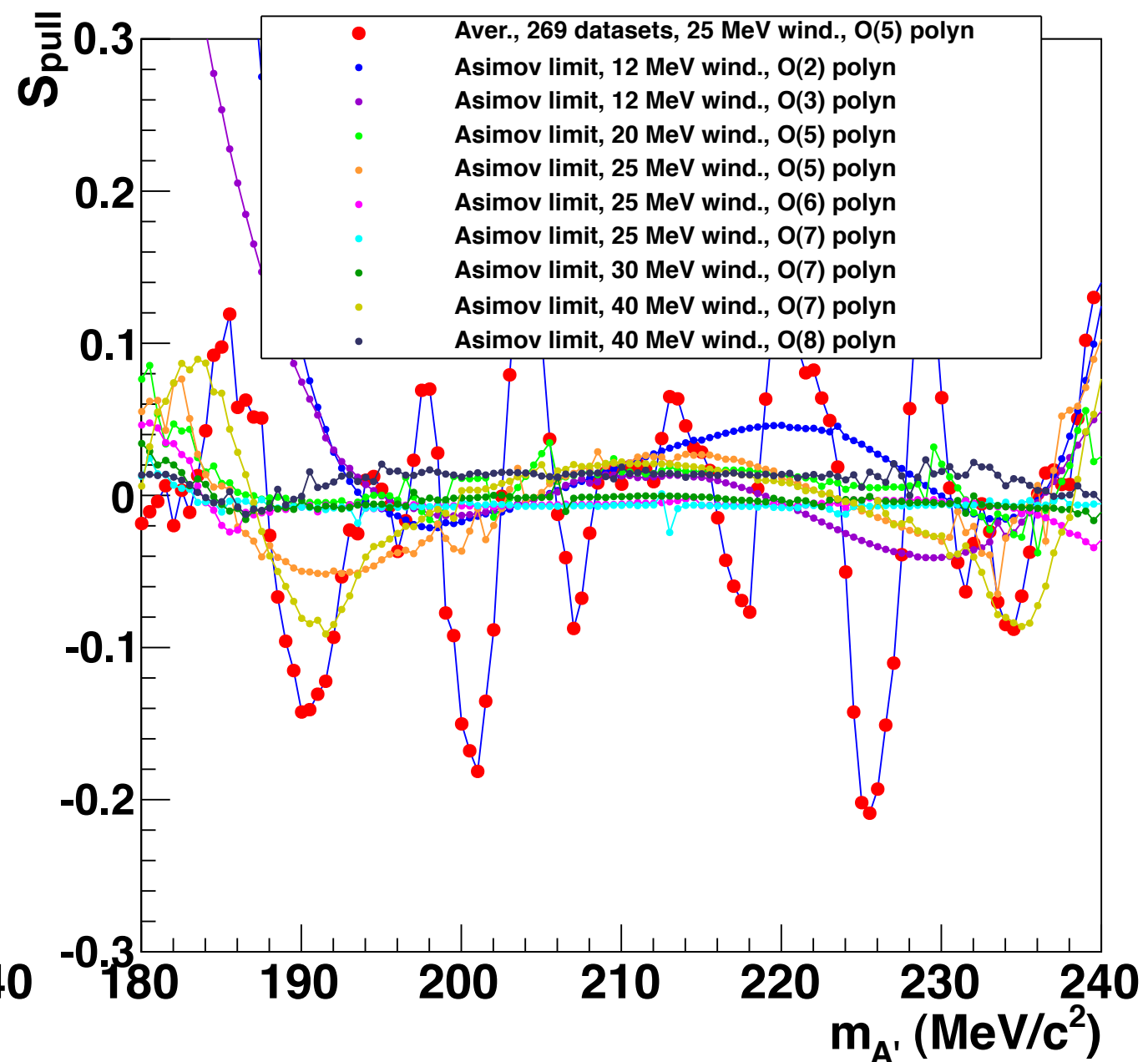
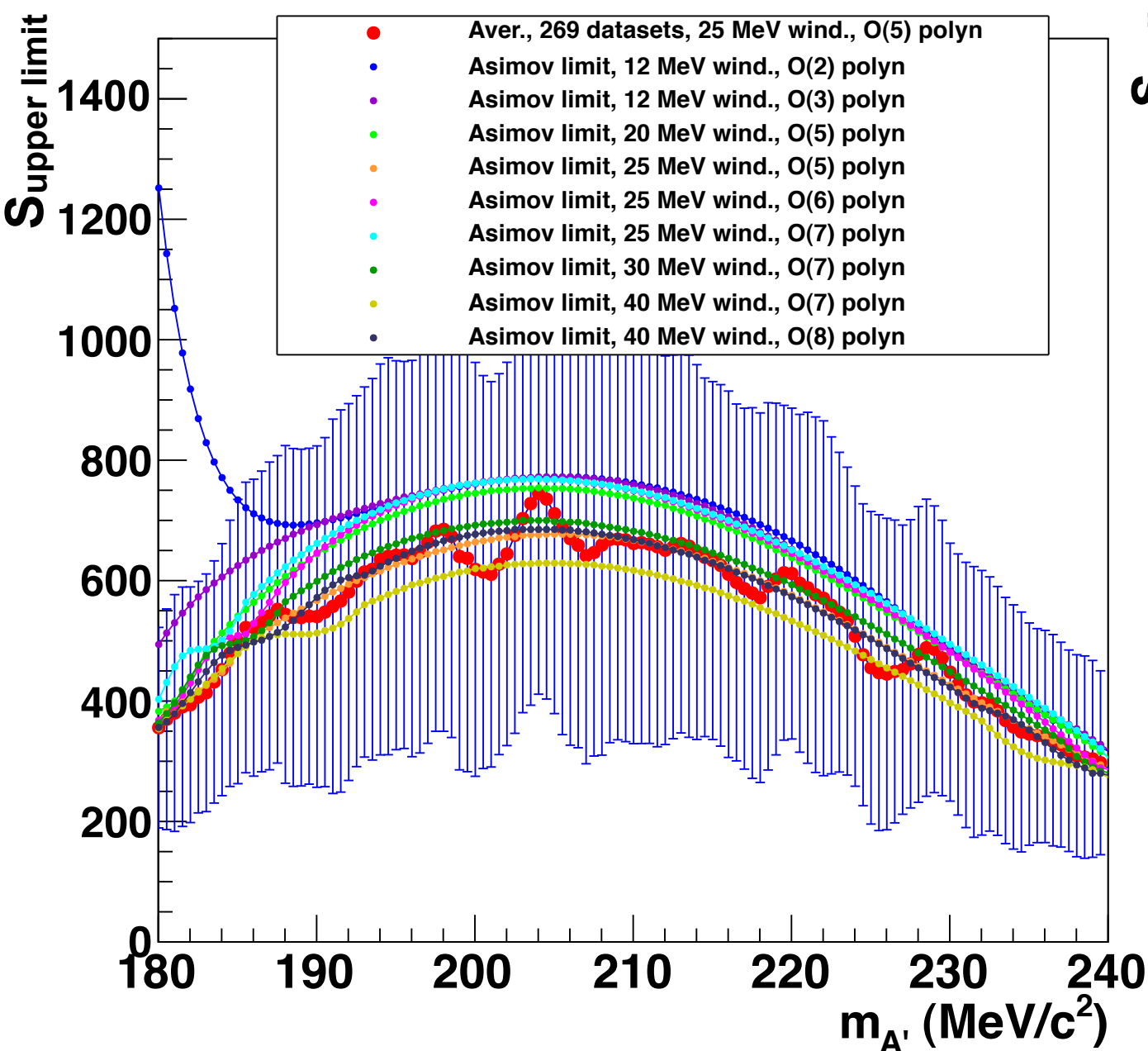
$$S_{\text{inserted}} = 1000$$

$$S_{\text{inserted}} = 2000$$



Representative pseudodata sets; not an ensemble

Setting limit on S: Window size and polynomial order



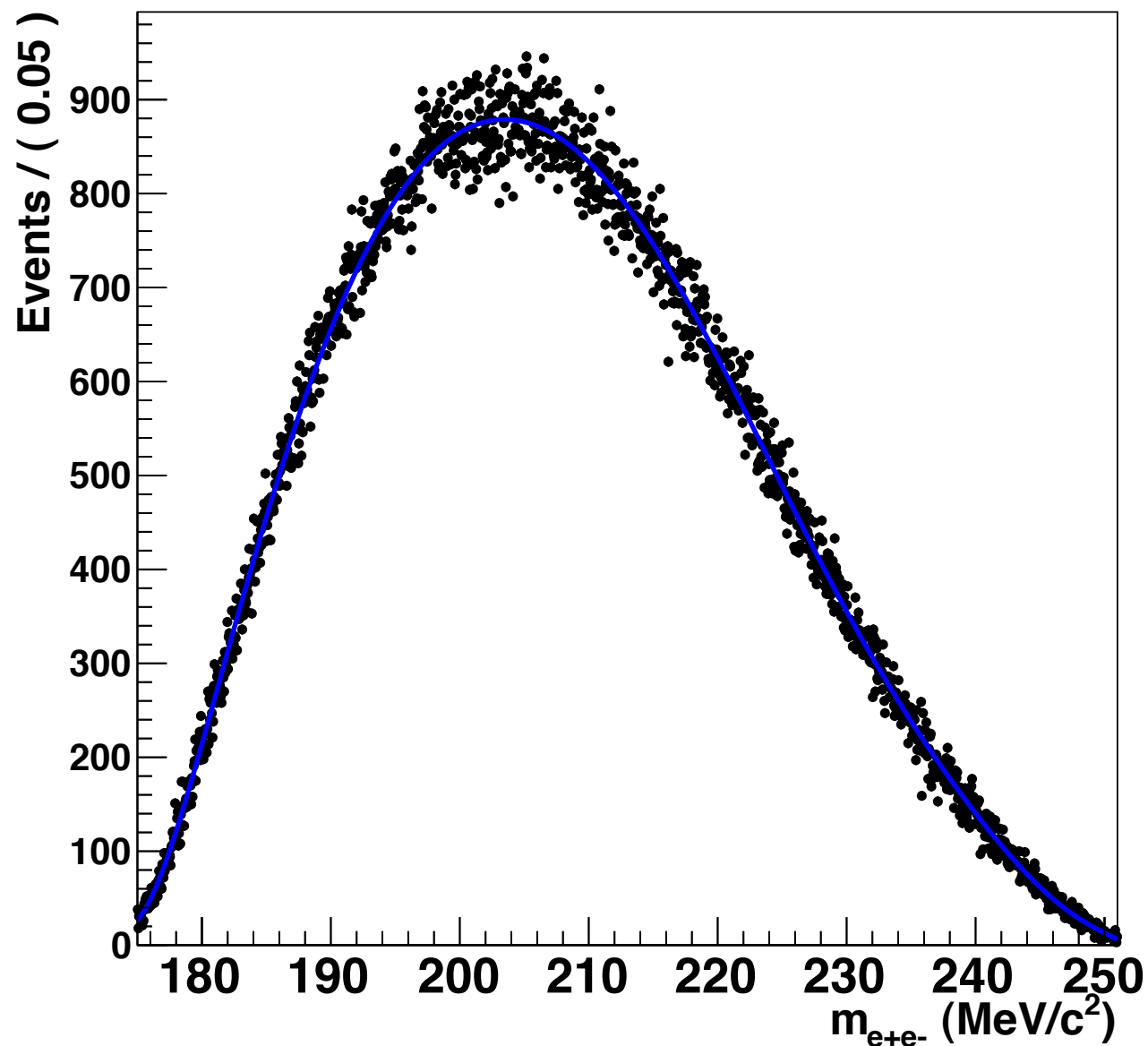
30 MeV window with 7th order polynomial (dark green line)
optimizes between sensitivity in S and minimal pull



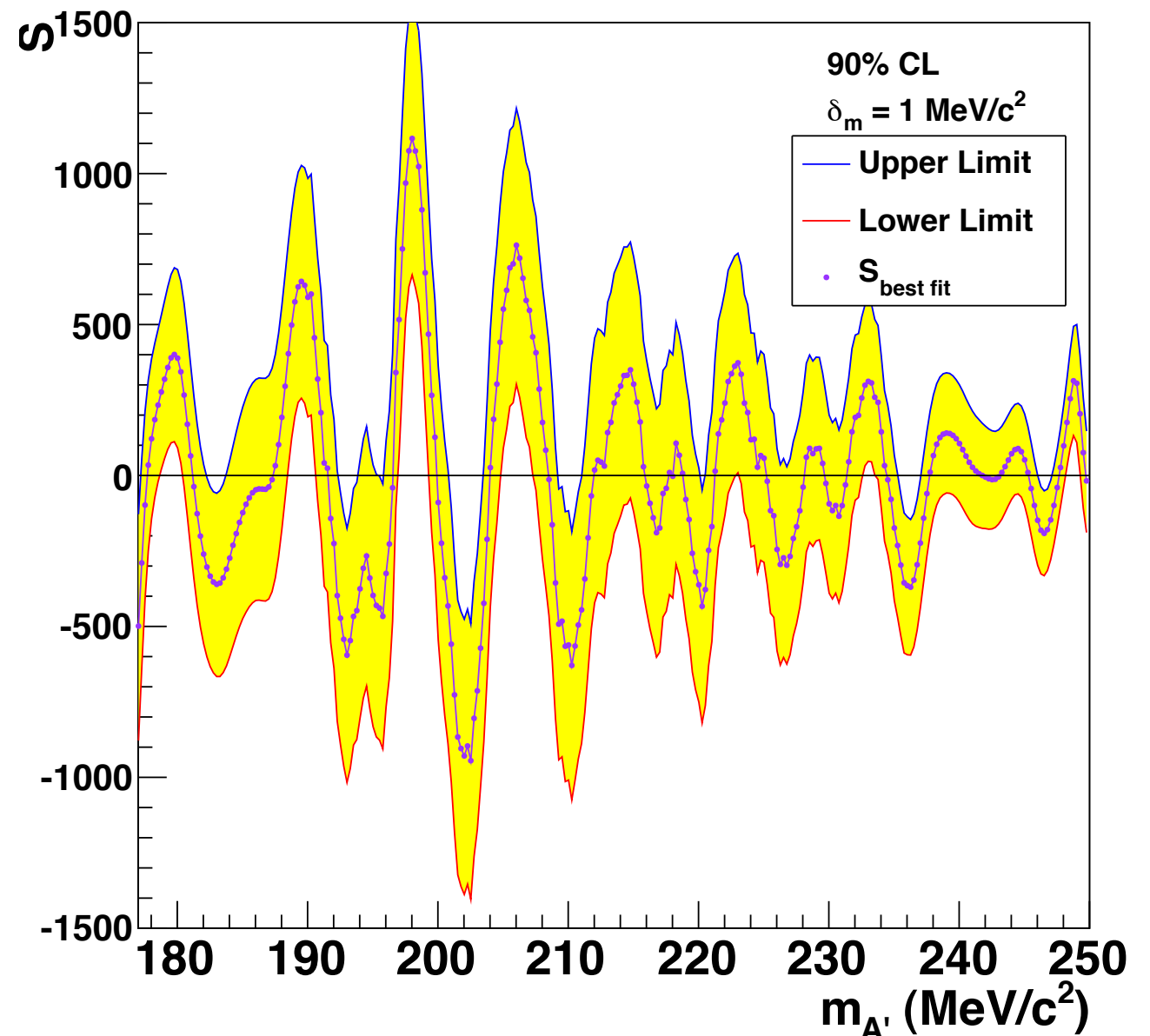
Characteristics of limit plots

Particular choice of polynomial basis and order captures the shape of the actual data or MC to a greater or lesser extent

Pseudodata set generated from 12th order polyn. pdf



Pseudodata from MC, global 12th order polyn. fit



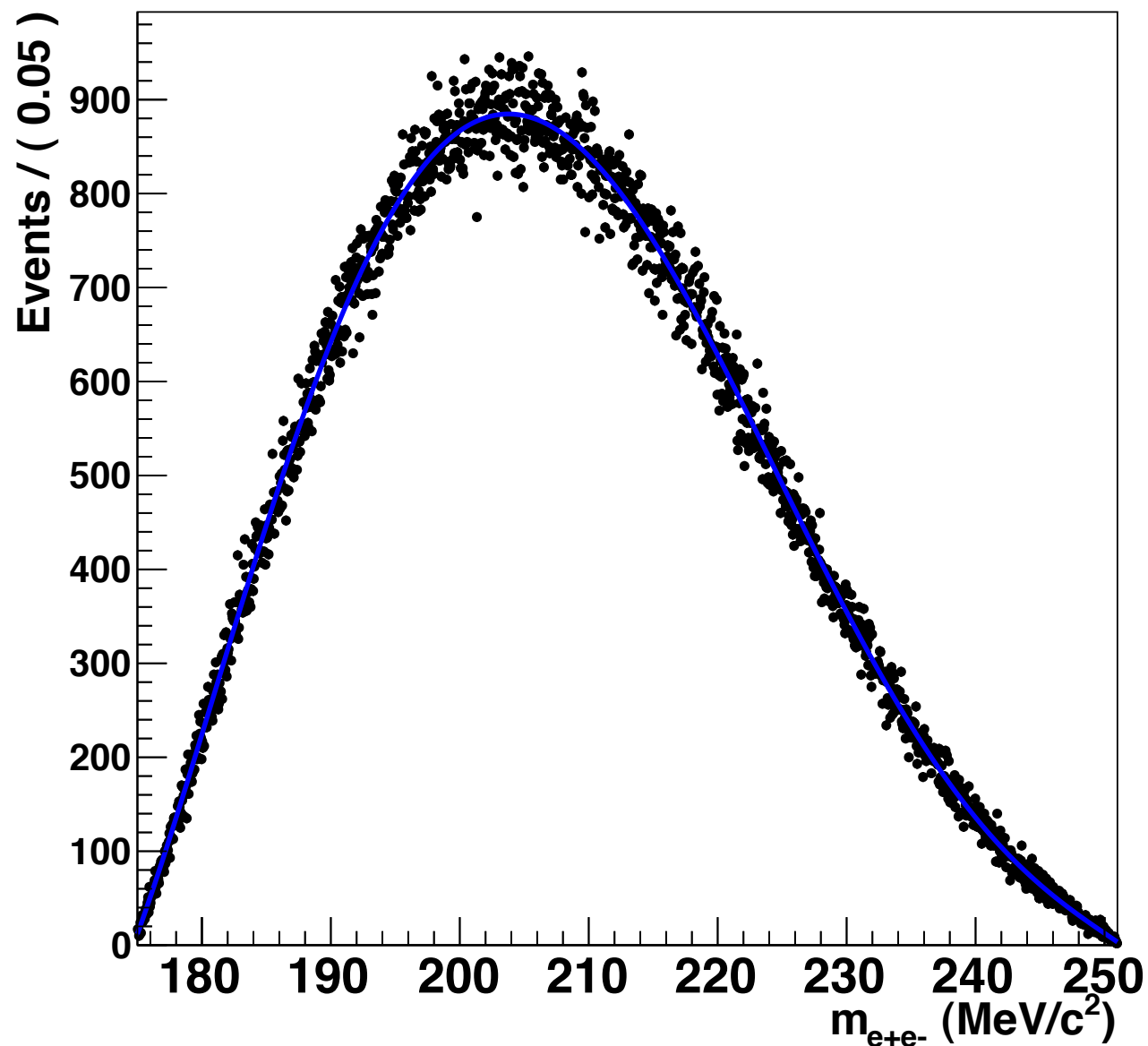
12th order global polyn. fit, 7th order polyn. background in window



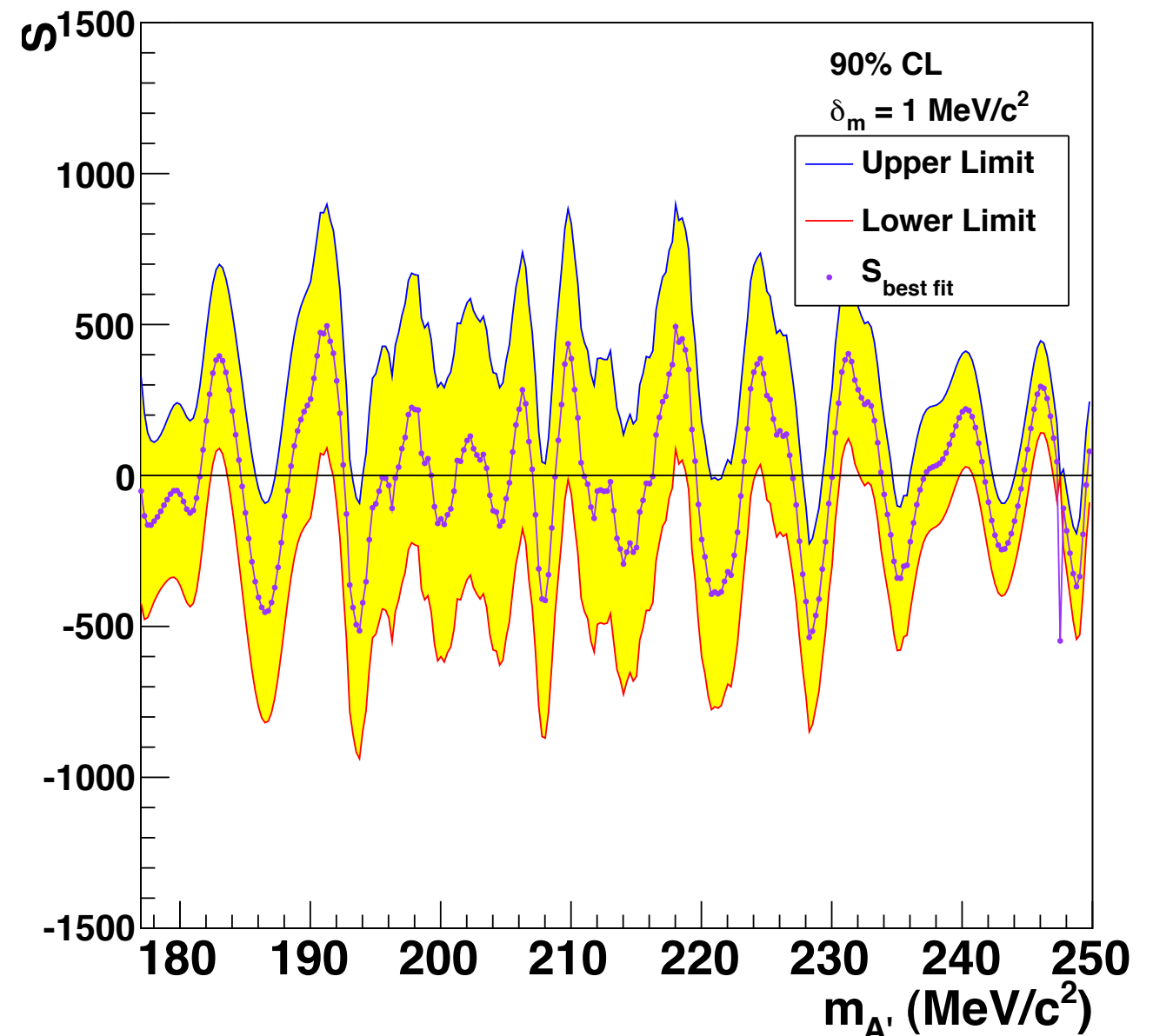
Characteristics of limit plots

Particular choice of polynomial basis and order captures the shape of the actual data or MC to a greater or lesser extent

Pseudodata set generated from 6th order polyn. pdf



Pseudodata from MC, global 6th order polyn. fit

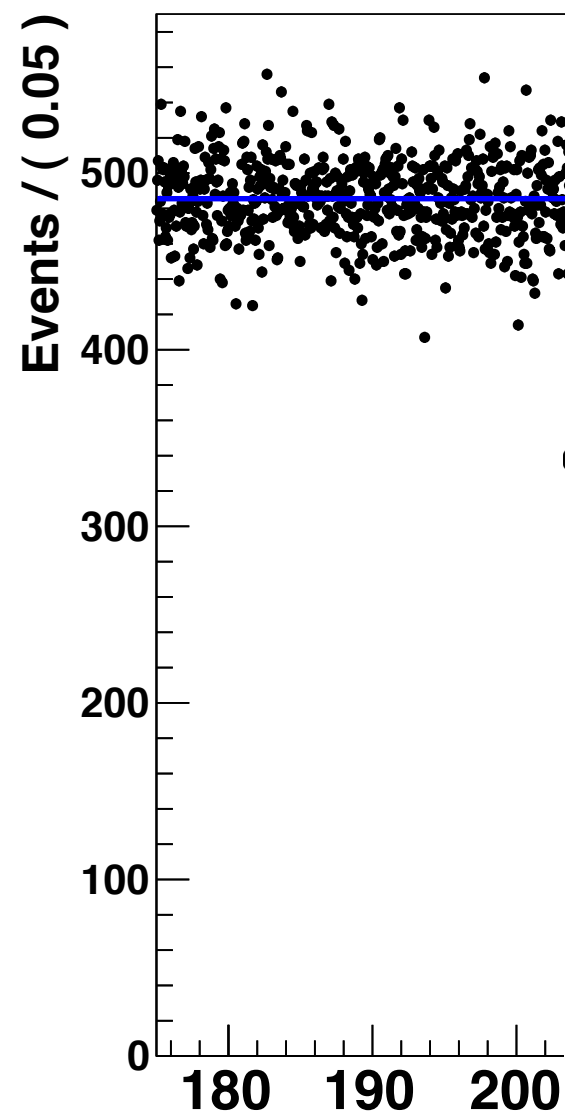


6th order global polyn. fit, 7th order polyn. background in window

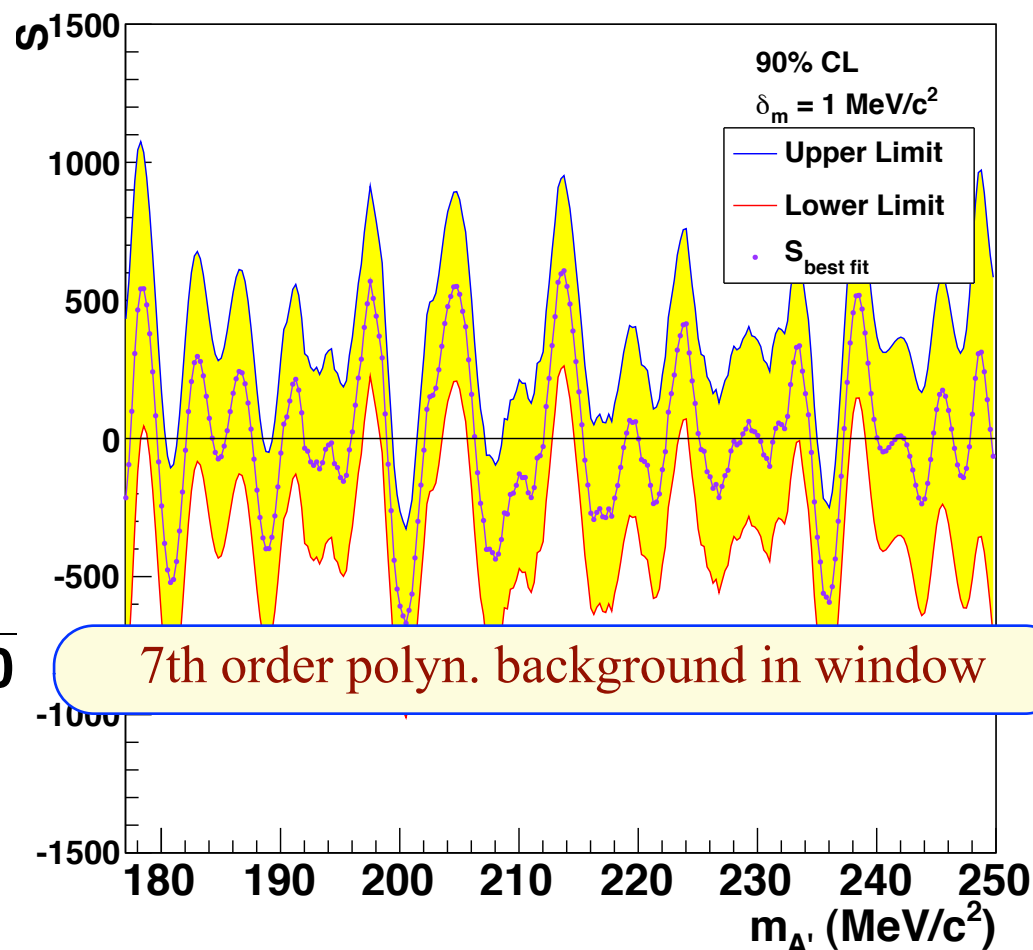


Characteristics of limit plots

Linear model pdf

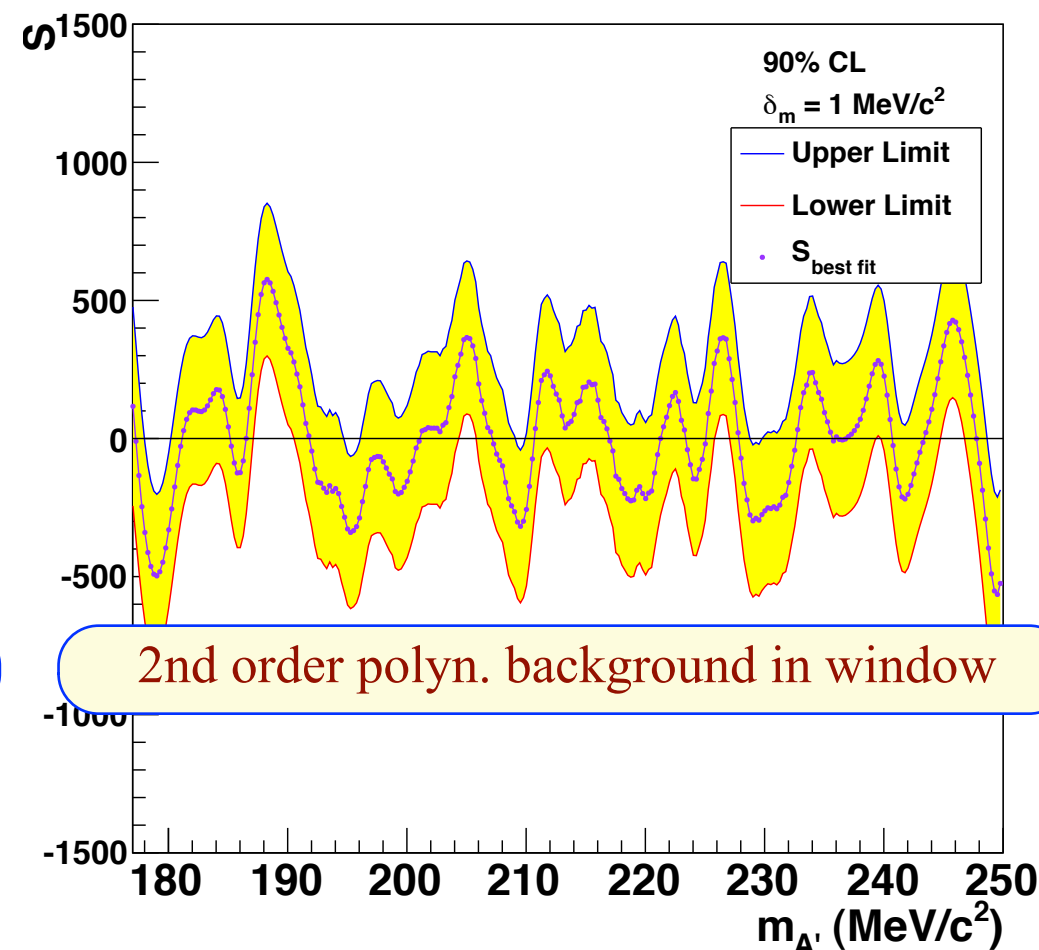


Pseudodata from linear model pdf



7th order polyn. background in window

Pseudodata from linear model pdf

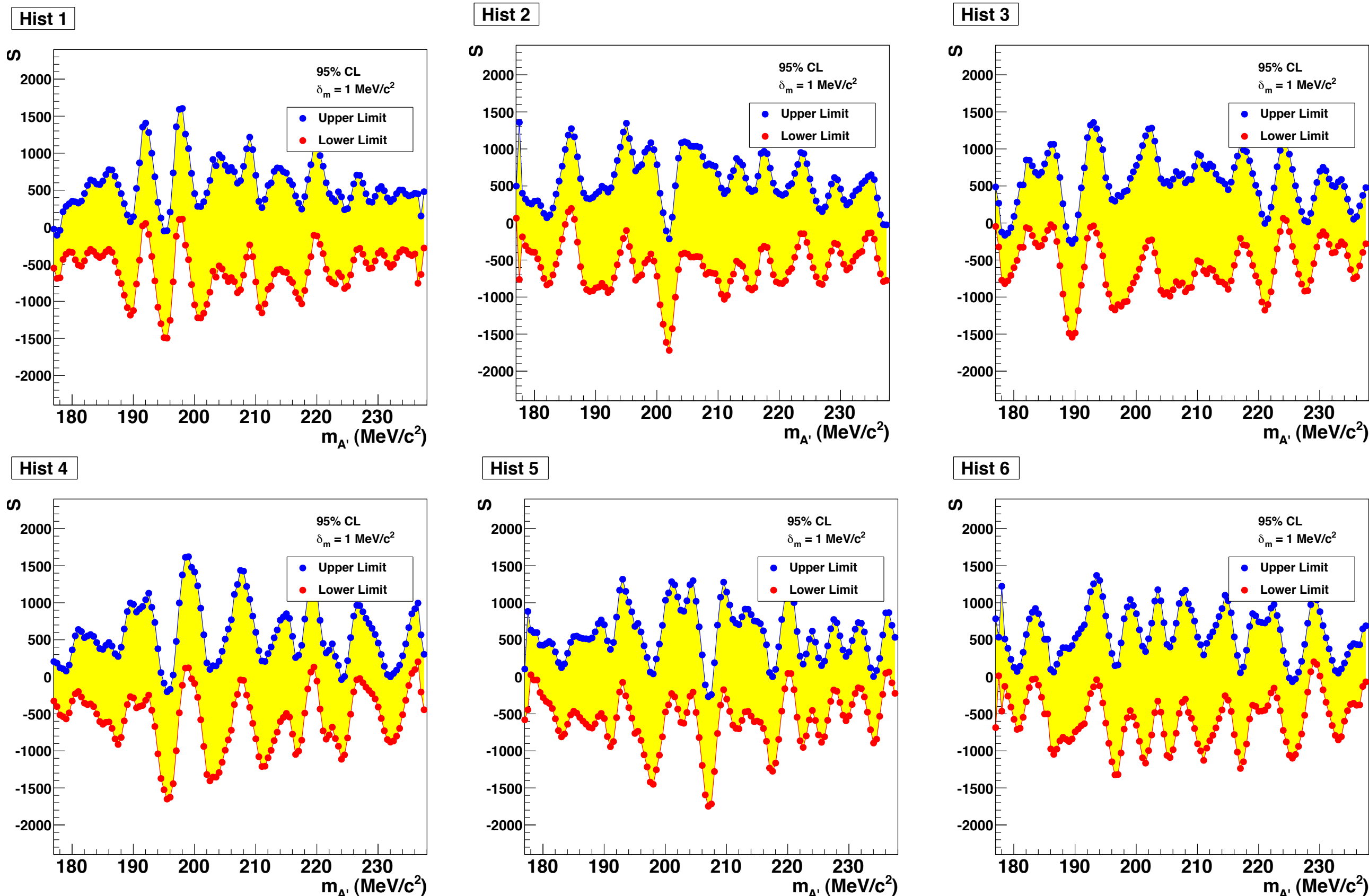


2nd order polyn. background in window

Other background parameterizations yield wiggles, too, though location and size varies; a systematic effect subdominant to statistical error

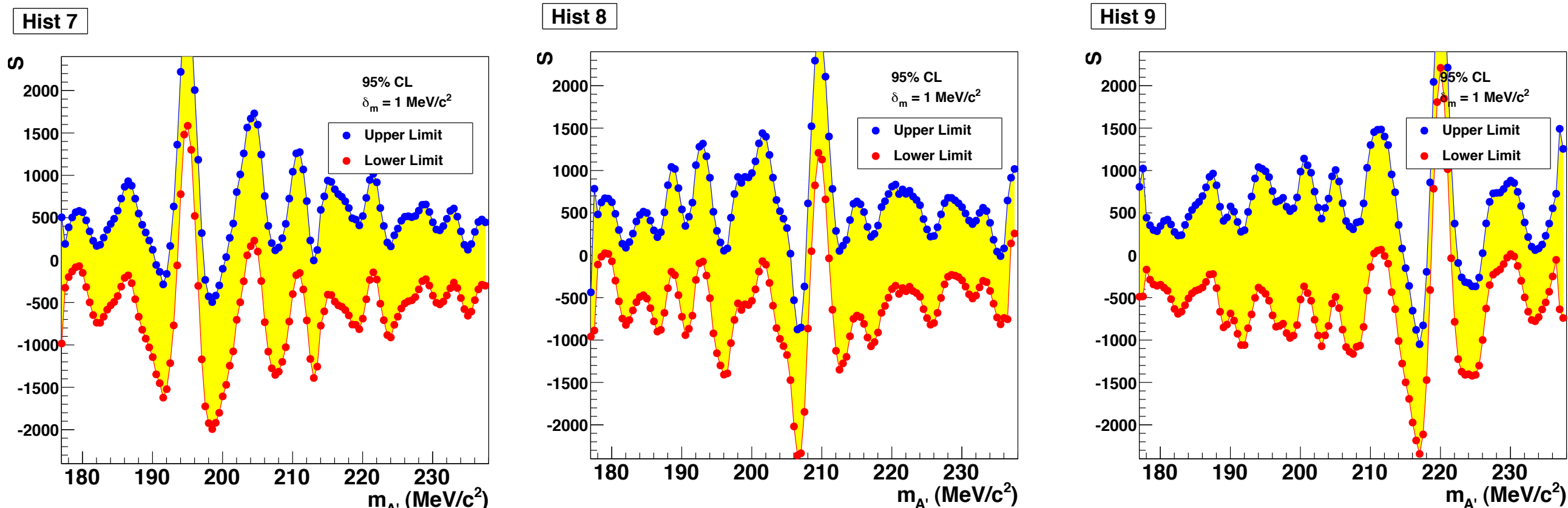


Blind analysis: Nine histograms; eight pseudodata, one real



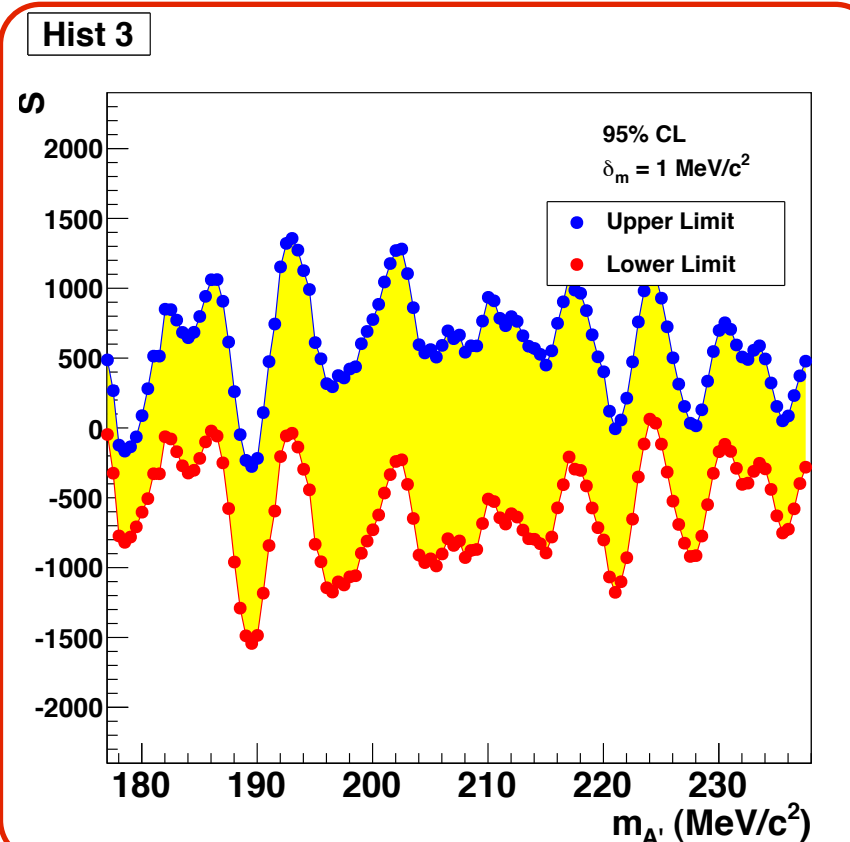


Blind analysis: Nine histograms; eight pseudodata, one real



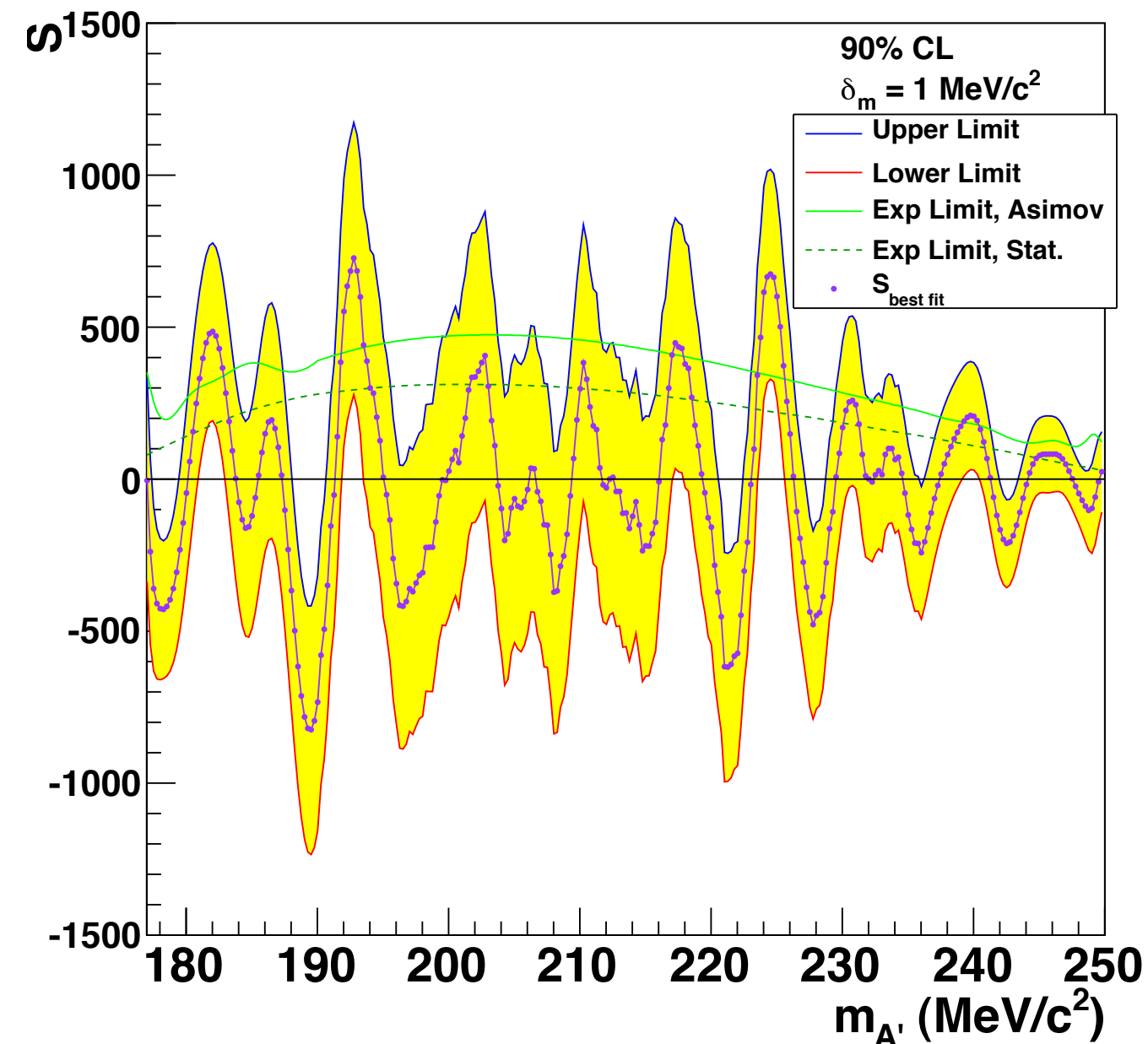
Unblinding: Histogram #3

Additionally, 40 similar histograms were prepared, with peaks of various significances inserted; method found all inserted excesses of nominal significance > 2 sigma

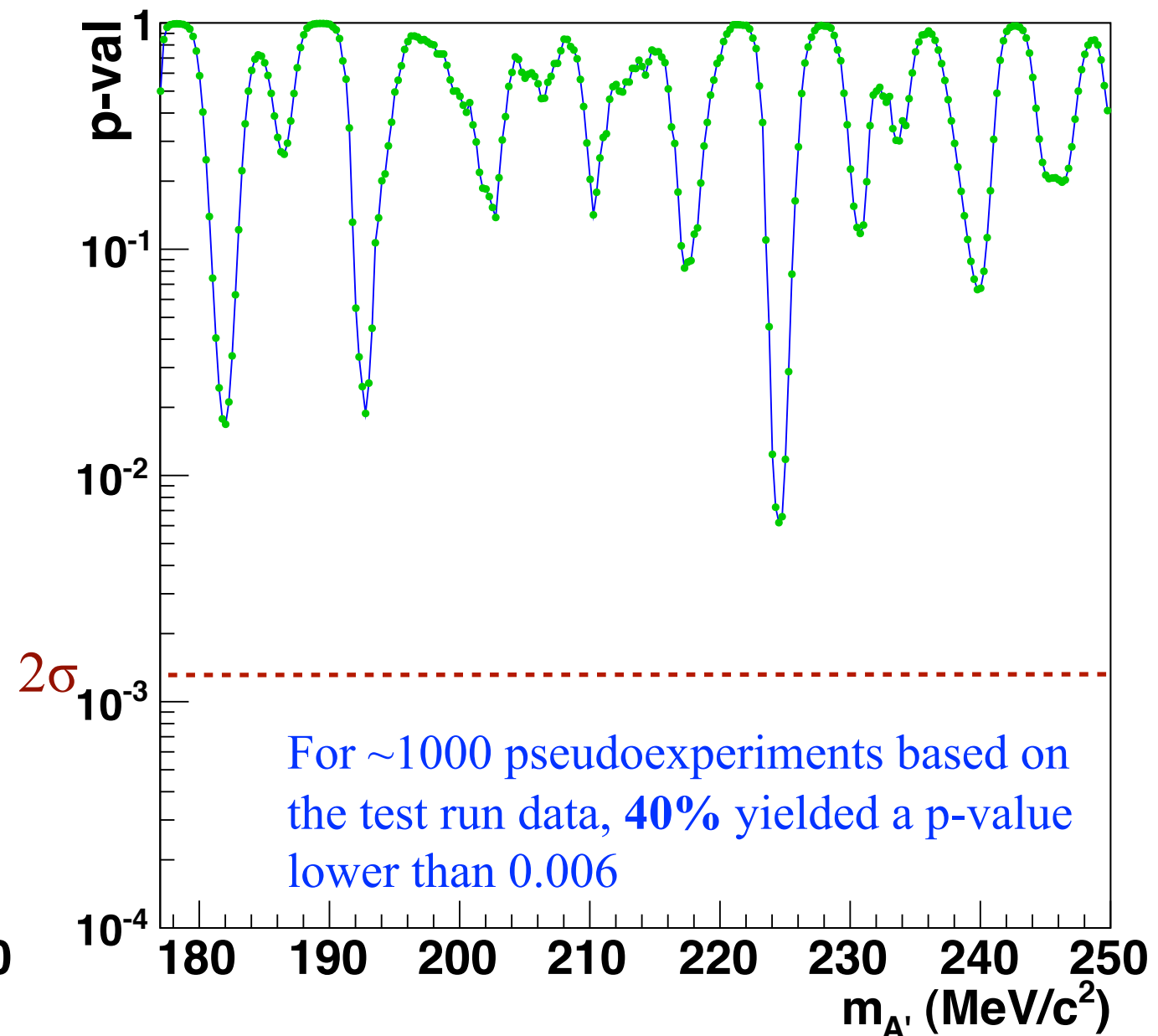


Results from scan of test run data: S, P-values

APEX Test Run Data, Two-Sided Central Limit



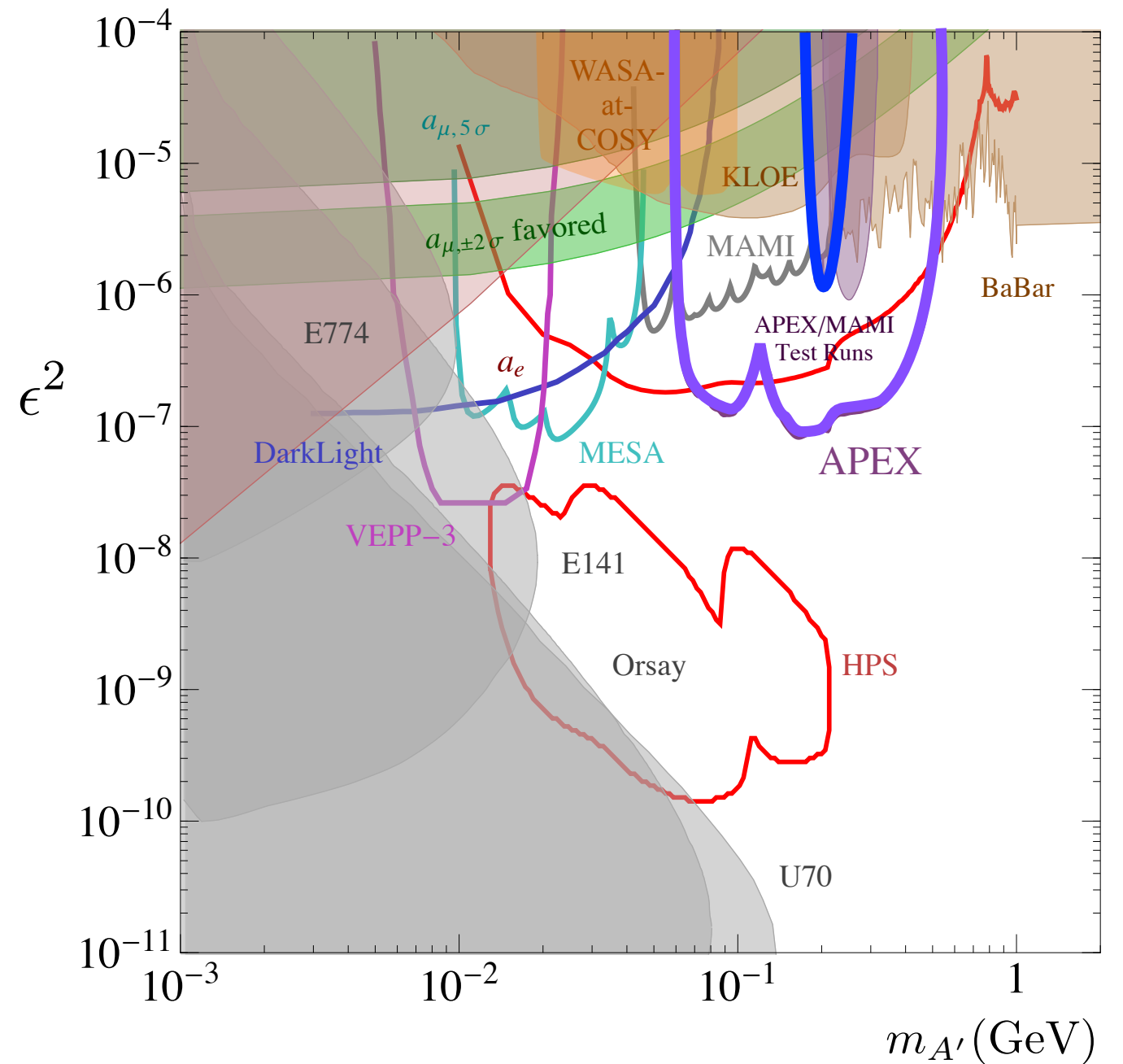
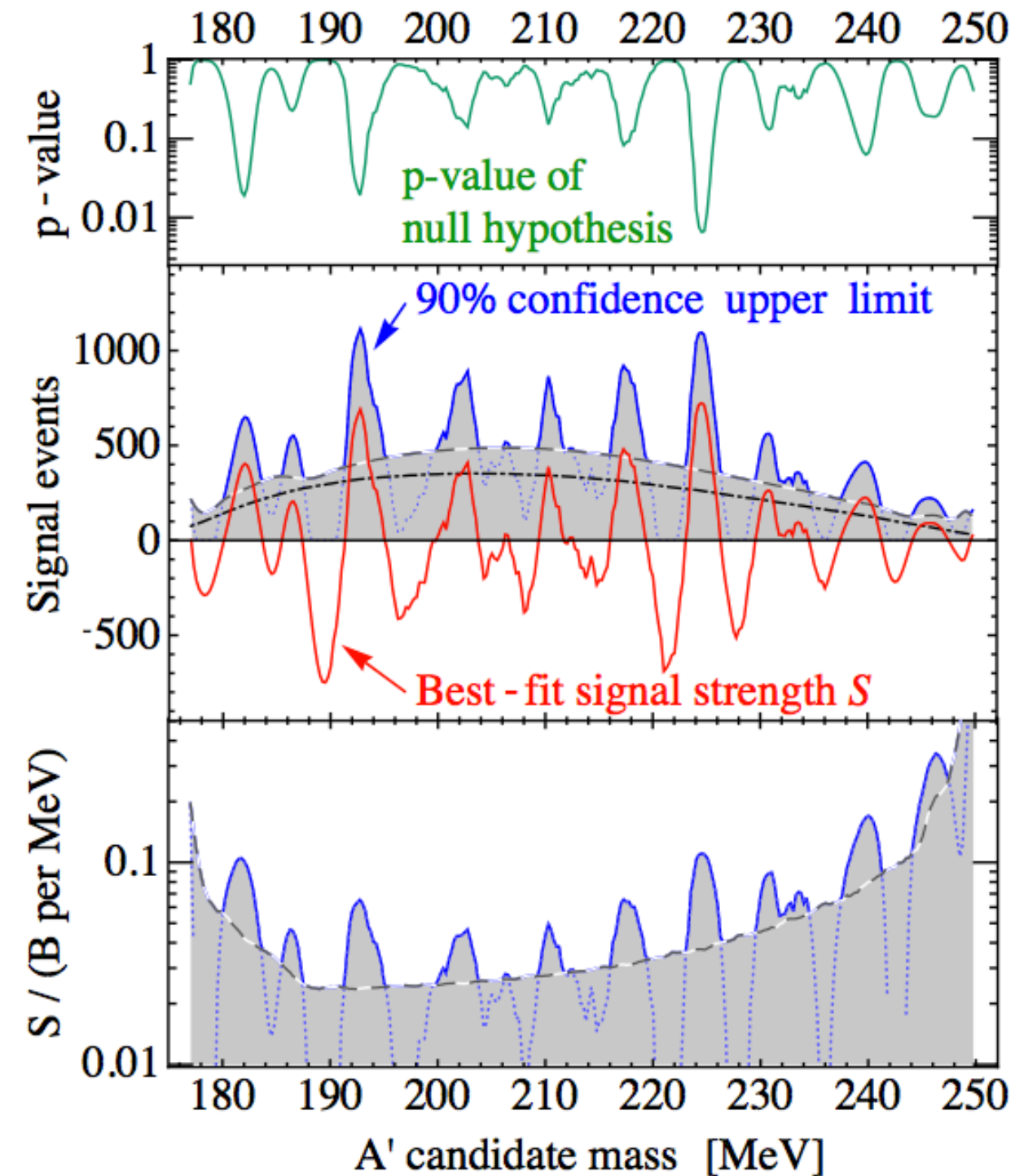
APEX Test Run Data, Raw Null P-values



Two-sided central limit; plot in paper is constrained to physical values of S



Upper limit on coupling





Plans for full run

Look closer at unbinned analysis

- Should be possible to address computational issues

Compare multiple methods for binned analysis

Investigate mass-dependent window sizes

Full MC needed to complement resonance search

End Part 1



In progress: Update of reach estimate for full run

Proposal-style reach calculation, with the inclusion of new LeRose acceptance

- Monte Carlo events are simulated for the Bethe-Heitler, radiative tridents, and the continuum trident background including the full interference effects between the diagrams. The latter background is computationally intensive, and only a small statistics sample is generated, sufficient to obtain the cross-section from MadEvent.
- The cross-section ratio of the full continuum background (with interference effects) to the sum of the Bethe-Heitler and radiative tridents is calculated, and represents a multiplicative factor by which the latter must be multiplied to get the background cross-section.
- The rates of all reactions impinging the spectrometer acceptance were calculated by integrating over a chosen target profile, which usually extended from 4.5 to 5.5 degrees. For Bethe-Heitler, radiative tridents, and the continuum trident background, the calculation of the rate was performed as a function of invariant mass.
- Using the expressions in Appendix B, we calculated the mass resolution δ_m . We then tiled the acceptance region with bins of size $2.5 \times \delta_m$ in invariant mass.
- As a function of α'/α , the total number of signal (S) and background (B) events was calculated with the help of the formulas in §3 for each bin.
- We then set $S/\sqrt{B} = 2$, and solved for α'/α .

APEX PAC37

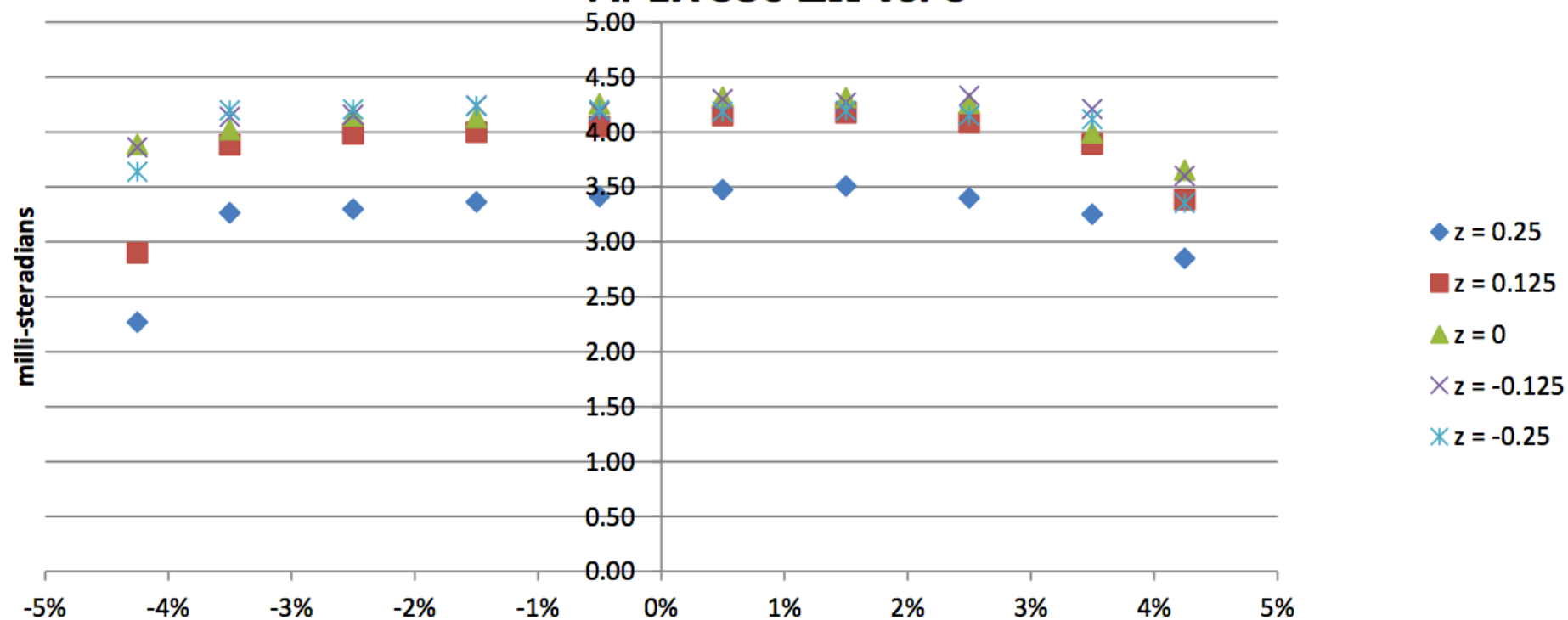
Changes from this prescription:

- Monte Carlo for radiative trident and Bethe-Heitler backgrounds, generated at four beam energies (not necessarily those being used for full run), including correction factors for interference effects
- Ten foils/angular settings between 4.5 and 5.5 degrees → Utilize MC multiple times
- Each “run” gives an unbinned dataset of invariant masses; tile the spectrum in bins of $2.5 \times (\text{mass resolution})$ and calculate 2σ reach in ϵ^2
- New here is the imposition of the LeRose S80 acceptance numbers
- The following is **pre-preliminary**



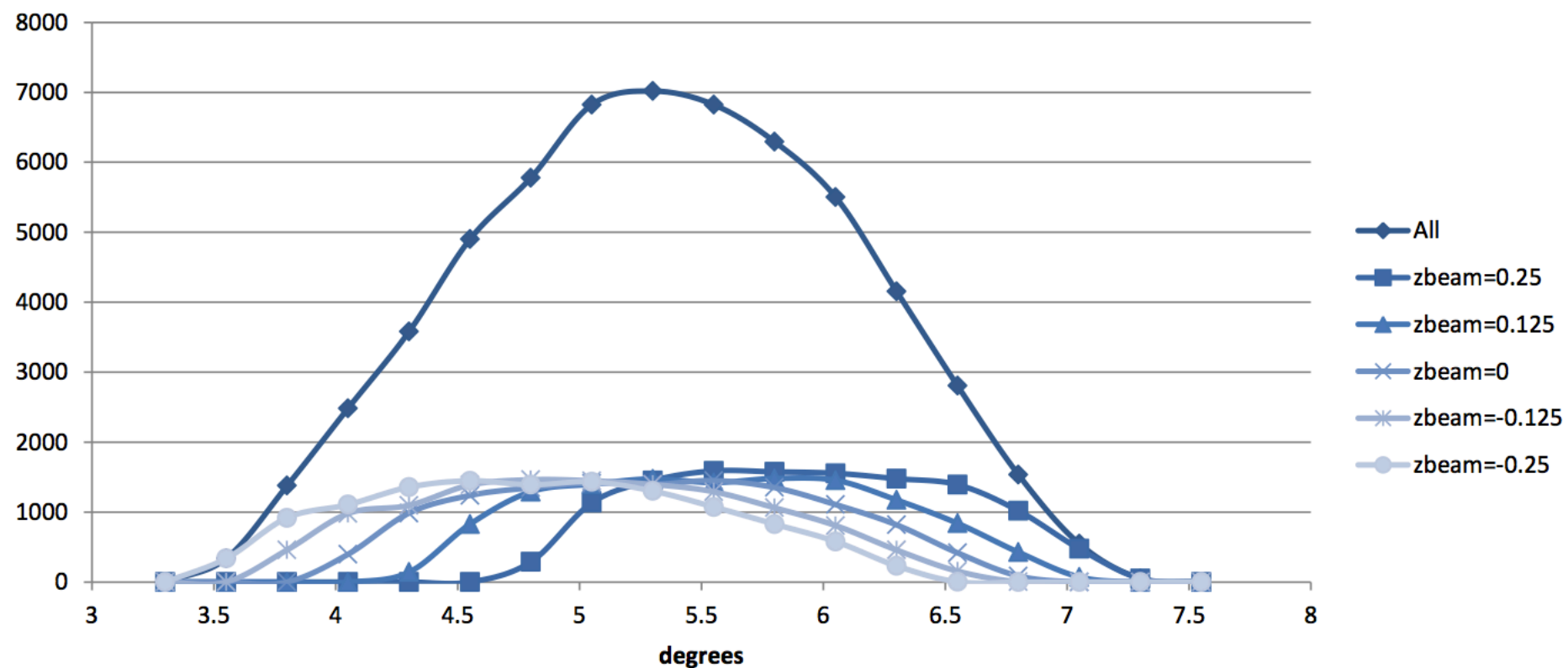
In progress: Update of reach estimate for full run

APEX S80 $\Delta\Omega$ vs. δ



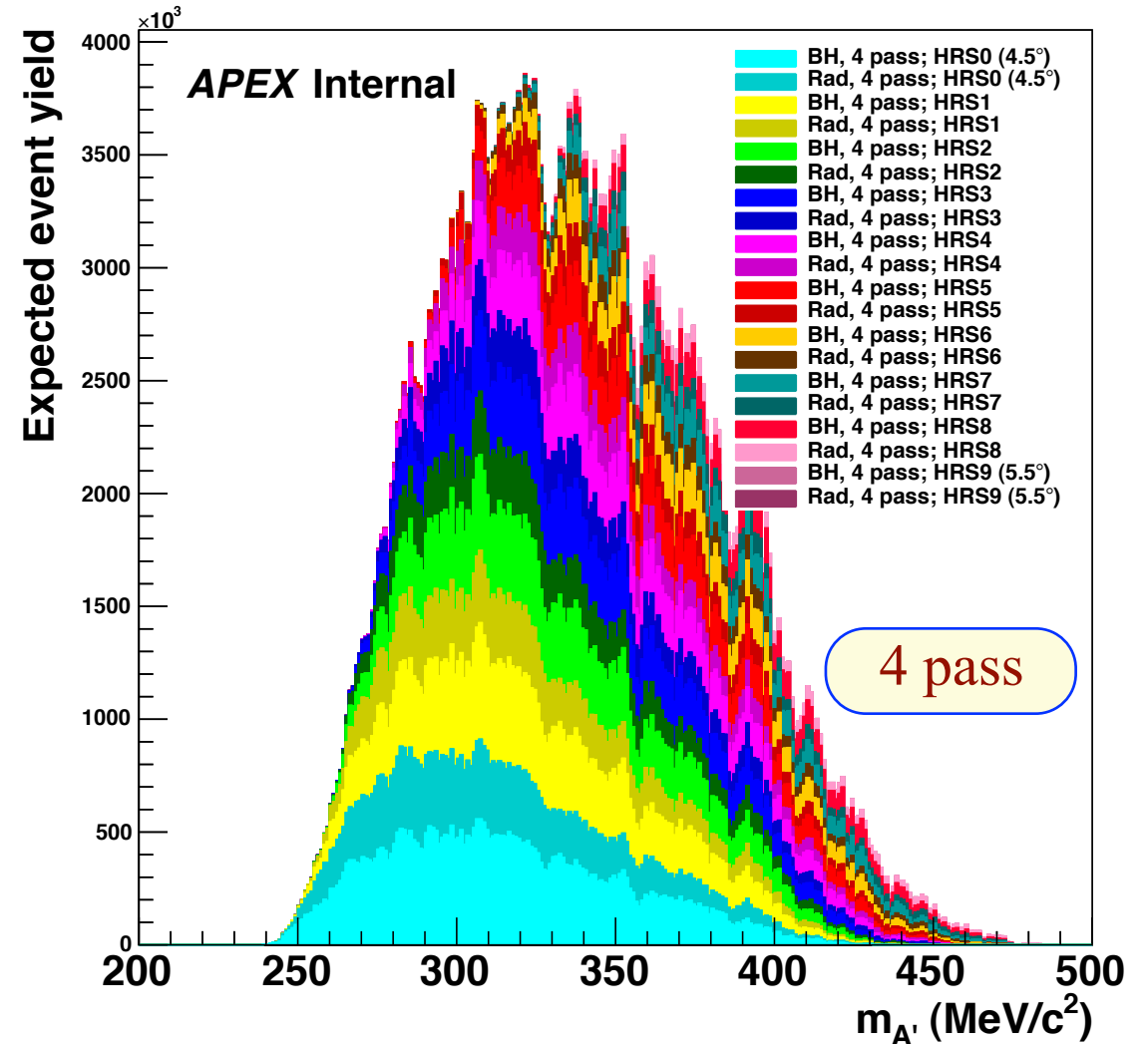
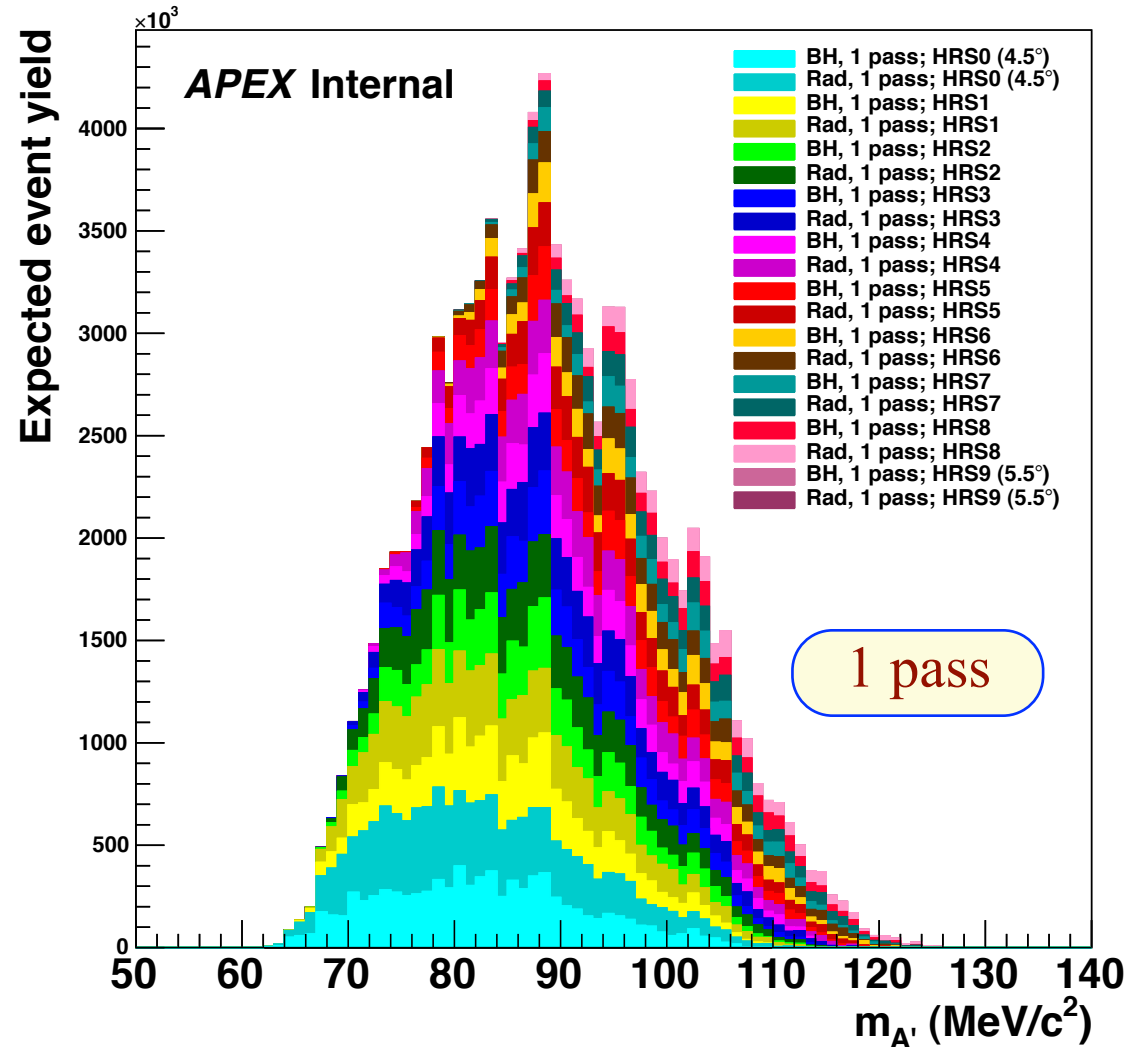
New acceptance for
septum magnets calculated
by John LeRose

APEX S80 Scattering angle coverage





In progress: Update of reach estimate for full run

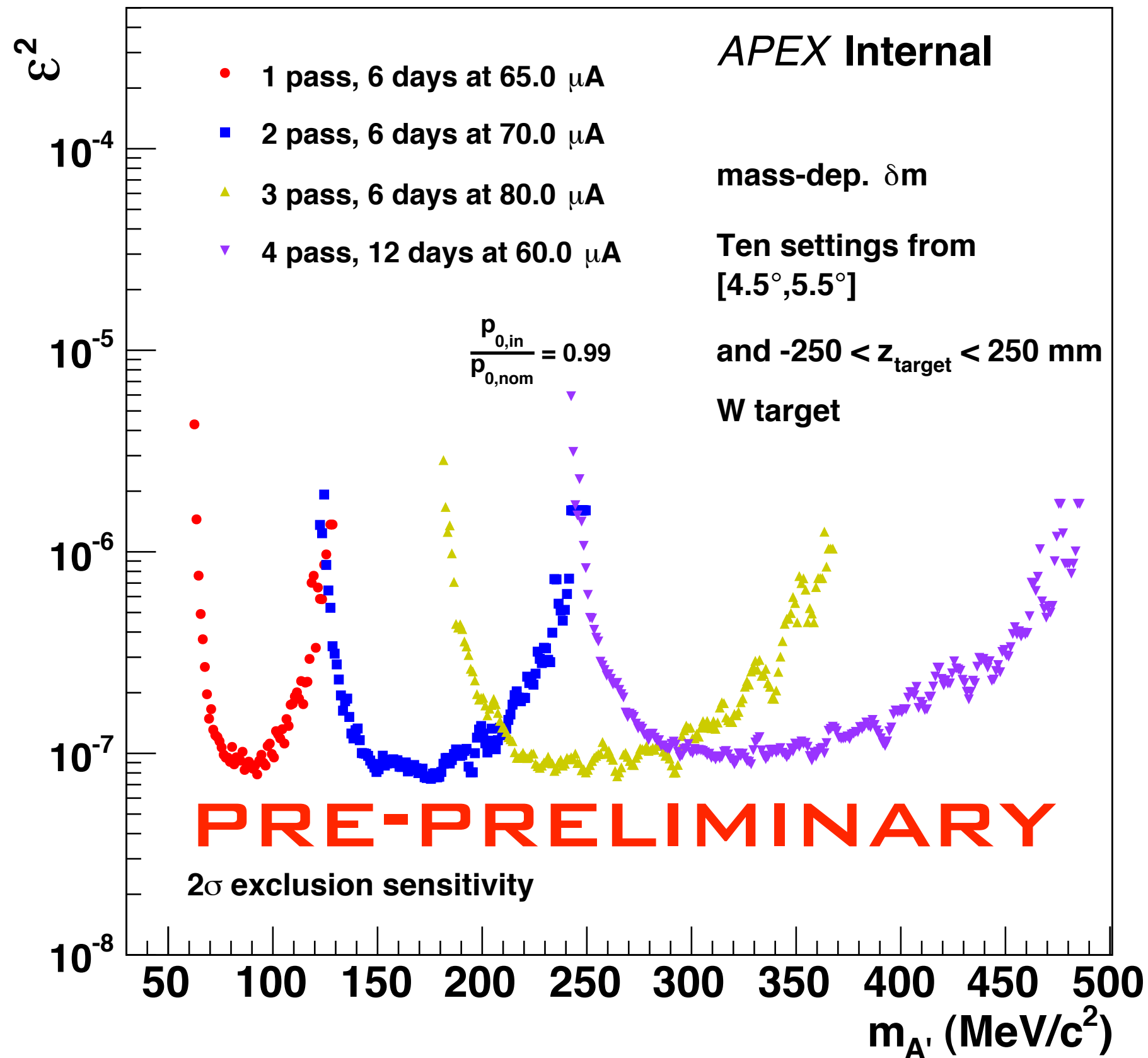


Monte Carlo for
radiative trident
and Bethe-Heitler
backgrounds

E_{beam}	$p_0 = 0.99 \cdot E_{\text{beam}}/2$	Time [days]	Current [μA]	T [rad. lengths]
1.056	0.523	6	65	0.007
2.056	1.018	6	70	0.04
3.056	1.513	6	80	0.08
4.056	2.008	12	60	0.08



In progress: Update of reach estimate for full run



Proposal-style
reach calculation
with LeRose S80
acceptance
numbers

**STILL VALIDATING;
WILL CHANGE;
MACHINERY IN
PLACE**

End

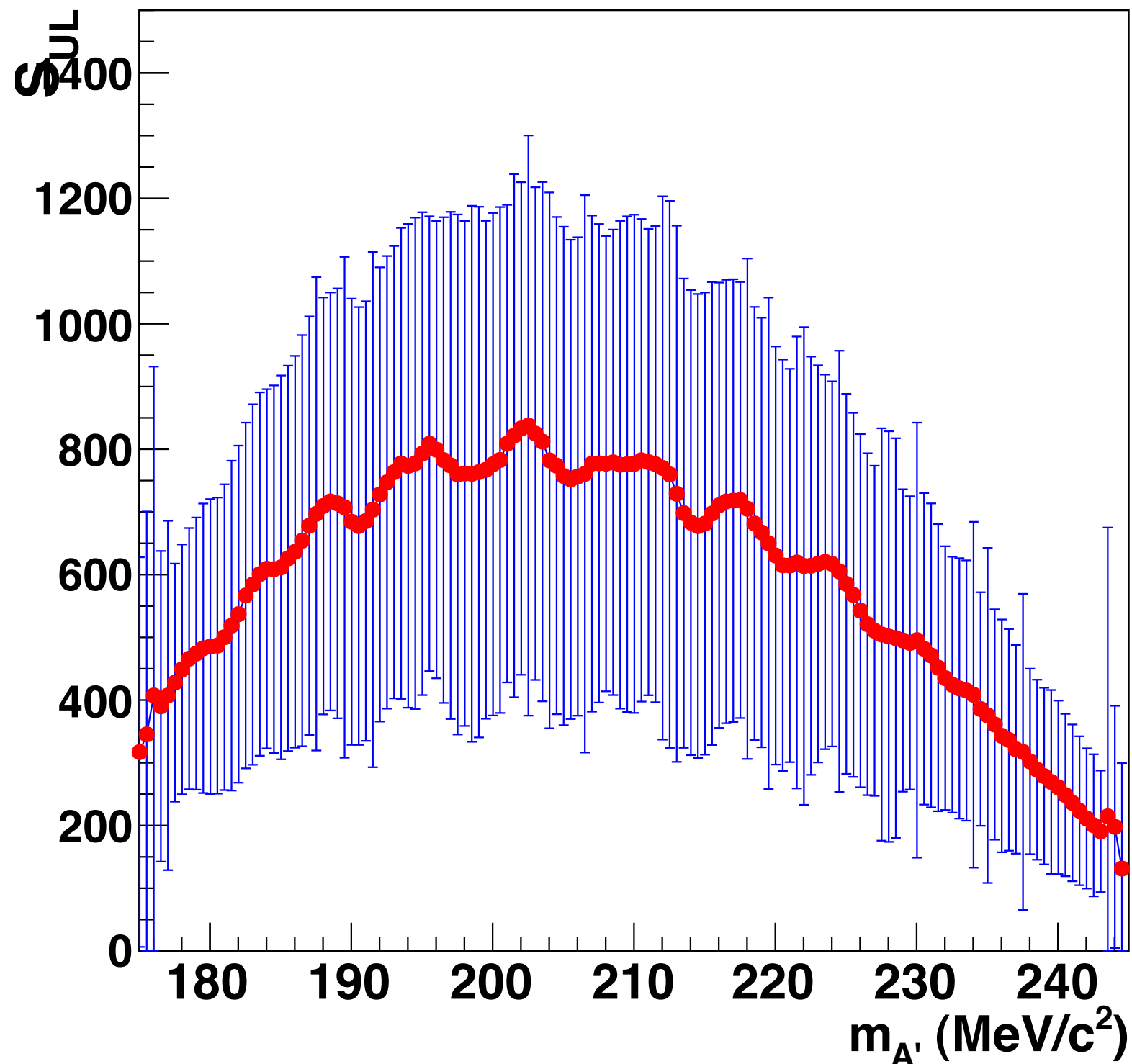


Backups



Asimov data is the limit of an ensemble

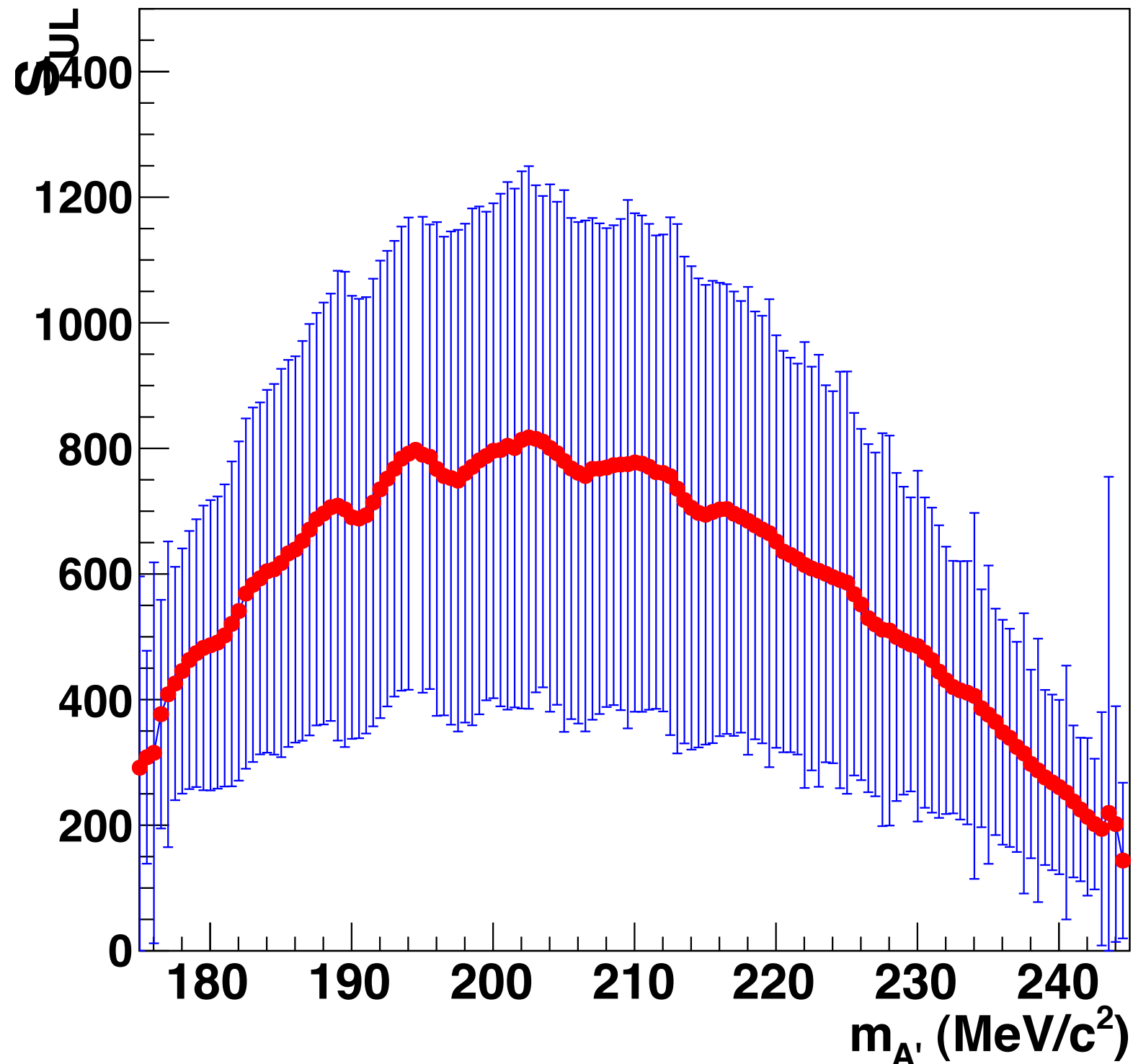
S_{UL} , 2500 datasets





Asimov data is the limit of an ensemble

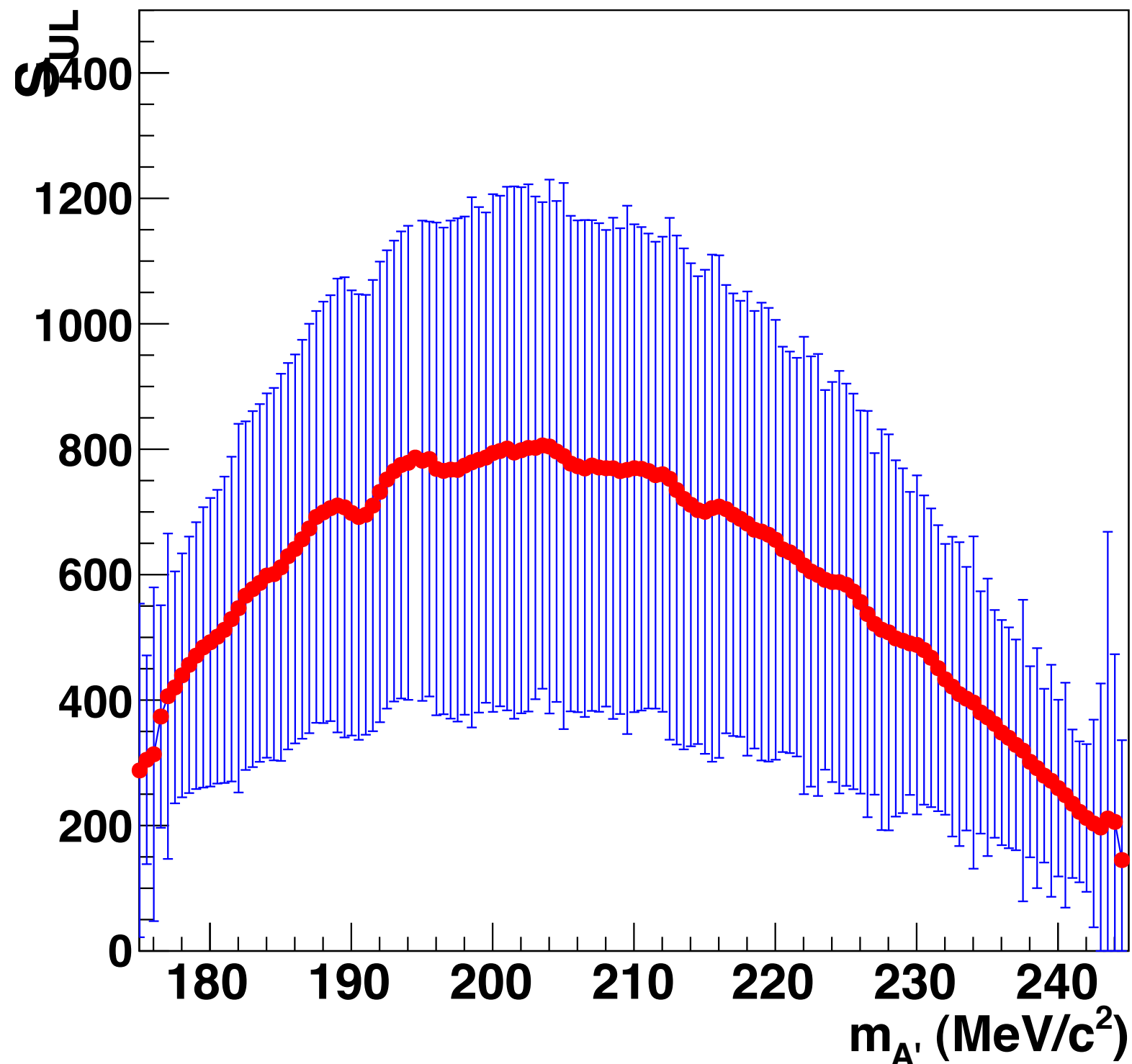
S_{UL} , 4087 datasets





Asimov data is the limit of an ensemble

S_{UL} , 6303 datasets





Asimov data is the limit of an ensemble

S_{UL} , 8000 datasets

