



¹ Massachusettes Institute of Technology, Cambridge, MA 02139, USA





Predictions of Pentaquarks

In 1997, according to soliton model D. Diaknov et al. predicted a $SU(3)_F$ flavor antidecuplet of pentaquarks. In the antidecuplet, what attracts people most is a narrow (5 MeV), exotic state, $\Theta^+(1540)$, which has quark component $uudd\bar{s}$. Following observations have shown the quite high possibility of its existance. However, in our experiment, the interests were focused on Θ^+ 's partners: $N^0(1650)$ and $\Sigma^0(1760)$, and they can be used to verify the existance of such antidecuplet. Also, if the Θ^+ has non-zero isospin we're anticipating to discover it's isospin partners, so here we also selected a strong interaction channel for the Θ^{++} . The pentaquarks we searched in the experiment are circled by blue dots.

Particle Identification

AEROGEL Detectors

In HRS-left we used 3 Cerenkov detectors to identify particles. Two of them are amplitude detectors using AEROGEL with index 1.015 and 1.055. So the thresholds of Cerenkov radiation in these detectors are 2.84(0.803) GeV/c for kaons(pions) and 2.8(0.415) GeV/c for protons(pions) repectively.





The third detector measuring the angle of Cerenkov radiation is called RICH. It has a very good angle resolution of 5 mrad(σ). Accompanied by a radiator of index n = 1.3, it allows us to identify kaons from pions clearly up to 3 GeV/c momentum.





With the help of RICH detector, Aerogel 1, Aerogel 2, and other conventional Hall A PID detectors: Pion Rejector and Gas Cherenkov, we can clearly identify the kaon peak in the coincidence time spectrum.

the HRS-right at a central angle of 6° with another septum magnet. The polarity of HRS-left was positive for the Σ_{10}^0 and N_{10}^0 , and negative for the Θ^{++} .





Cross-Section Spectra

According to accumulated luminosity and other efficiency factors, we combined the acceptance corrected spectra from different kinematic settings and normalized the scale into cross-section. In left hand side, from top to bottom there are missing mass spectra for Σ^0 , Θ^{++} , N^0 search. In the Σ^0 spectrum, the properties of a known big resonance $\Lambda^0(1520)$ have also been measured. Note that some spectra may have negtive values because of accidental background substraction.

Kinematic Parameters

 $E_{\gamma}~=~3~GeV$, $Q^2~\sim~0.1~(GeV/c)^2$, $heta_{\gamma K(\pi)}^{Lab(CM)}~\sim~2(6)^{\circ}$

Parameters of $\Lambda^0(1520)$

$$M_{\Lambda^{\circ}(1520)} = 1520.2 \pm 0.5 \; MeV$$

 $\Gamma_{\Lambda_{0}(1500)} = 16.6 + 1.3 MeV$

Acceptance Correction

Missing Mass at Forward Angle

 $M_X = \sqrt{(p_e^{\mu} + M - p_{e'}^{\mu} - p_{K(\pi)}^{\mu})^2} \approx C' - p_R^0 \delta_R - p_L^0 \delta_L$



For particular missing mass, it's a straight line in the momentum acceptance and $\Delta E'$ is proportional to the length in the 2-dimensional acceptance plot.

Momentum Acceptance Correction



$$\frac{d\sigma}{d\Omega_K}|_{Lab}(\gamma p \to \Lambda^0 K^+) \approx 350 \ nb$$

Cross-Section Upper Limits for Pentaquarks

Since there's **NO SIGNIFICENT NARROW PEAK** found in our searching range, the upper limits in Lab system for differential scross-sections of 5 MeV resonances are given: (**PRELIMINARY**)

 $\frac{d\sigma}{d\Omega_K}(\gamma p \to \Sigma^0 K^+) < 11 \ nb = 3.1\% \ \frac{d\sigma}{d\Omega}(\Lambda(1520))$

 $\frac{d\sigma}{d\Omega_K}(\gamma p \to \Theta^{++}K^-) < 3 \ nb = 0.9\% \ \frac{d\sigma}{d\Omega}(\Lambda(1520))$

$$\frac{ao}{d\Omega_K}(\gamma p \to N^0 K^+) < 8 \ nb = 2.3\% \ \frac{ao}{d\Omega}(\Lambda(1520))$$



With above knowledge, we are able to normalize the missing mass spectrum and calculate cross-section $\frac{d^3\sigma}{dE'd\Omega_{e'}d\Omega_{K(\pi)}}$. The plot shows the pion missing mass spectra after such correction: the triangle shape of accidental events has been removed and the resonance $\Delta(1232)$ got its proper shape.