

Nucleon Form Factors

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Proton and Neutron FFs in
space-like domain
(focus on new development)

EINN-2007

Milos

September 12-15, 2007

Outline

- Introduction – EM FFs,
 - » New parameterizations
- New experimental results –
 - » low Q^2
 - » high Q^2
- Theory links observations and ideas –
 - » Charge density
 - » GPDs fit from FFs
- Forthcoming experiments – GEP-III, ..., 12 GeV
- Summary

Reviews and Analysis

Reviews

Ch. Hyde-Wright and C. de Jager, Ann. Rev. Nucl. Part.Sci. 54, 217 (2004)
H. Gao, Int. J. Mod. Phys A20, 1595 (2005)
Ch. Perdrisat, V. Punjabi, and M. Vanderhaeghen, [hep-ph/06012014](#)
J. Arrington, C.D. Roberts, and J.M. Zanotti, [nucl-th/0611050](#)

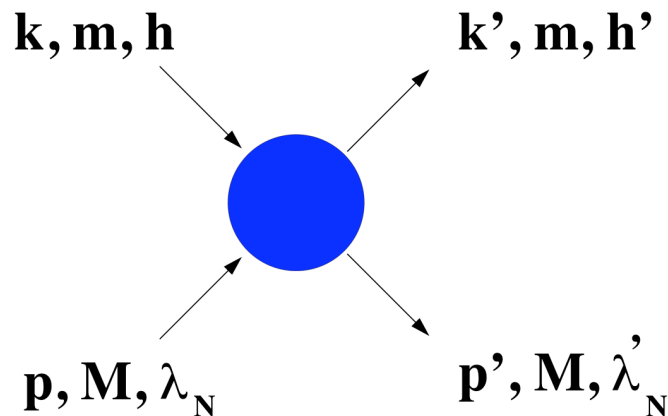
Parameterizations

P. Bosted, PRC 64, 409 (1995) - Fit of pre-JLab data
E. Brush et al. PRC 65, 051001(R) 2002 - Fit with JLab GEP/GMP
J.J. Kelly, PRC 66, 065203 (2002) - Breit frame densities
[BABB](#), [arXiv:hep-ex 0708.1946](#) - Fit with local duality constraints
J. Arrington, W.Melnitchouk, J.A.Tjon, [nucl-ex/0707.1861](#) - 2-gamma

GPD

M.Diehl et al., Eur.Phys.J. C39 (2005) 1-39, [GPDs from FFs data](#)
M.Guidal et al., PRD 72, 054013 (2005) , [FFs from GPDs](#)

Lepton-Nucleon scattering



$$l(k, h) + N(p, \lambda_N) \rightarrow l(k', h') + N(p', \lambda'_N)$$

h, h', λ_N , and λ'_N are helicities

$$P = \frac{p+p'}{2}, K = \frac{k+k'}{2}, q = k - k' = p' - p$$

$$s = (p + k)^2, t = q^2 = -Q^2, u = (p - k')^2$$

$$T_{\lambda'_N, \lambda_N}^{h', h} \equiv \langle k', h'; p', \lambda'_N | T | k, h; p, \lambda_N \rangle$$

Total 16 amplitudes.

Parity invariance \rightarrow number of independent helicity amplitudes from 16 to 8.

Time reversal invariance \rightarrow to 6.

When neglect the lepton mass \rightarrow to 3.

Three complex amplitudes

$$T_{+,+}^{+,+}; \quad T_{-,-}^{+,+}; \quad T_{-,+}^{+,+} = T_{+,-}^{+,+}$$

which are functions of $(s - u)$ and t .

Electro-Magnetic Form Factors

One-photon approximation, $\alpha_{em} = 1/137$, hadron current

$$\mathcal{J}_{hadronic}^\mu = ie\bar{N}(p') \left[\gamma^\mu F_1(Q^2) + \frac{i\sigma^{\mu\nu}q_\nu}{2M} F_2(Q^2) \right] N(p) \quad \text{Rosenbluth (1950)}$$

Full expression for \mathcal{M} has three complex functions, F_1, F_2, F_3 Guichon & Vanderhaeghen

$$\mathcal{M} = \frac{4\pi\alpha}{Q^2} \bar{u}' \gamma_\mu u \cdot \bar{N}' \left(\tilde{F}_1 \gamma^\mu - \tilde{F}_2 [\gamma^\mu, \gamma^\nu] \frac{q_\nu}{4M} + \tilde{F}_3 K_\nu \gamma^\nu \frac{P^\mu}{M^2} \right) N \quad \text{Afanasev et al.}$$

$$\tilde{G}_M = \tilde{F}_1 + \tilde{F}_2 \quad \tilde{G}_E = \tilde{F}_1 - \tau \tilde{F}_2$$

\tilde{F}_i are functions of $(s - u)$ and t

Blunden et al.

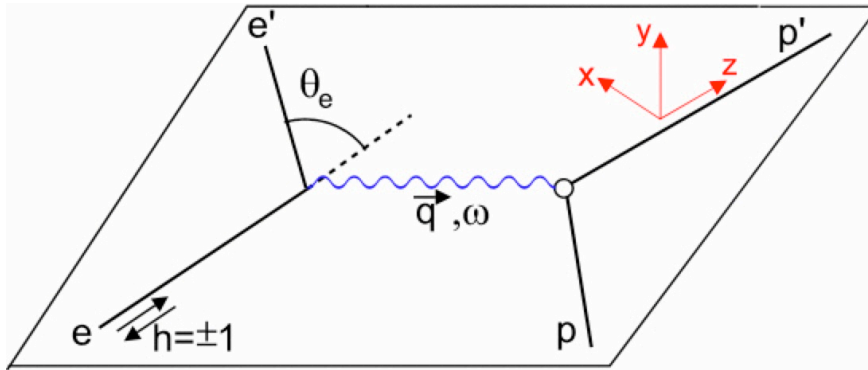
$$d\sigma = d\sigma_{NS} \left\{ \varepsilon (\tilde{G}_E + \frac{s-u}{4M^2} \tilde{F}_3)^2 + \tau (\tilde{G}_M + \varepsilon \frac{s-u}{4M^2} \tilde{F}_3)^2 \right\}$$

old $G_{E,M}$ are real functions of $t=-Q^2$

$$\sigma_R = \varepsilon G_E^2 + \tau G_M^2 + 2\tau G_M \operatorname{Re} \left(\delta \tilde{G}_M + \varepsilon \frac{s-u}{M^2} \tilde{F}_3 \right) + 2\varepsilon G_E \operatorname{Re} \left(\delta \tilde{G}_E + \frac{s-u}{M^2} \tilde{F}_3 \right)$$

Extra terms contribute less than few % to σ_R

Double Polarized Observables



$$N(\vec{e}, e' \vec{N})$$

Akhiezer et al., (1958)

Arnold et al., (1981)

Guichon &

Vanderhaeghen (2003)

Two-photon correction, $\delta \sim 0.02$
at typical values $\varepsilon = 0.3 \text{ --} 0.8$

$$P_x = -2\sqrt{\tau(\tau+1)}G_E^p G_M^p \tan \frac{\theta_{e'}}{2} / I_0$$

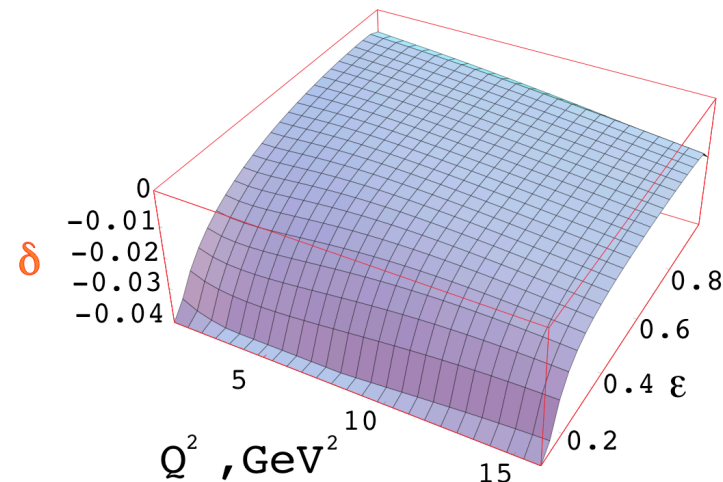
$$P_z = \frac{E_e + E_{e'}}{M_p} \sqrt{\tau(\tau+1)} (G_M^p)^2 \tan^2 \frac{\theta_{e'}}{2} / I_0$$

$$I_0 \propto \epsilon (G_E^p)^2 + \tau (G_M^p)^2$$

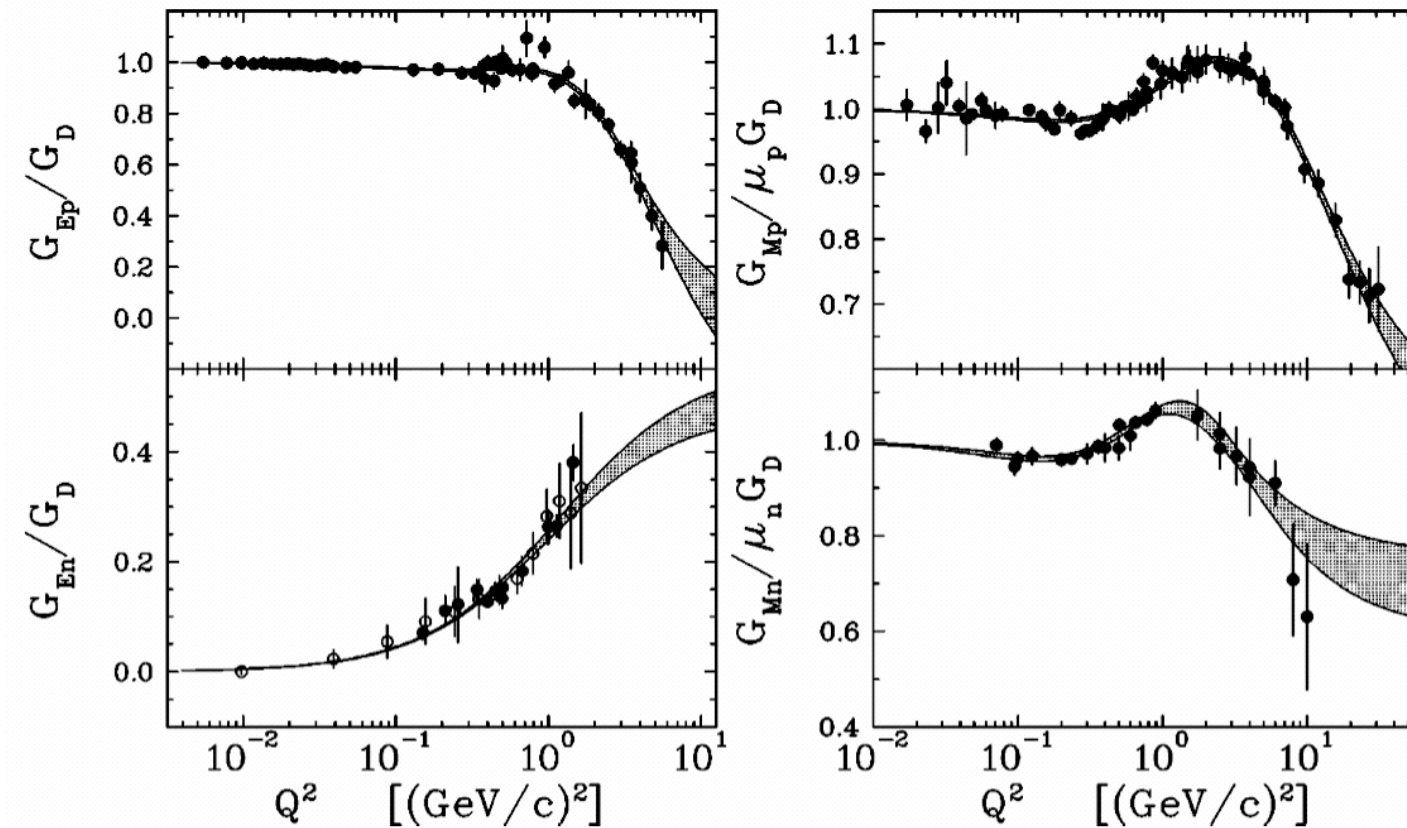
$$\left. \frac{G_E^p}{G_M^p} \right|_{1-\gamma} = -\frac{P_x}{P_z} \frac{E_e + E_{e'}}{2M_p} \tan(\theta_e/2)$$

$$\mu \left. \frac{G_E^p}{G_M^p} \right|_{1,2-\gamma} = \mu \left. \frac{G_E^p}{G_M^p} \right|_{1-\gamma} + \delta$$

Similar result for polarized target case



Kelly's Parameterization

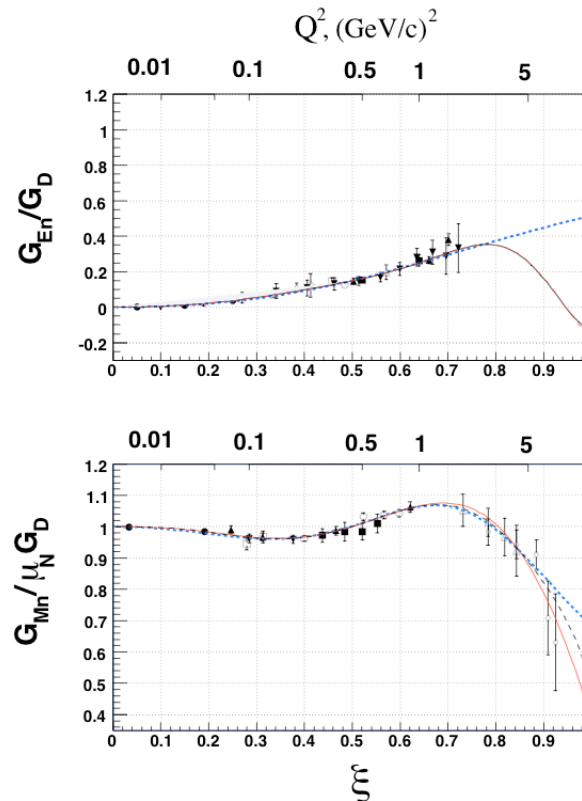
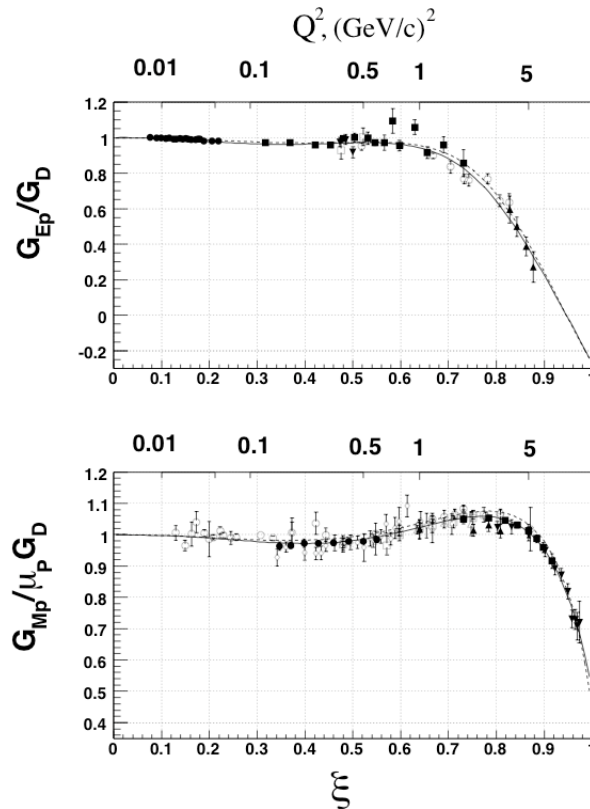


J. Kelly,
PRC 70,
068202
(2004)

$$G(Q^2) = \sum_{k=0}^n a_k \tau^k / (1 + \sum_{k=1}^{n+2} b_k \tau^k)$$

scaling constraint: $Q \rightarrow \infty, G \sim Q^{-4}$

Duality constrained parameterization



Bodek,
Avvakumov,
Bradford, Budd
arXiv:hep-ex
0708.1946

ξ is Nachtmann
scaling variable

Two constraints
QCD motivated

Kelly's

$(d/u) = 0, 0.2$

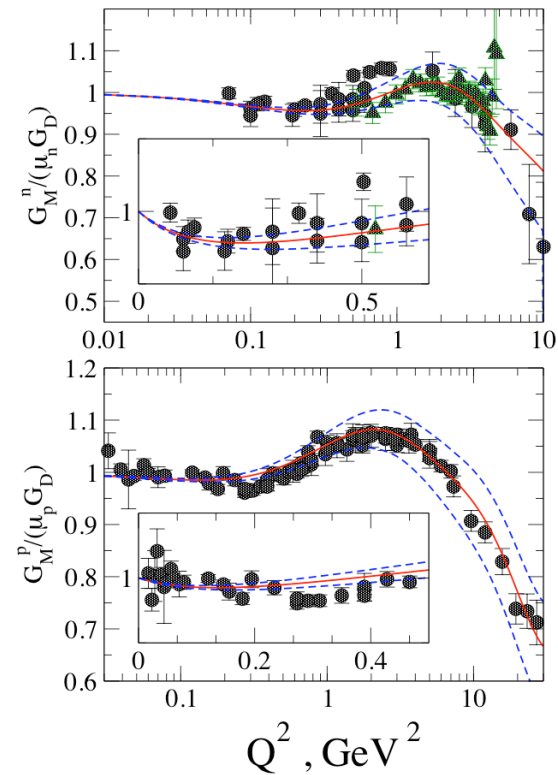
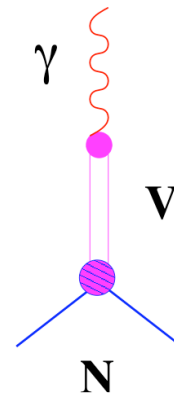
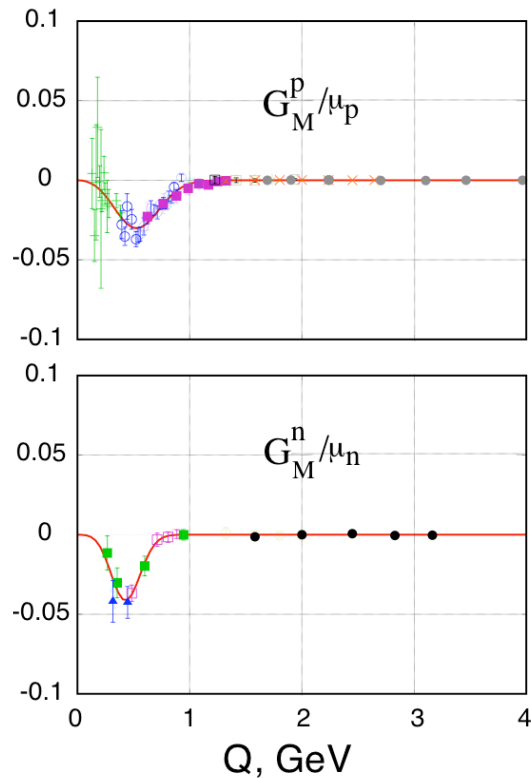
$$\xi^{p,n} = 2/(1 + \sqrt{1 + \tau_{p,n}^{-1}})$$

constrains: at $\xi \rightarrow 1$

$$G_{BABB} = A(\xi) \times G_{Kelly}(Q^2)$$

$$1) \frac{G_{Mn}^2}{G_{Mp}^2} = \frac{1+4(d/u)}{4+(d/u)} \quad 2) \frac{G_{En}^2}{G_{Mp}^2} = \frac{G_{Ep}^2}{G_{Mp}^2}$$

Form Factors at low Q^2



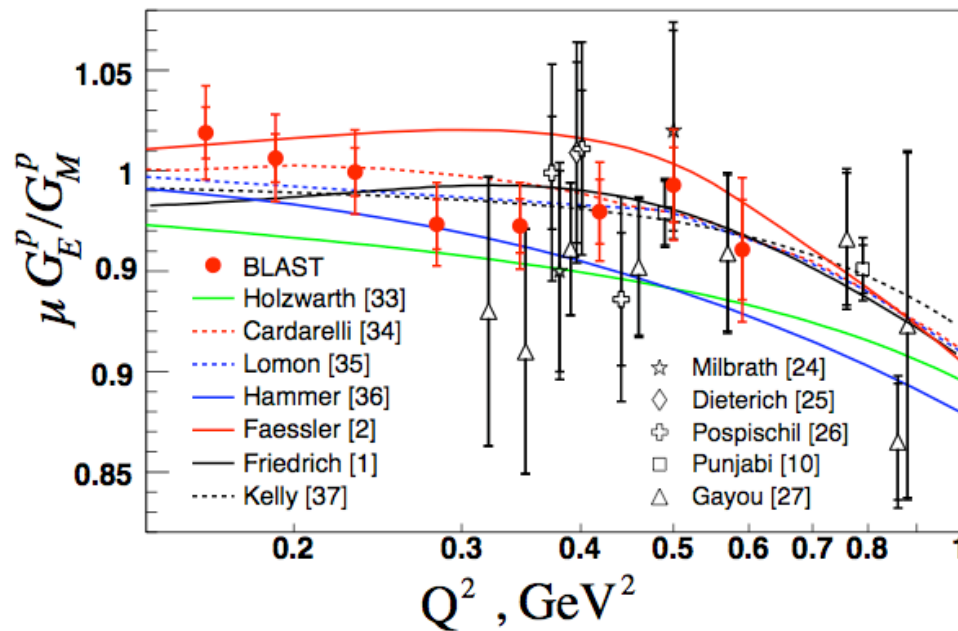
Friedrich&Walcher, EPJ (2003)
have found a bump in all four FFs
relatively to a two-dipoles fit

Belushkin et al, PRC (2007) get a good
description of most data with **dispersion
analysis** including meson continua

Revived interest in precision experiments at Q^2 range below 1 GeV²

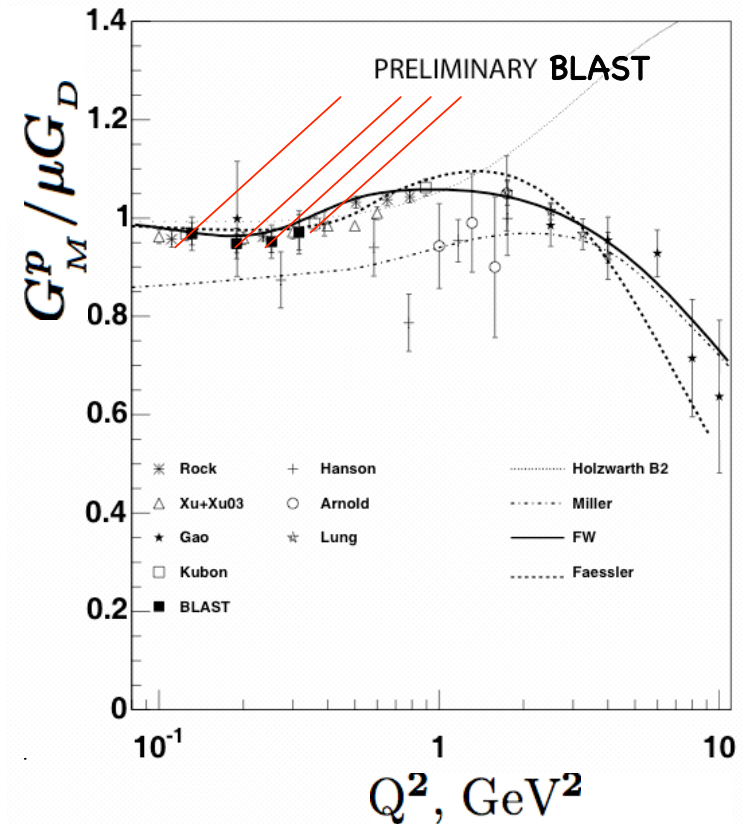
Low Q^2 FFs from BLAST

Crawford et al., PRL 98, 052301 (2007)



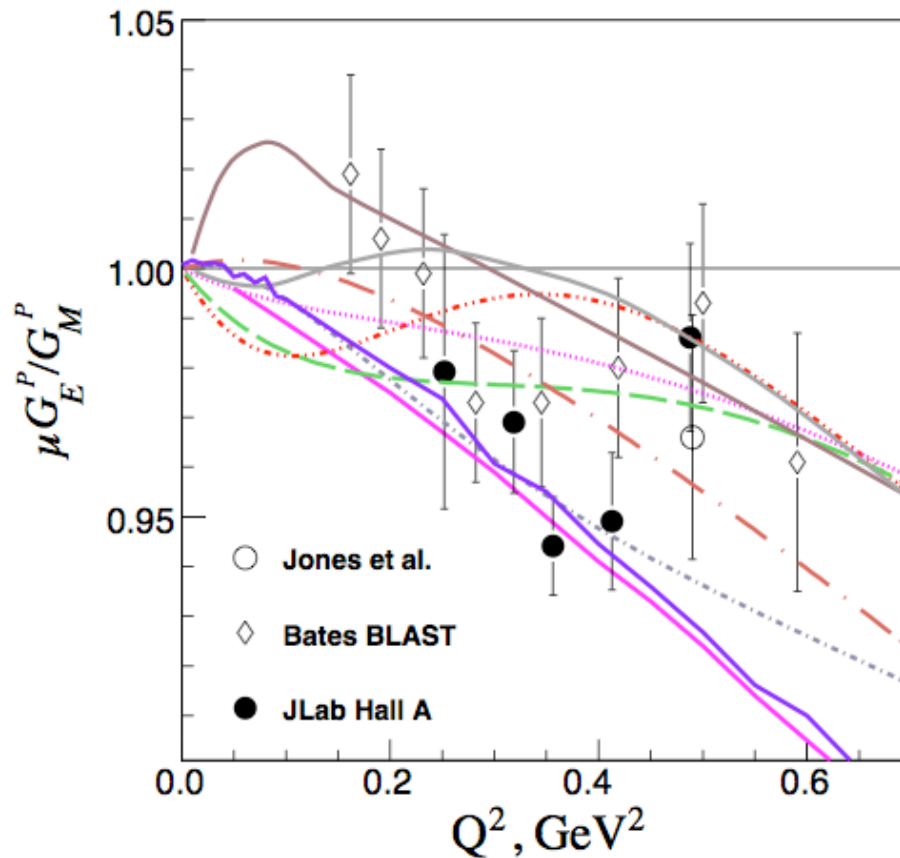
Consistent with unity in 0.2-:-0.6 GeV^2

Alarcon, Eur.Phys.J. A32,477 (2007)



Show a dip at 0.2 GeV^2

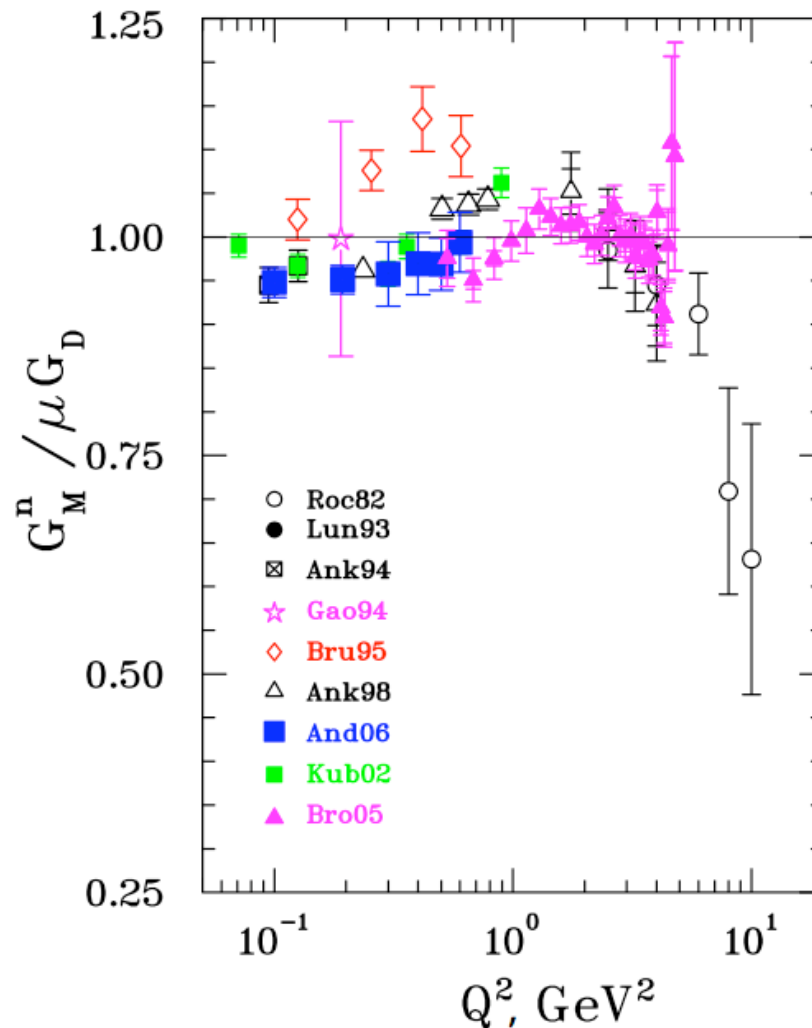
Low Q^2 GEP/GMP



New polarization transfer data
from JLab Hall A:
G.Ron at al., arXiv:0706.0128

A ratio less than unity in range from 0.2 to 0.5 GeV^2

Low Q^2 GMN inconsistency



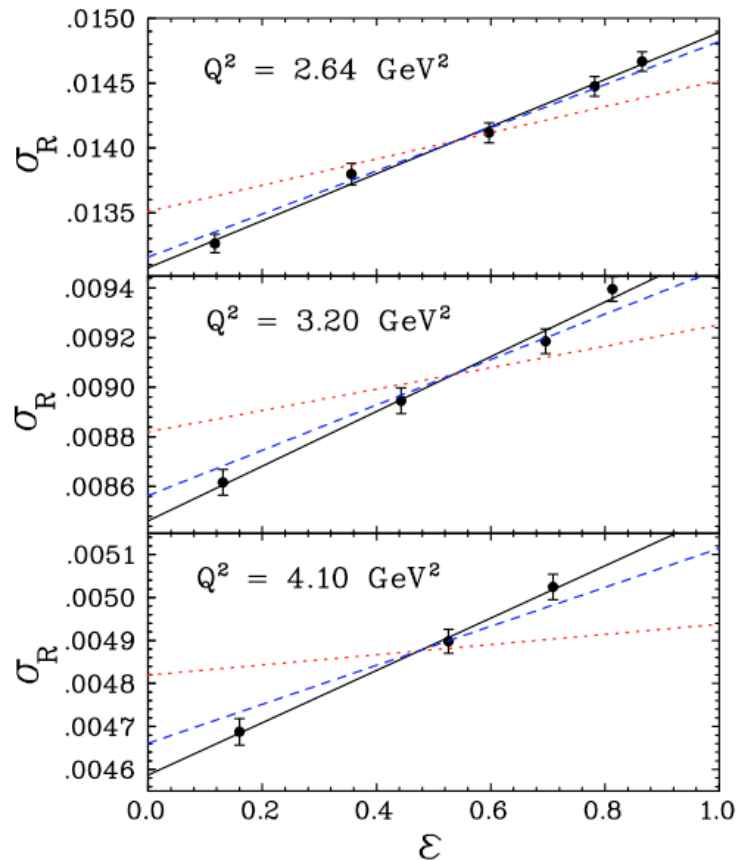
A systematic difference of several % between results (■ ■ ▲) in Q^2 range 0.4 – 0.8 GeV^2 . A final analysis and paper from CLAS is coming soon.

Reminder that at least two independent experiments are always needed.

Super ratio Rosenbluth

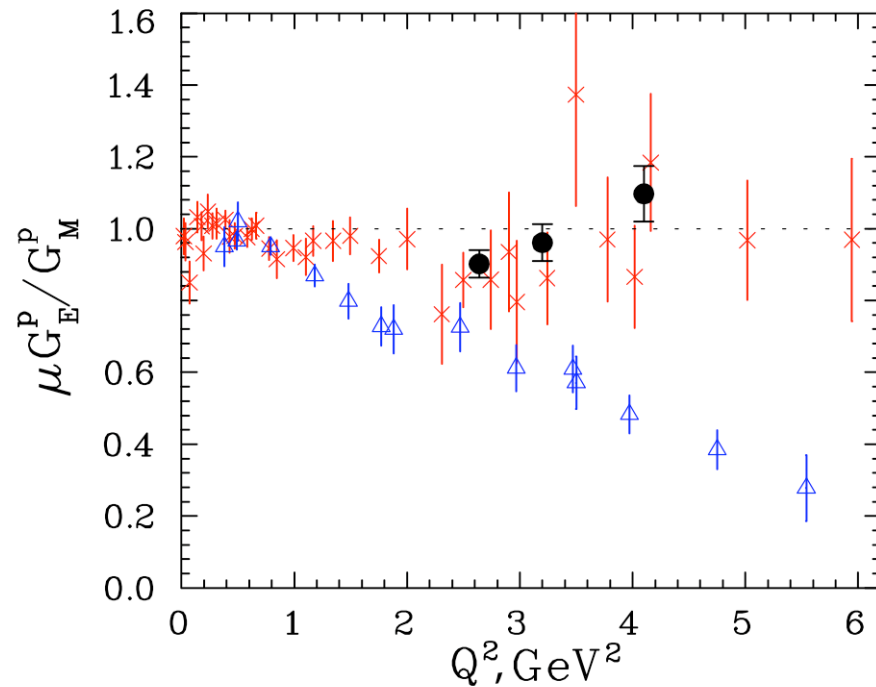
Arrington & Segel

$$\sigma_{\varepsilon 1}/\sigma_{\varepsilon 2}/\sigma_{\varepsilon 3} \dots$$



$$H(e, p)e'$$

Qattan et al., PRL
94, 142301 (2005)



New L/T experimental results consistent with old set

High Q^2 GEN

- ✓ Since 1984, when Blankleider&Woloshin suggested ${}^3\vec{H}e(\vec{e}, e'n)$ several experiments of this type were performed at NIKHEF-K and Mainz (A1, A3) for Q^2 up to 0.7 GeV^2 , big success in part due **new accurate 3-body calculation possible at low Q^2** (Glockle et al.)
- ✓ At Q^2 **above $1\text{--}2 \text{ GeV}^2$** Glauber method becomes sufficiently accurate (Sargsian)
- ✓ Electron-polarized neutron luminosity and high polarization of ${}^3\text{He}$ target made **measurement about 10 times** more effective than with ND_3 . In combination with a large acceptance electron spectrometer the total enhancement is **more than 100**, which allows to reach 3.5 GeV^2

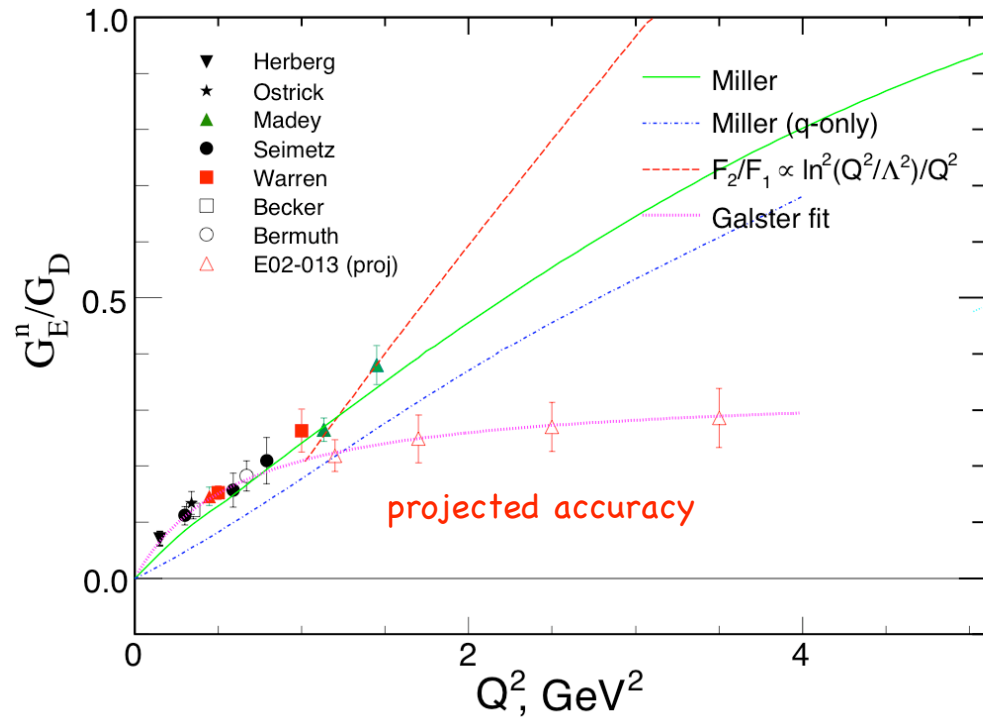
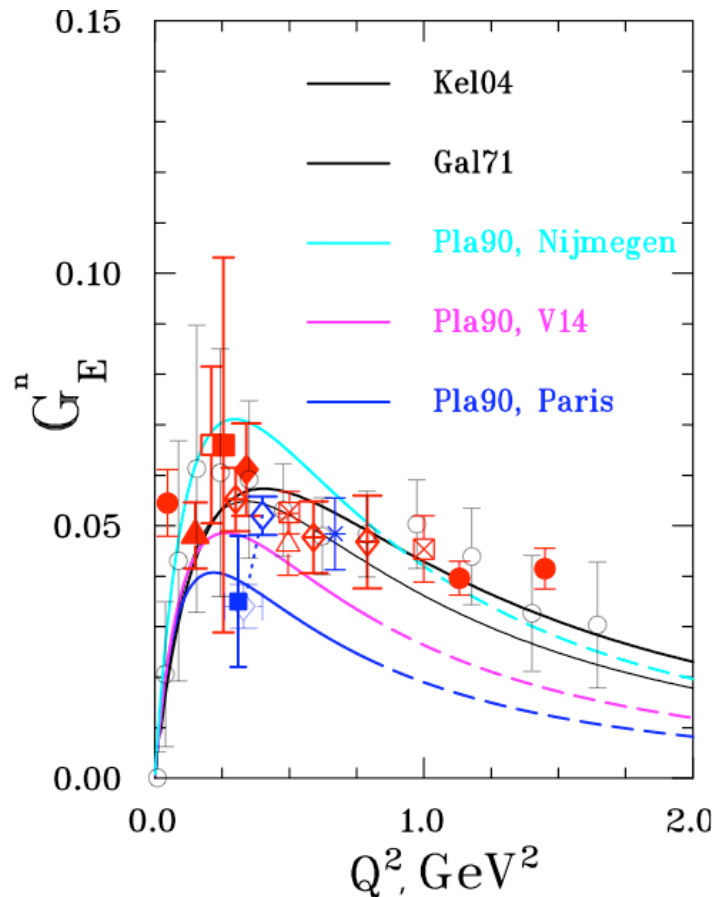
Require sup.

- Polarized target
- Electron spectrometer
- Neutron detector

Hall A GEN (E02-013)

G.Cates, N.Liyanage, and BW

$${}^3\vec{H}e(\vec{e}, e'n)$$

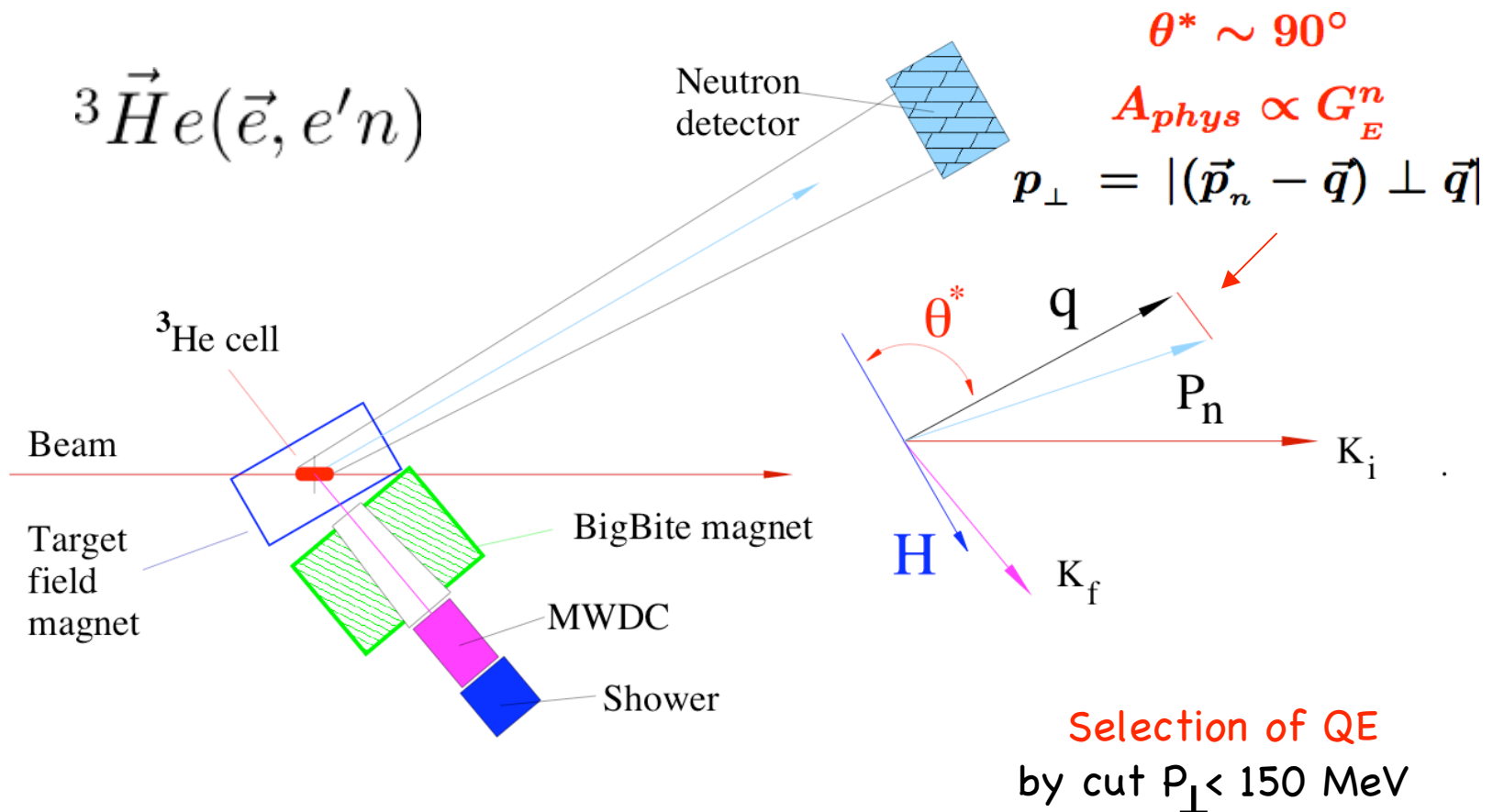


Factor 2 increase of Q^2 require
a 100 larger Figure-of-Merit

E02-013 scheme

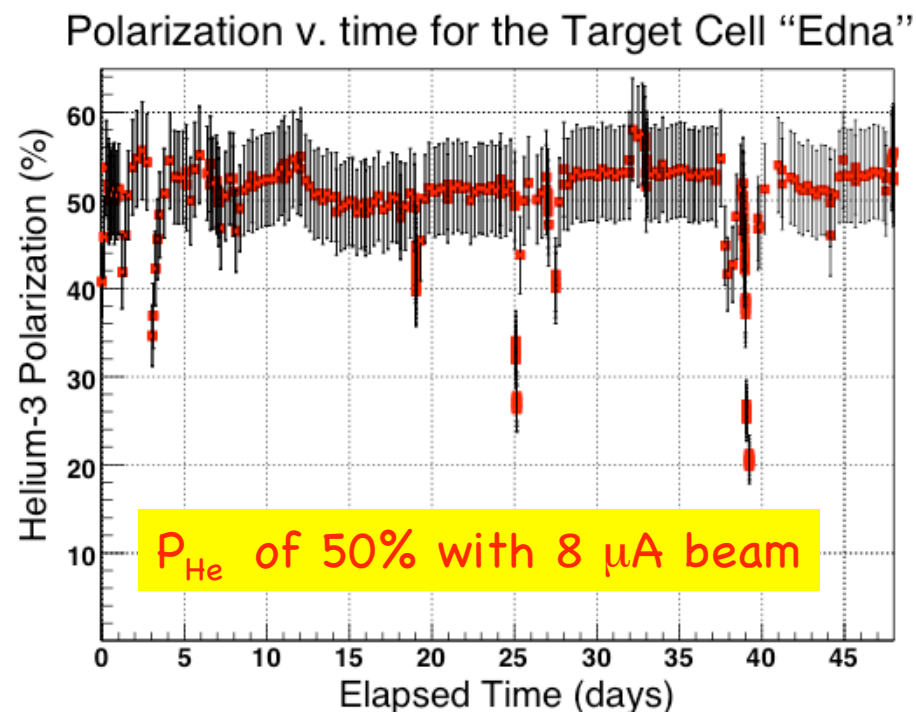
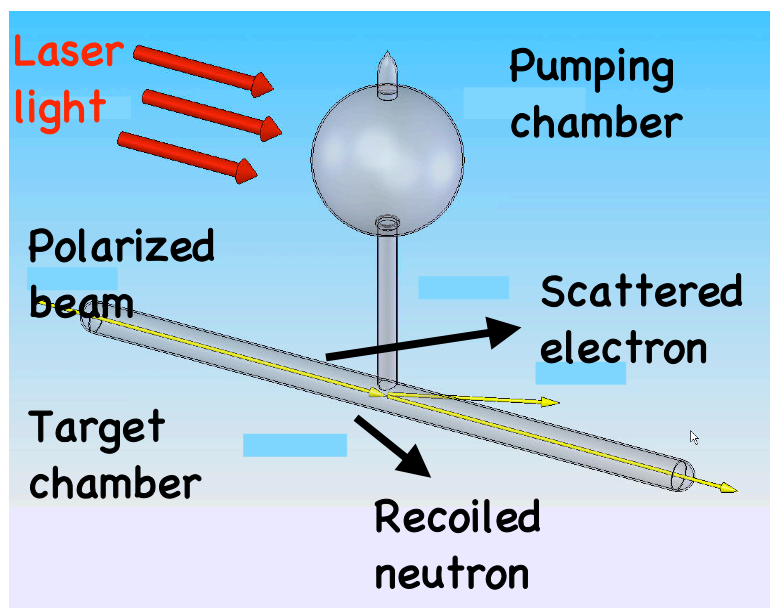
$$A_{phys} = -C1_{kin} \cdot G_E^n G_M^n \sin \theta^* \cos \phi^* - C2_{kin} \cdot (G_M^n)^2 \cos \theta^*$$

$${}^3\vec{H}e(\vec{e}, e' n)$$



Polarized target

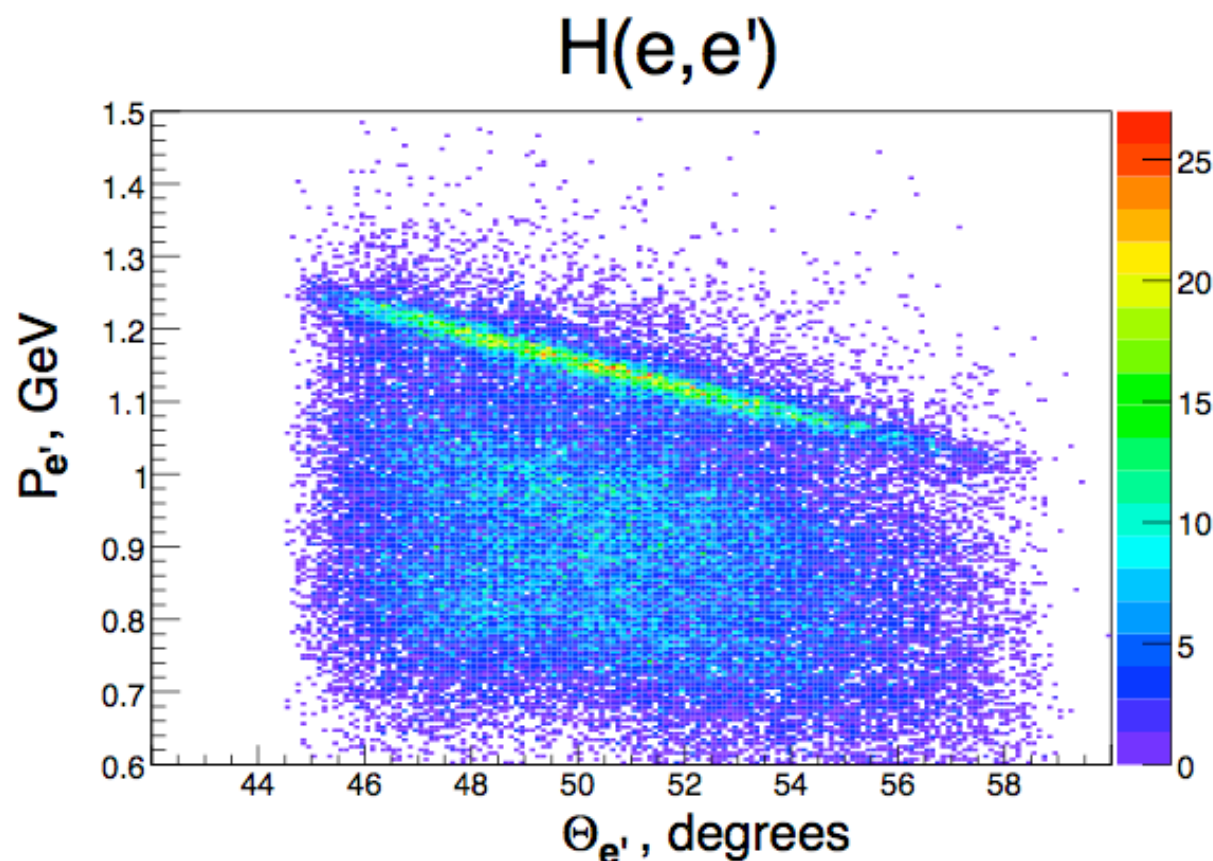
$${}^3\text{He} = p + p + n$$
$$S + S' + P \text{ waves}$$
$$P_n = 0.86 P_{\text{He}}$$



Rb + K mixture have shorten spin-up time to 5-8 hours. Hybrid method for the first time in actual target.

Electron Spectrometer

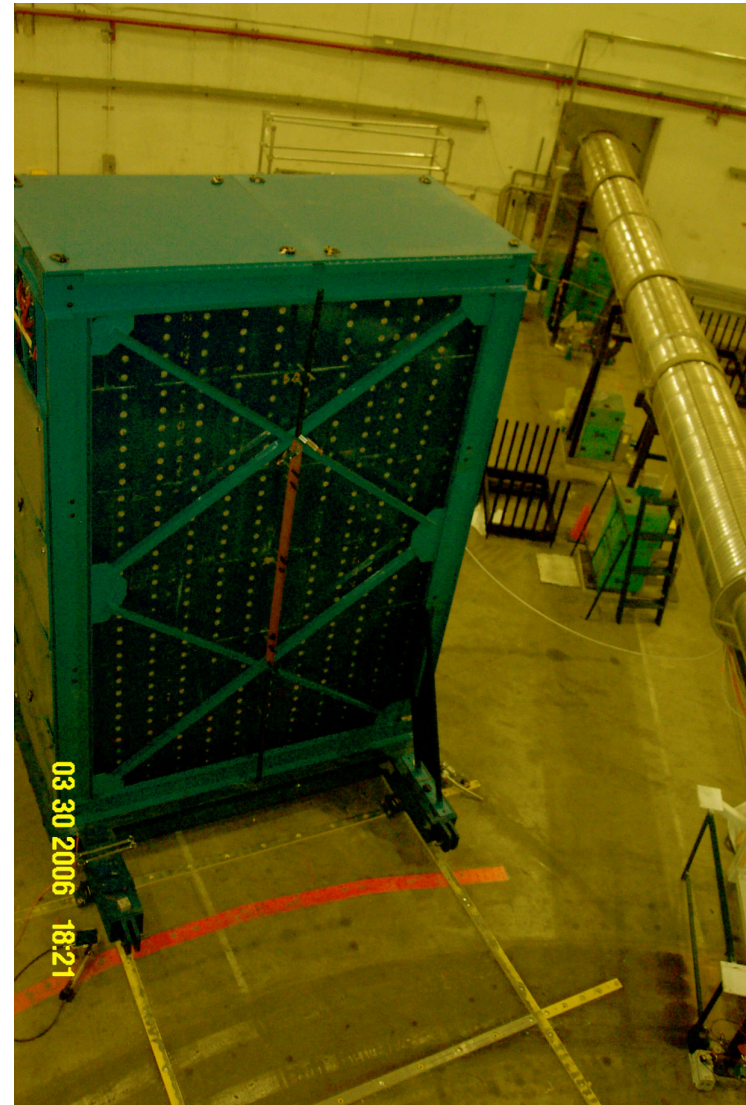
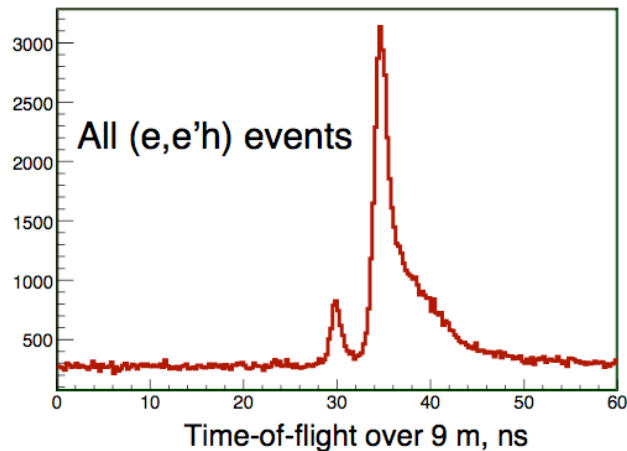
Useful $\Delta Q^2/Q^2 \sim 0.1$ and $\max \Omega$ leads to a large aspect ratio, limited just of 30° for the polar. target. BigBite was designed at NIKHEF for aspect ratio $\Delta\theta/\Delta\phi = 1/5$. Spectrometer has solid angle up to **95 msr**.



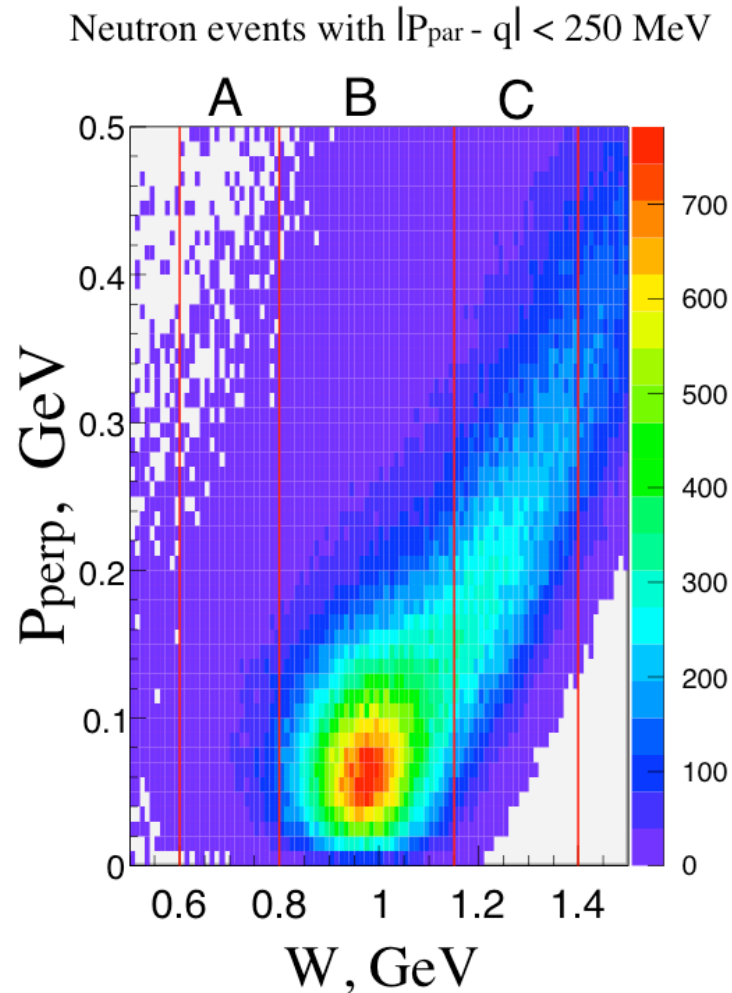
With luminosity of JLab polarized target, $10^{37} \text{ cm}^{-2}/\text{s}$, the open geometry - a dipole spectrometer - works well when all MWDCs located behind the magnet.

Neutron Detector

- Match BigBite solid angle for QE kinematics
- Flight distance ~ 10 m
- Operation at $3 \cdot 10^{37}$ cm²/s
- 1.6×5 m² active area
- 6–7 layers (~ 250 bars)
- 2 veto layers (~ 200)
- 0.38 ns time resolution

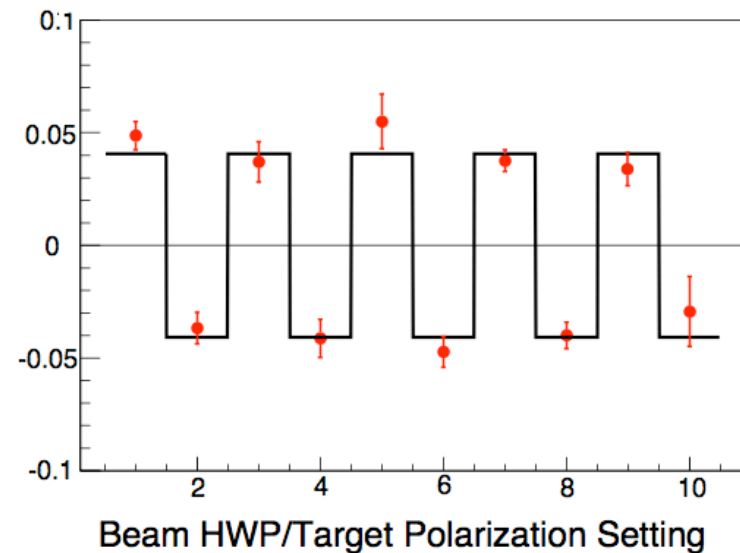


Data analysis



Selection of QE ($e, e'n$) events

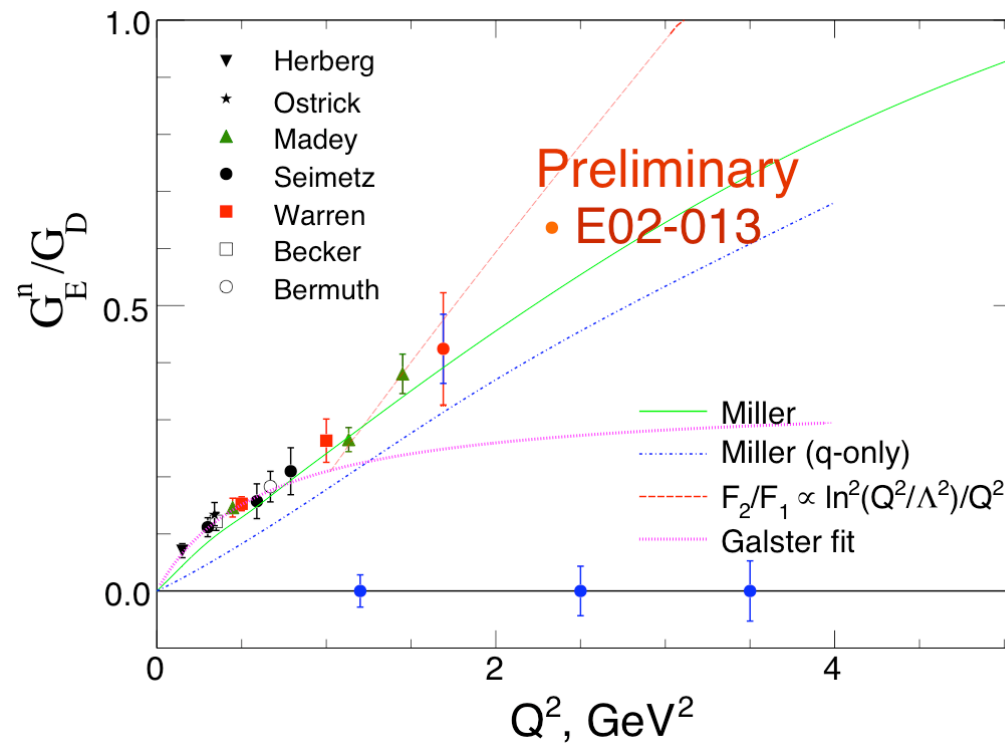
Observed Asymmetry for Quasi-elastic Neutrons



Asymmetry then corrected for

1. p-n identification
2. A_{\parallel} contribution
3. FSI for $e, e'n$ process
4. Target, beam polarizations

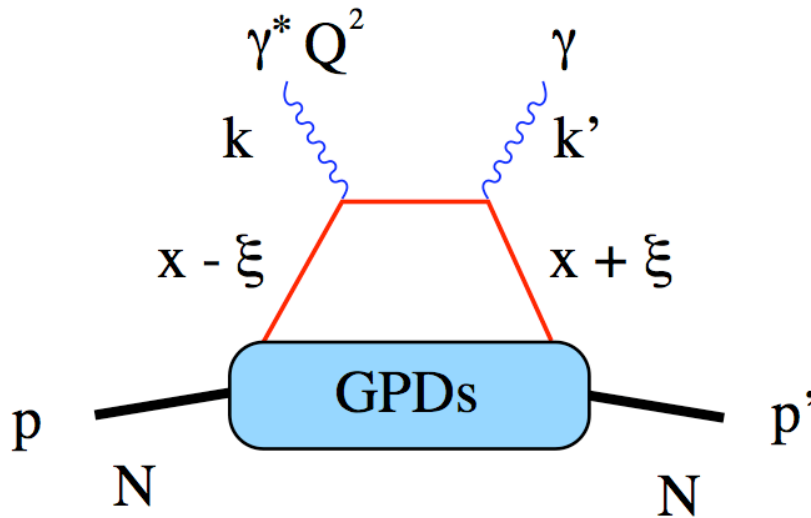
First physics result from Hall A GEN



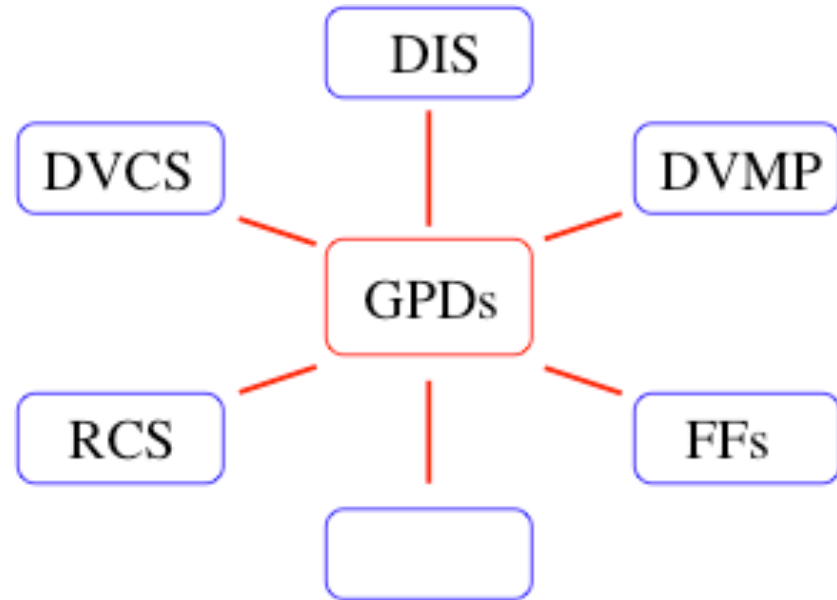
- Result is well above Galster.
- Nuclear corrections include neutron polarization and estimate of Glauber (~5%).
- Present error (~20%) dominated by preliminary “neutron dilution factor”, and is expected to be ~7% stat. and 8% syst. with further analysis.
- 3.4 GeV² result to be released in October DNP meeting at Newport News.

GPDs of nucleon

Muller (94), Ji (97), Radyushkin (97)



where $\xi = (p_q^+ - p_q'^+)/ (p_q^+ + p_q'^+)$



Quark dynamics of nucleon encoded in GPD functions

$H(x, \xi, t)$, $\tilde{H}(x, \xi, t)$ conserve hadron helicity;

$E(x, \xi, t)$, and $\tilde{E}(x, \xi, t)$ flip hadron helicity

GPDs and Form Factors

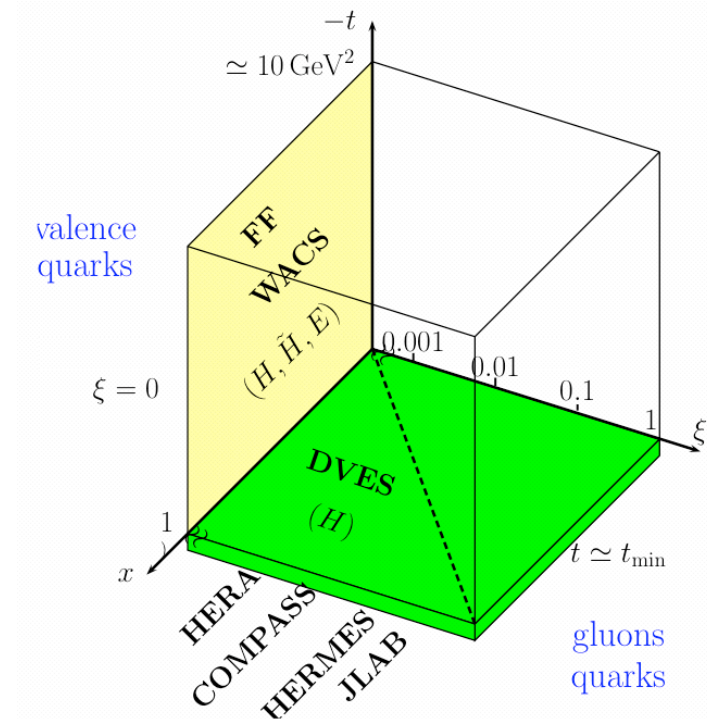
Reduction formulas at $\xi = t = 0$
for DIS and $\xi = 0$ for FFs

$$H^q(x, \xi = 0, t = 0) = q(x)$$

$$\tilde{H}^q(x, \xi = 0, t = 0) = \Delta q(x)$$

$$\int_{-1}^{+1} dx H^q(x, 0, Q^2) = F_1^q(Q^2)$$

$$\int_{-1}^{+1} dx E^q(x, 0, Q^2) = F_2^q(Q^2)$$



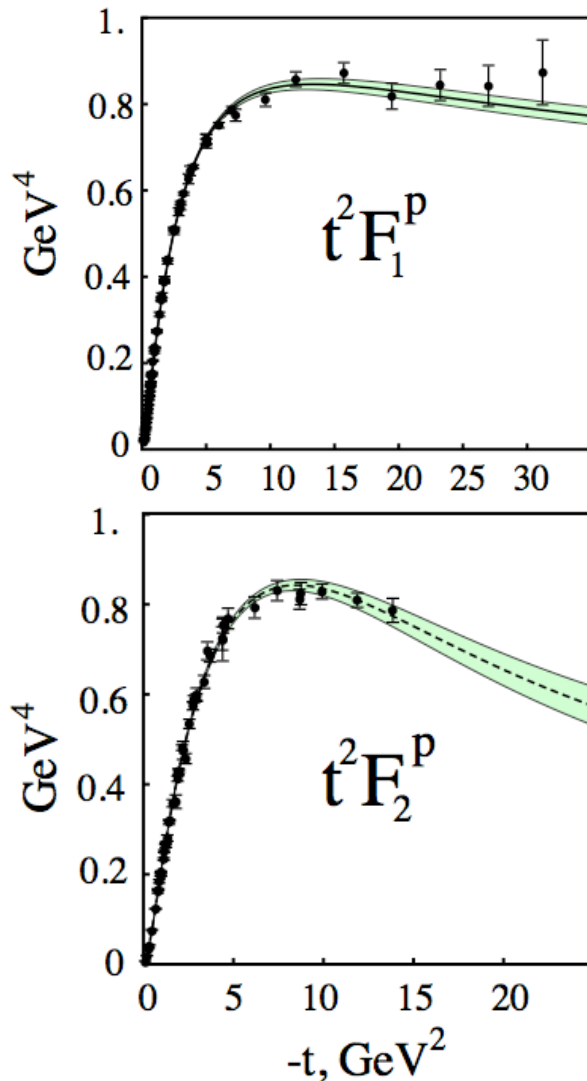
Ji's sum rule for quark orbital momentum

$$\langle L_v^q \rangle = \frac{1}{2} \int_0^1 dx [x E_v^q(x, \xi = 0, t = 0) + x q_v(x) - \Delta q_v(x)]$$

DVCS will access low t , large Q^2 kinematics

FFs presently is the main source for E_v^q

Model of GPD and Form Factors



Diehl et al (2005),
Guidal et al (2005)

use **all** available data on $G_M^p, G_M^n, G_E^p, G_E^n, F_A$

and CTEQ6 parametrization of $q(x), \Delta q(x)$

in order to determine $H_v^{u,d}, \tilde{H}_v^{u,d}, E_v^{u,d}$

ANSATZ: $H_v^q(x, t) = q_v(x) \exp[f_q(x)t]$

$$f_q = [\alpha' \log(1/x) + B_q] (1-x)^{n+1} + A_q x (1-x)^n$$

fixed $\alpha' = 0.9 \text{ GeV}^{-2}, n = 1, 2;$

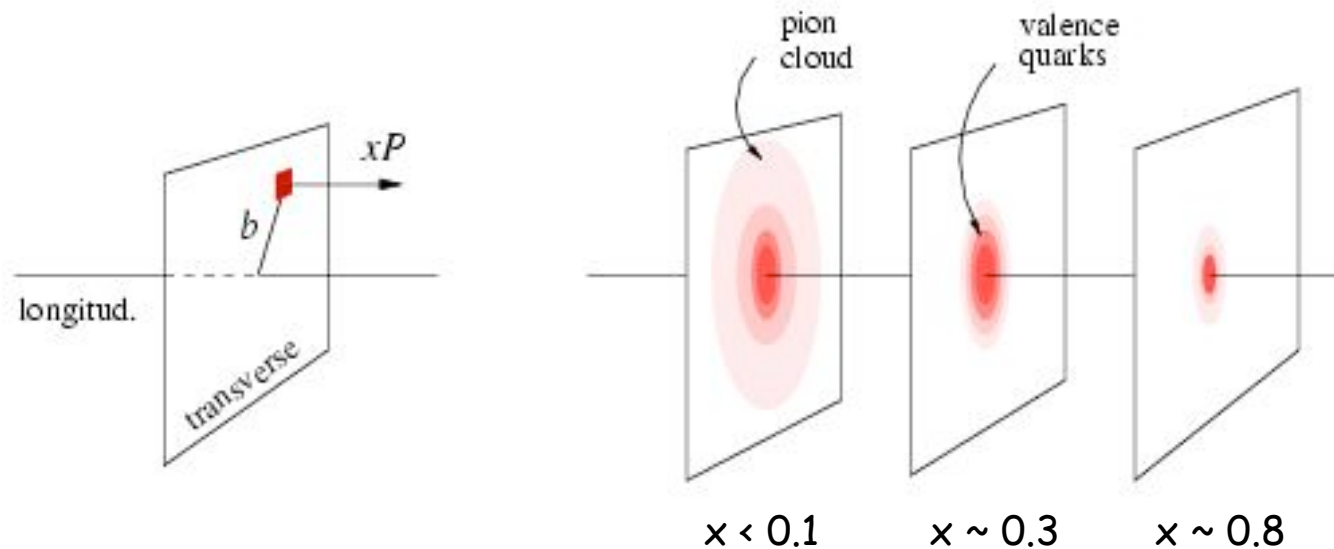
only A_q , and B_q parameters are fitted

GPDs and impact parameter

Transverse momentum invariance
allows frame independent Fourier
transform from $H(x, \xi, t)$ to $q(x, \xi, b)$
with impact parameter b defined
relative to center of momentum

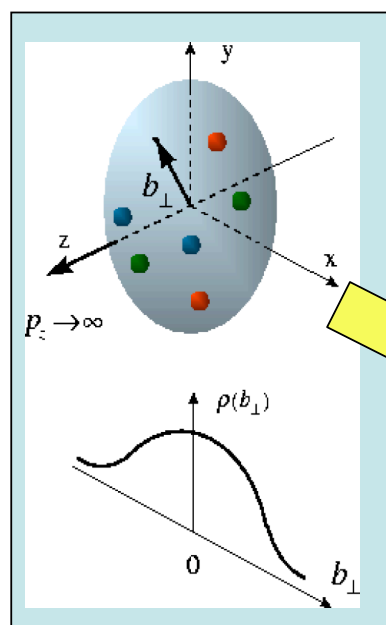
Burkardt, Int. J. Mod.
Phys. A 18, 173 (2003)

Fourier transform in momentum transfer

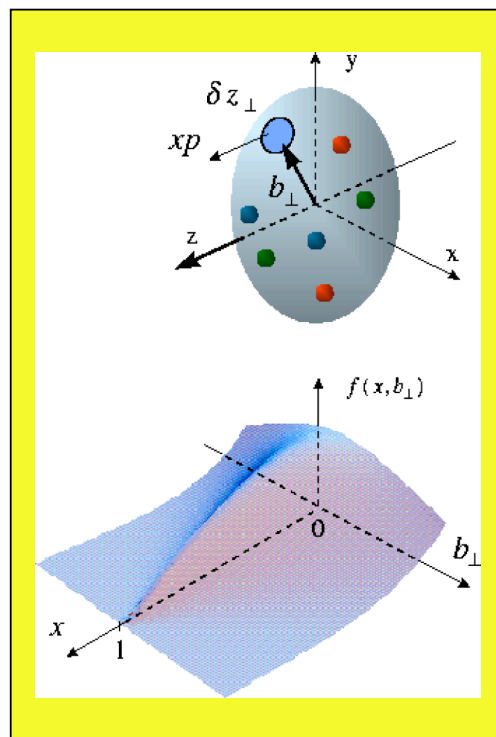


gives transverse size of quark (parton) with longitud. momentum fraction x

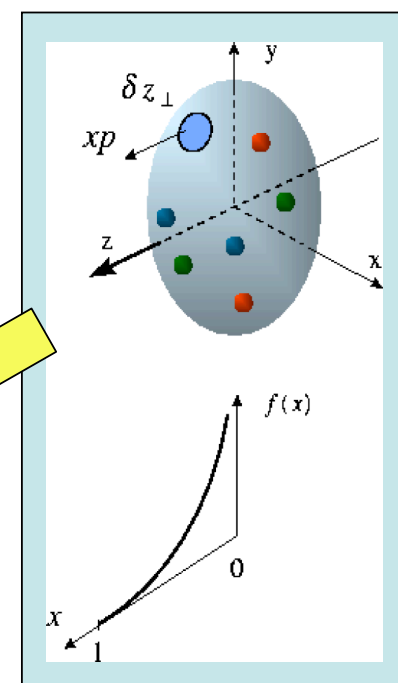
3-d picture of the nucleon



Proton form factors,
transverse charge &
current densities



Correlated quark momentum
and helicity distributions in
transverse space - GPDs



Structure functions,
quark longitudinal
momentum & helicity
distributions

Nucleon Density from GPD

$$F_1(t) = \sum_q e_q \int dx H_q(x, t)$$

Muller, Ji, Radyushkin

$$q(x, b) = \int \frac{d^2 q}{(2\pi)^2} e^{i q \cdot b} H_q(x, t = -q^2)$$

M.Burkardt

$$\rho(b) \equiv \sum_q e_q \int dx q(x, b) = \int d^2 q F_1(q^2) e^{i q \cdot b}$$

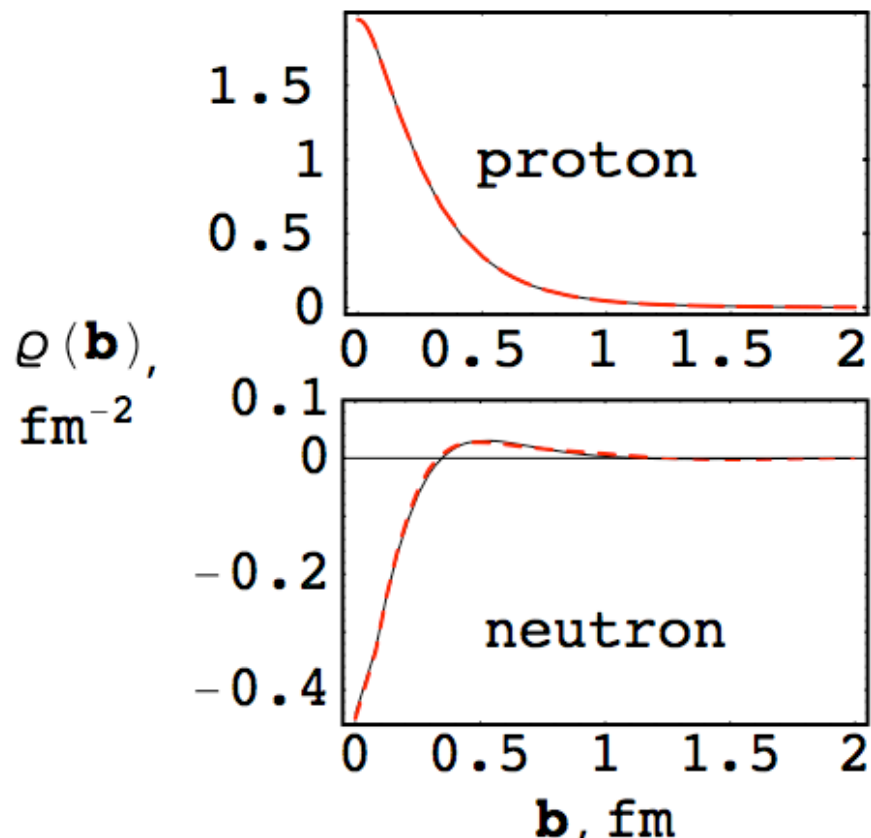
$$\rho(b) = \int_0^\infty \frac{Q \cdot dQ}{2\pi} J_0(Qb) \frac{G_E(Q^2) + \tau G_M(Q^2)}{1 + \tau}$$

G.Miller,
arXiv:0705.2409

center of momentum $\mathbf{R}_\perp = \sum_i \mathbf{x}_i \cdot \mathbf{r}_{\perp,i}$

b is defined relative to \mathbf{R}_\perp

Neutron is negative inside



G.Miller, arXiv:0705.2409

using FFs from

Kelly's fit

BBBA's fit

Does it contradict intuition ?

Static picture: a neutron is
a proton in the center plus π^-

Intuitive picture of static
constituent quarks is not
applicable for large Q^2
where quark DFs play role

- Negative density at low b in a neutron \Rightarrow d quarks dominate
- High Q^2 elastic process in Feynman mechanism requires a large x quark \Rightarrow d quarks dominate at large x , in agreement with DIS

GEP-III: G_E^p/G_M^p for 8.6 GeV^2

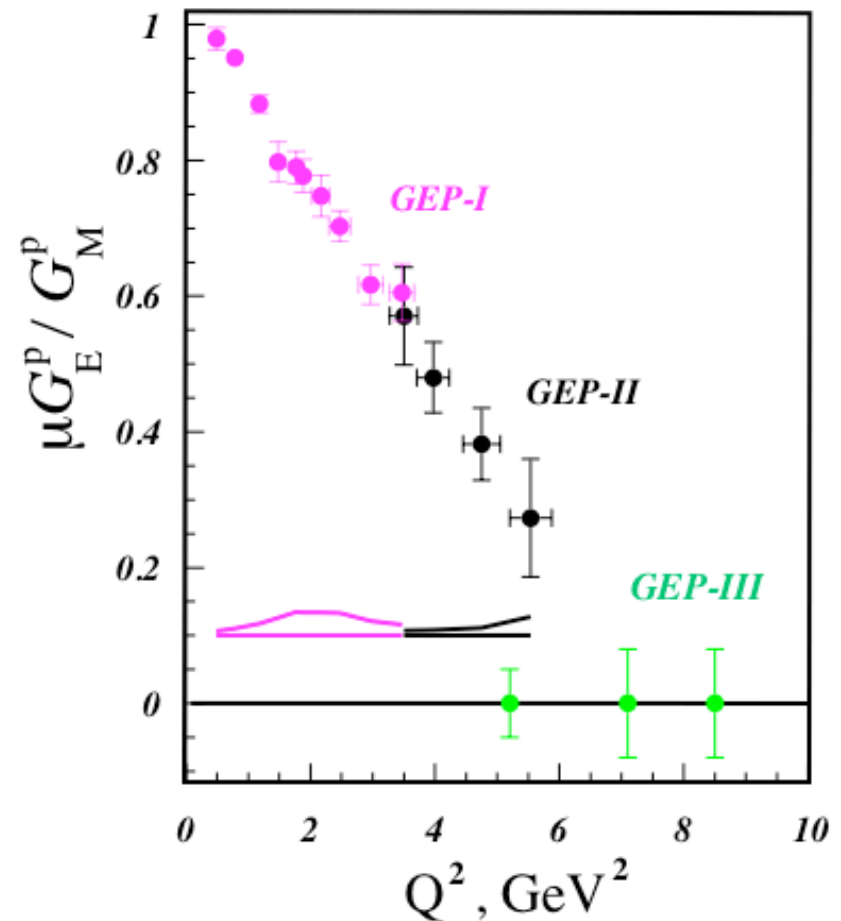
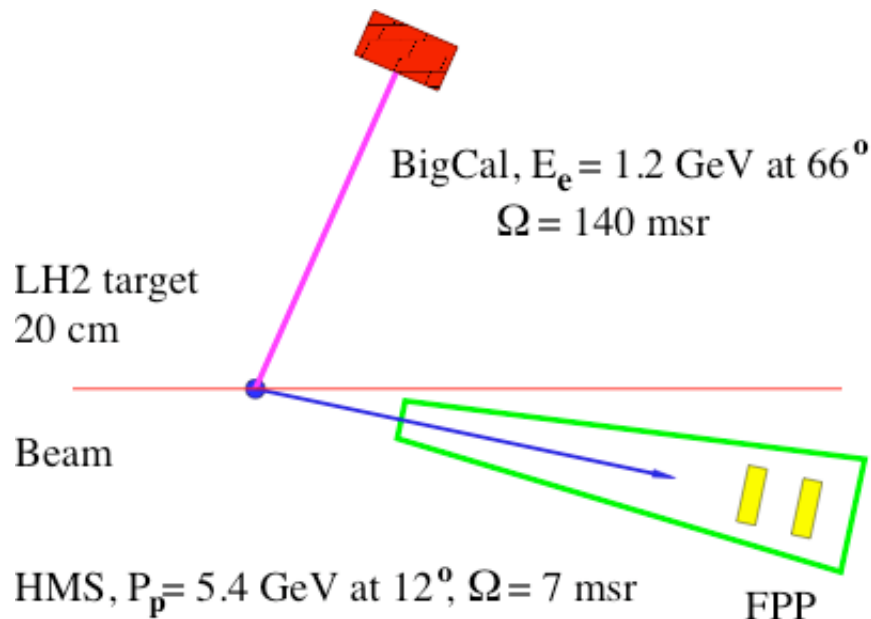
Brash, Jones, Perdrisat, Punjabi

Polarization transfer in $H(\vec{e}, e' \vec{p})$

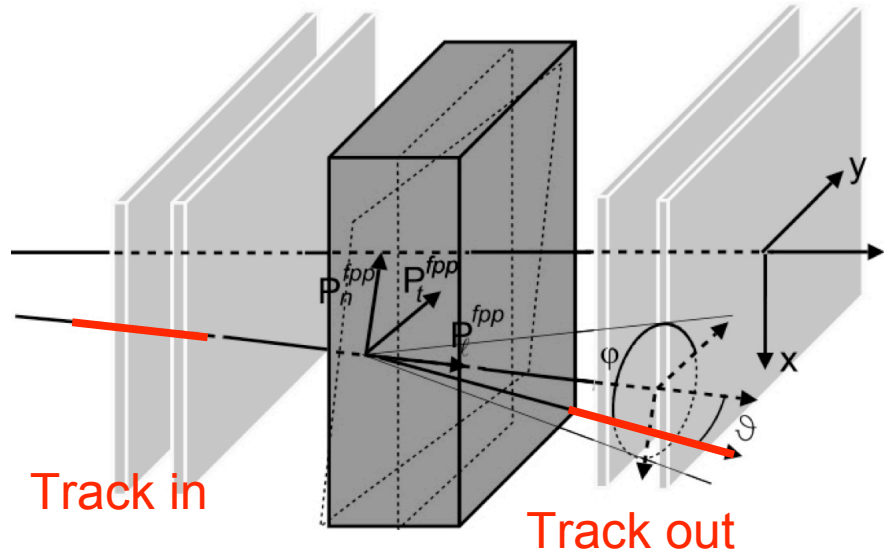
- New detector for scattered electron
- New polarimeter for recoiled proton

Commissioning will start next month

High Q^2 data will be taken by 5/08



Focal Plane Polarimeter

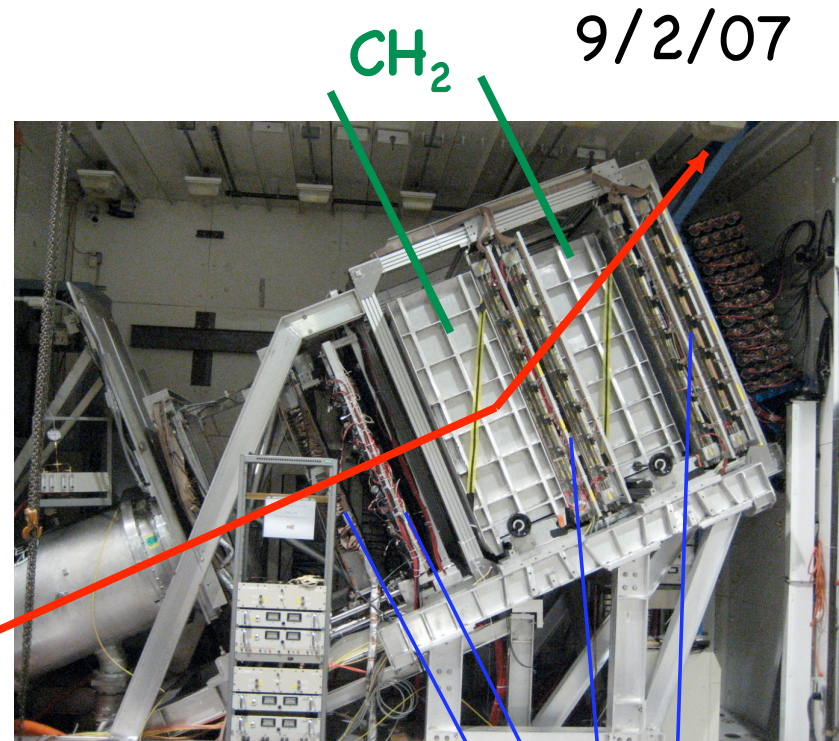


$$f^{\pm}(\vartheta, \varphi) = \frac{\epsilon(\vartheta, \varphi)}{2\pi} \left[1 \pm A_y (P_x^{fpp} \sin \varphi - P_y^{fpp} \cos \varphi) \right]$$

where \pm refers to electron beam helicity

$$A = \frac{f^+ - f^-}{f^+ + f^-} = A_y (P_x^{fpp} \sin \varphi - P_y^{fpp} \cos \varphi)$$

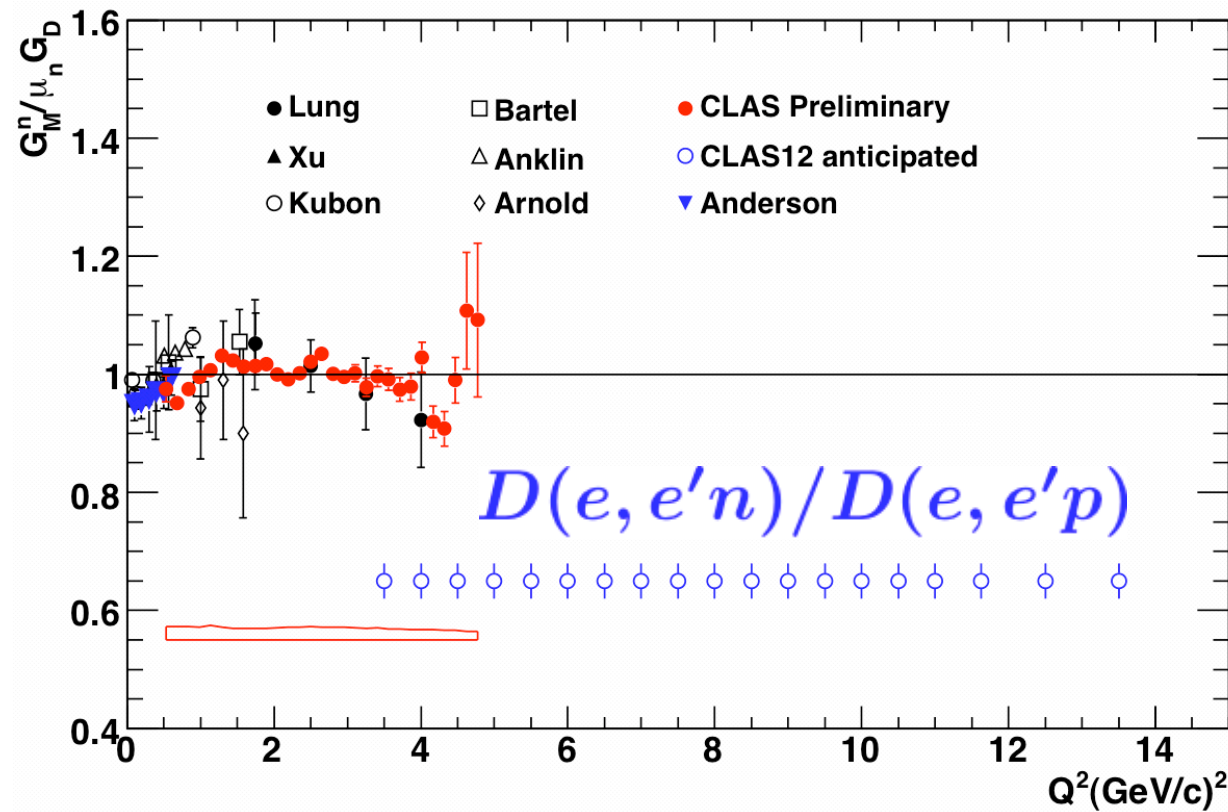
$$\mu_p \frac{G_E^p}{G_M^p} = -\mu_p \frac{E_e + E'_e}{2M_p} \tan \frac{\theta_e}{2} \left(\frac{P_x^{fpp}}{P_y^{fpp}} \sin \chi_\theta + \gamma_p (\mu_p - 1) \Delta \phi \right)$$



MWDCs

GMN-14 with CLAS++

Gilfoyle, Brooks, Vineyard, Hafidi, Lachniet

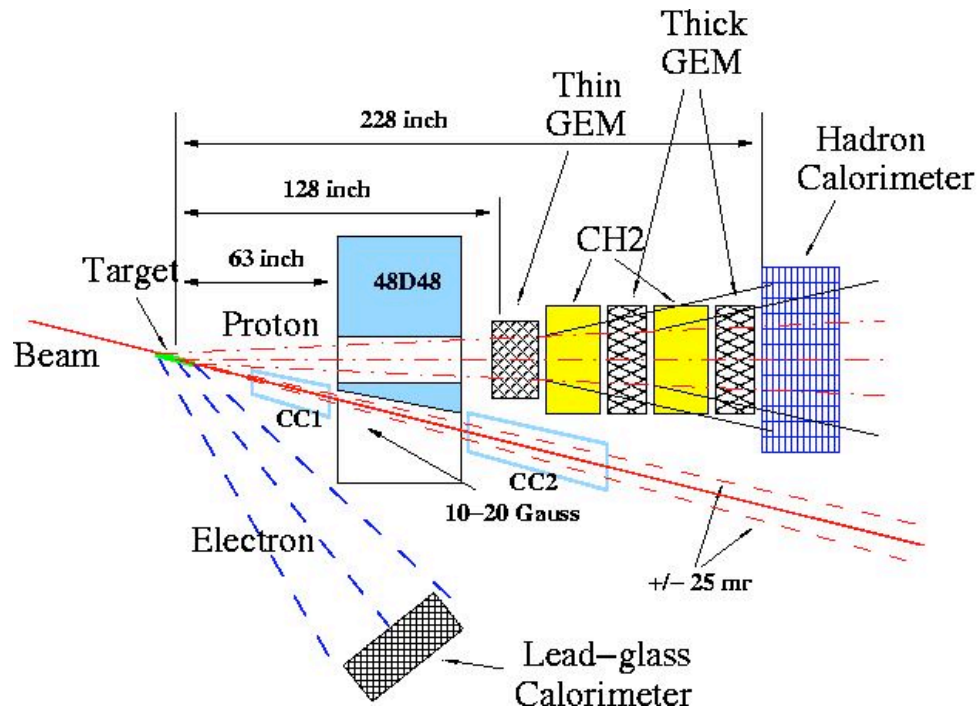


approved by PAC32
for 12 GeV program

GEP-15: G_E^p/G_M^p up to 15 GeV^2

Perdrisat, Pentchev, Cisbani, Punjabi, BW

$$H(\vec{e}, e' \vec{p})$$



Beam: $75 \mu\text{A}$, 85% polarization

Target is 40 cm liquid H_2

Electron arm at 37° , covers

$$Q^2 = 12.5 \text{ to } 16 \text{ GeV}^2$$

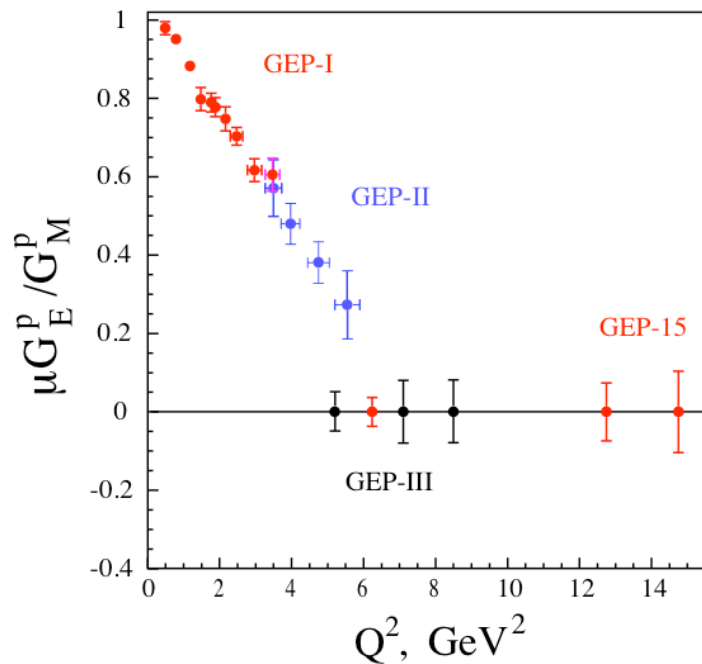
Proton arm at 14° , $\Omega \sim 35 \text{ msr}$

58 days of production time
resulting accuracy:

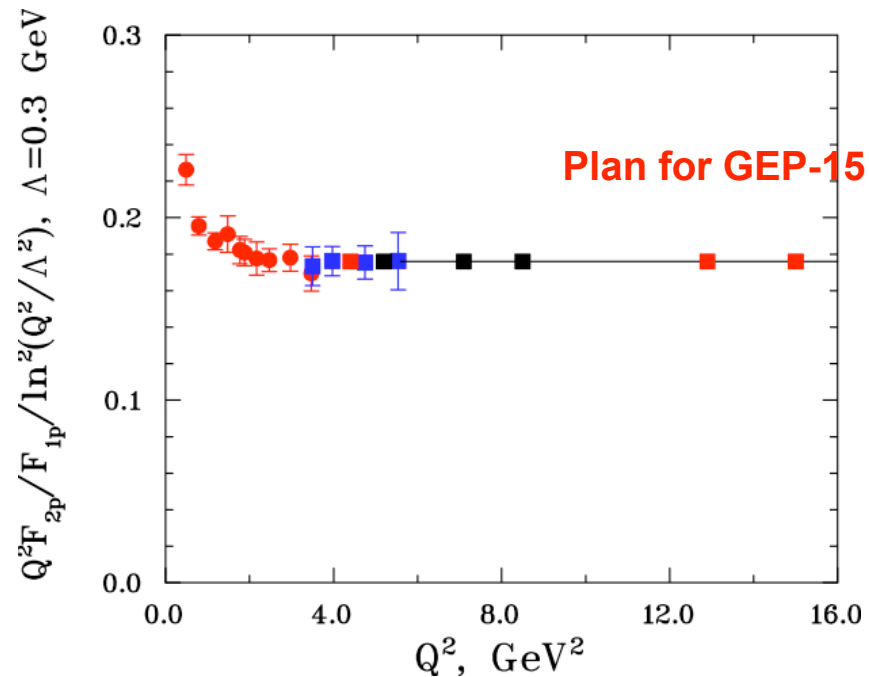
$$\Delta(\mu G_E^p/G_M^p) = \pm 0.10$$

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GEP-15: Projected accuracy



$$\Delta(\mu G_E^p / G_M^p) = \pm 0.10$$



$\Delta(F_2/F_1)/(F_2/F_1)$ accuracy will be **3%**

$$\text{compare to } \frac{\ln^2(Q^2=10/\Lambda^2)}{\ln^2(Q^2=15/\Lambda^2)} = 0.85$$

GMP-18: New measurement of G_M^p

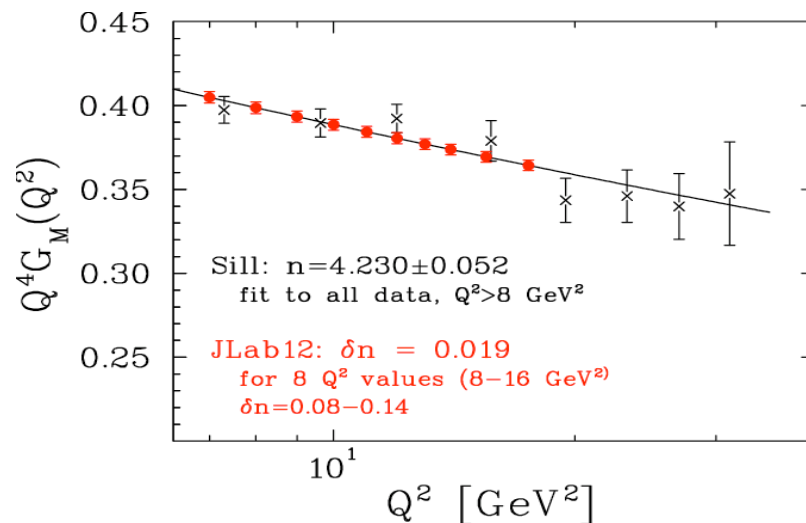
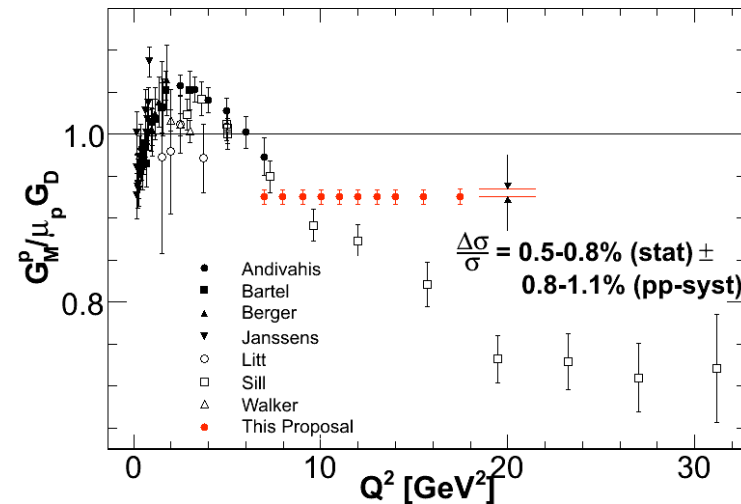
Moffit, Gilad, Arrington, BW

The cross section of $H(e,e')p$.

By using two existing Hall A High Resolution Spectrometers with several new ideas for improved control of systematic.

With 11 GeV beam in 31-day run.

approved by PAC32
for 12 GeV program



Summary

- ❑ Experiment and theory have created improved base for understanding of nucleon
- ❑ Future experiments will provide precision FFs data for Q^2 up to 7/14/15/18 GeV^2
- ❑ GPD approach, as expected, shed light into to the nucleon structure
- ❑ Expectation of Lattice QCD results is boiling

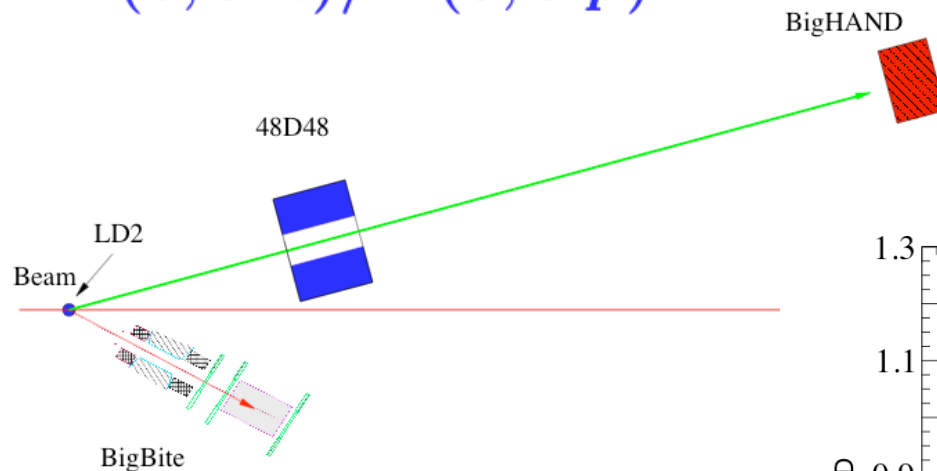
It is exciting time for nucleon FF physics

backup
slides
after this

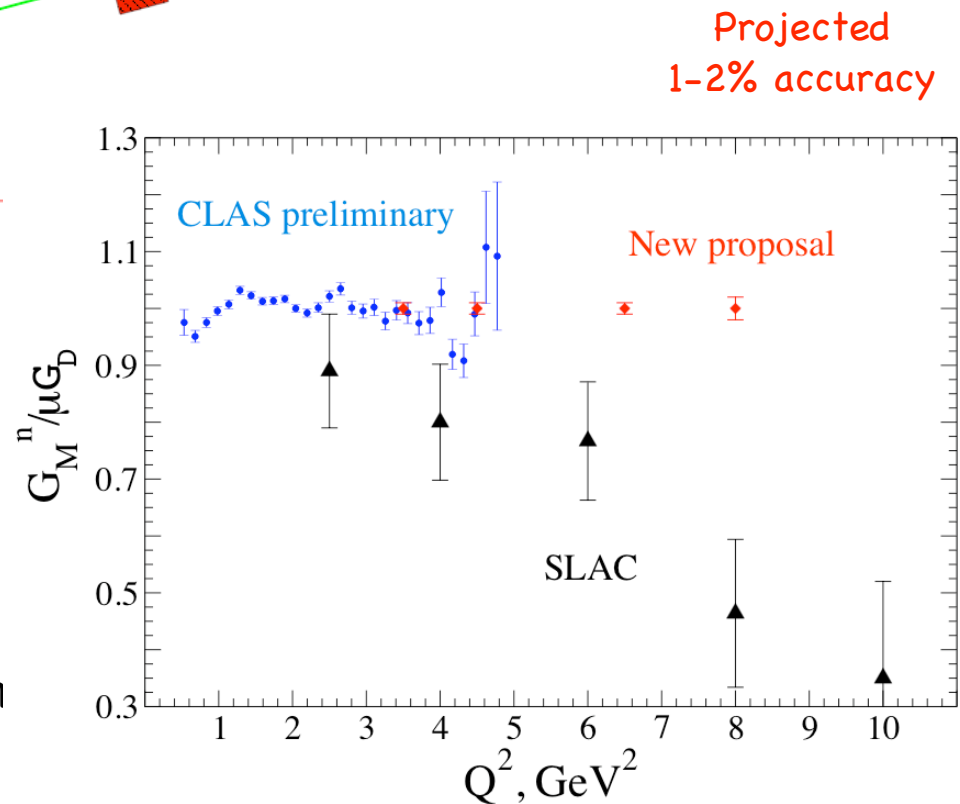
Near future: GMN-8 (for PAC33)

B.Quinn and BW

$$D(e, e'n)/D(e, e'p)$$



- 4–6 GeV beam on 4 cm LD₂
- BigBite spectrometer
- BigHAND neutron detector
- Deflection magnet for proton
- Luminosity $1 \cdot 10^{37} \text{ cm}^{-2}/\text{s}$



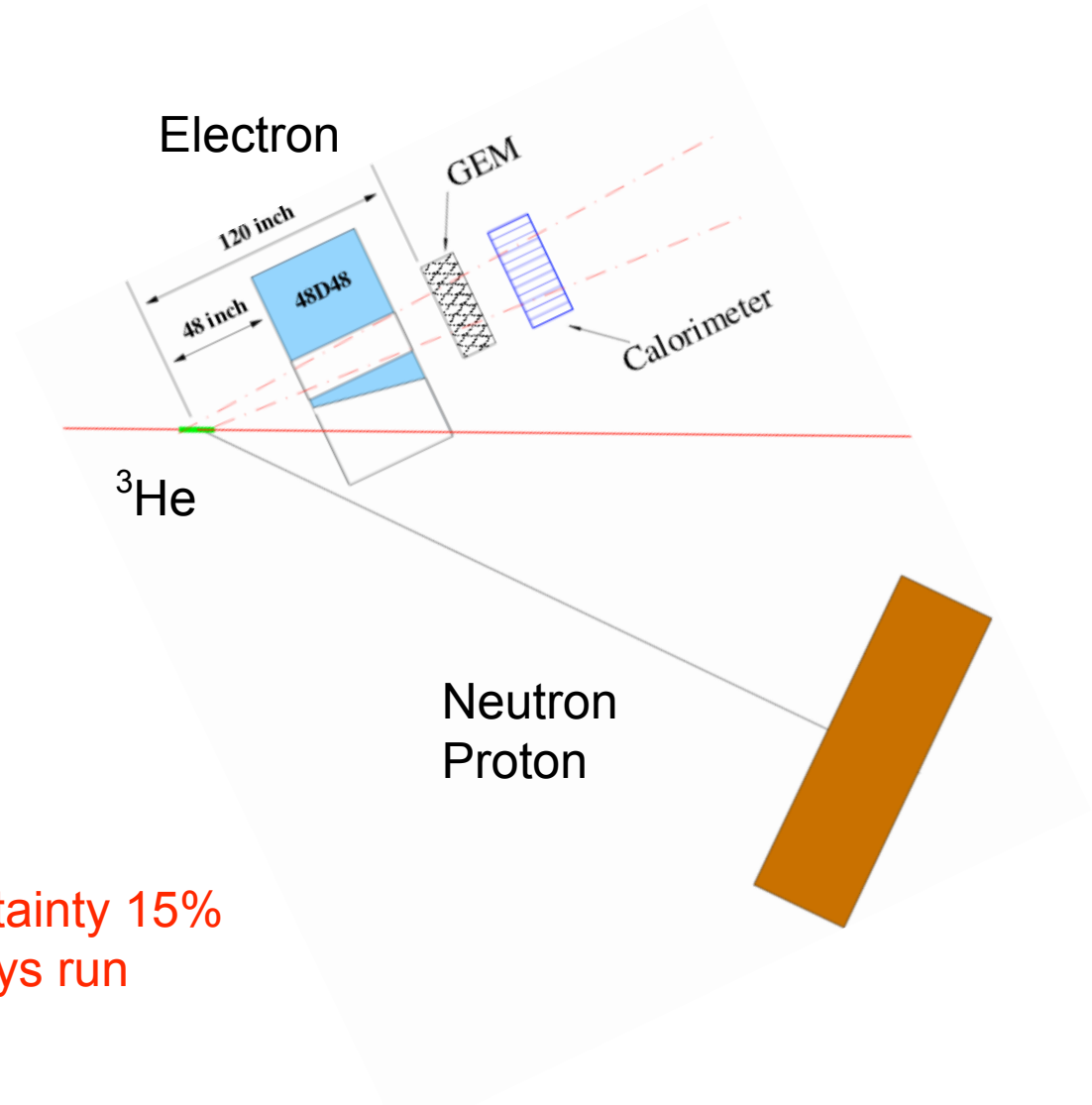
Perspective: G_E^n up to 7 GeV^2

The plan for GEN-7 is:

$${}^3\vec{H}e(\vec{e}, e'n)$$

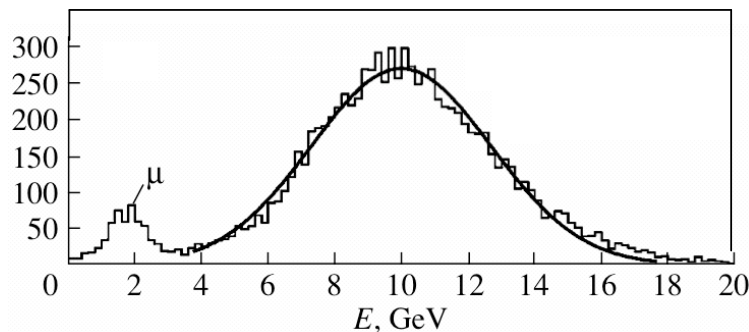
- Beam at 8.8 GeV
- Resolution σ_p/p for electron -
BNL magnet, GEM
- He-3 cell in vacuum, lower
background in neutron arm
- Hybrid He-3 cell with narrow
pumping laser line

G_E^n at 7 GeV^2 with uncertainty 15%
of Miller's value in 30-days run



GEP-15: Proton Arm

- Magnet: 48D48 - 46 cm gap, 3 Tm field integral, 100 ton
- solid angle is 35 msr for GEP, could be ~70 msr at larger angle
GEM chambers for tracking with 70 μm resolution
- momentum resolution is 0.5% for 8.5 GeV/c proton
- angular resolution is 0.3 mrad
- trigger threshold is 4 GeV from
hadron calorimeter



Calorimeter response for 10 GeV protons from test for Compass experiment

