Form of QED Elastic Cross Sections

Seamus Riordan

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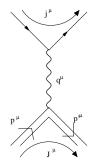
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- Want to find most general form of cross section with QED considerations
- Transition amplitude from $i \rightarrow f$:

$$T_{fi} = -e \int j_{\mu}^{fi} A^{\mu} d^4 x = -e \int j_{\mu}^{fi} \frac{J^{\mu}}{q^2} d^4 x$$
(1)

• First, we must describe general e.m. current J^{μ} for:



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Four vectors available:

$$\boldsymbol{p}^{\mu}, \boldsymbol{p}^{\prime \mu}, \gamma^{\mu} \tag{2}$$

where

$$\gamma^{0} = \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix}$$
(3)
$$\vec{\gamma} = \begin{pmatrix} 0 & \vec{\sigma} \\ -\vec{\sigma} & 0 \end{pmatrix}$$
(4)

Tensors:

$$\sigma^{\mu\nu} = \frac{i}{2} \left(\gamma^{\mu} \gamma^{\nu} - \gamma^{\nu} \gamma^{\mu} \right) \tag{5}$$

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Have to consider all linear combinations

For 4×4 matricies, 16 linearly independent quantities:

$$\mathbb{I}, \gamma^{\mu}, \gamma^{5}, \sigma^{\mu\nu}, \gamma^{\mu}\gamma^{5}$$
(6)

where

$$\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3 \tag{7}$$

Parity operator can be represented $P = \gamma^0$, so:

$$\left\{\gamma^5, P\right\} = 0\tag{8}$$

so currents with $\gamma^{\rm 5}$ ruled out by parity conservation.

- p and p' are the only available 4-vectors
- q^2 is the only independent scalar as all combinations of p and p' can be reduced to m_N and q.
- The most general form of J^{μ} is then:

$$J^{\nu} = e\bar{u}(p') [K_1(q^2)\gamma^{\nu} + (p^{\nu} + p'^{\nu})K_2(q^2) + (p^{\nu} - p'^{\nu})K_3(q^2) + i\sigma^{\nu\mu}(p_{\mu} - p'_{\mu})K_4(q^2) + i\sigma^{\nu\mu}(p_{\mu} + p'_{\mu})K_5(q^2)]u(p)e^{i(p-p')\cdot x}$$
(9)

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Gordon decomposition identity:

$$\bar{u}\gamma^{\mu}u = \frac{1}{2M}\bar{u}\left(p^{\mu} + p'^{\mu} + i\sigma^{\mu\nu}\left(p'_{\nu} - p_{\nu}\right)\right)u \qquad (10)$$

eliminating p + p' terms.

• Current conservation says $\partial_{\mu}J^{\mu} = q_{\mu}J^{\mu} = 0$, so terms that do not vanish must have K = 0

▶ For *K*₃ term:

$$q^{\mu}q_{\mu}K_{3} = q^{2}K_{3} \neq 0 \tag{11}$$

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so $K_3(q^2) = 0$.

• Rewriting J^{μ} in different terms:

$$J^{\nu} = e\bar{u}(p') \left[F_1(q^2)\gamma^{\nu} + \frac{\kappa}{2M} q_{\mu} \sigma^{\nu\mu} F_2(q^2) \right]$$
(12)

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• Transition amplitude from state $i \rightarrow f$:

$$T_{fi} = -i \int j^{\mu} \frac{1}{q^2} J_{\mu} d^4 x$$
 (13)

► Transition rate:

$$W_{fi} = \frac{|T_{fi}|^2}{T} \tag{14}$$

• relates to differenatial cross section $d\sigma$:

$$d\sigma = \frac{W_{fi}}{\Phi} dQ \tag{15}$$

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$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E^2 \sin^4 \frac{\theta}{2}} \frac{E'}{E} \left[\left(F_1^2 - \frac{\kappa^2 q^2}{4M^2} F_2^2 \right) \cos^2 \frac{\theta}{2} - \frac{q^2}{2M^2} (F_1 + \kappa F_2)^2 \sin^2 \frac{\theta}{2} \right]$$
(16)

References:

- My thesis
- Halzen and Martin, Quarks & Leptons: An Introductory Course in Modern Particle Physics, Wiley, 1984

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