Jefferson Lab
EXPLORING THE NATURE OF MATTER

January 2015

Latest Developments on Parton Distribution Functions

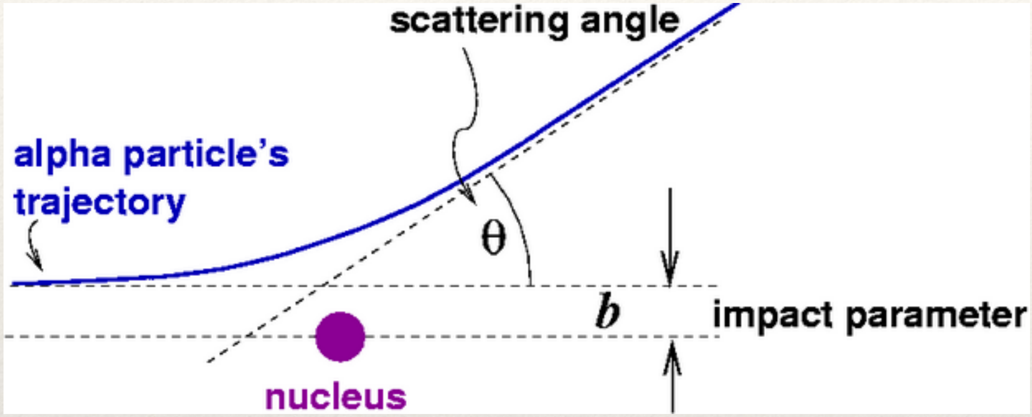
Voica Radescu
Physikalisches Institut Heidelberg

Abstract:

In the era of searching for new discoveries, it is crucial to achieve a higher level of precision in understanding the proton structure to allow for most unambiguous interpretations of the high energy, luminous data ahead. The knowledge of proton´s constituents mainly comes from the deep inelastic scattering HERA, fixed target, Tevatron and, increasingly precise LHC data. A guided tour of the road-map that marks the most recent measurements from past, present, and possible future experiments sensitive to the proton constituents will be presented here.

Early ideas …

- Rutherford's gold foil experiment 1909 (performed by Geiger and Marsden)



Geiger and Rutherford

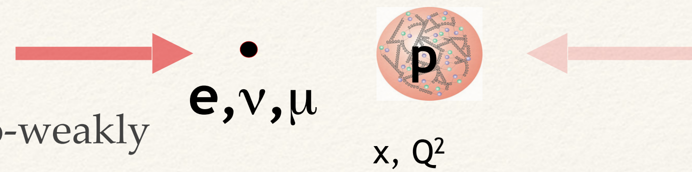
Rutherford's gold foil experiment set the scene for a century of ever-deeper and more precise resolution of the constituents of the atom, the nucleus and the nucleon.

➔ **Ideas for detecting quarks were formulated:**

To probe the interiors of target, point-like and easily produced particle needed to be used.

Probing the Proton Structure

- **Proton can be probed via elementary particles as:**
  - neutrinos (fixed target experiments) - interact only weakly
  - electrons (fixed target and collider experiments) - interact electro-weakly



- **Deep Inelastic Scattering (DIS)** as a tool to study the substructure of nucleon
  - scattering of a lepton off the quarks within the proton resulting into a hadronic shower and a lepton

- **Kinematic Lorenz Invariant Variables:**
  - virtuality of exchanged boson

$$Q^2 = -q^2 = -(k - k')^2$$

  - proton momentum fraction of the scattered quark (Bjorken scaling variable)

$$x = \frac{Q^2}{2p \cdot q}$$

  - inelasticity parameter:

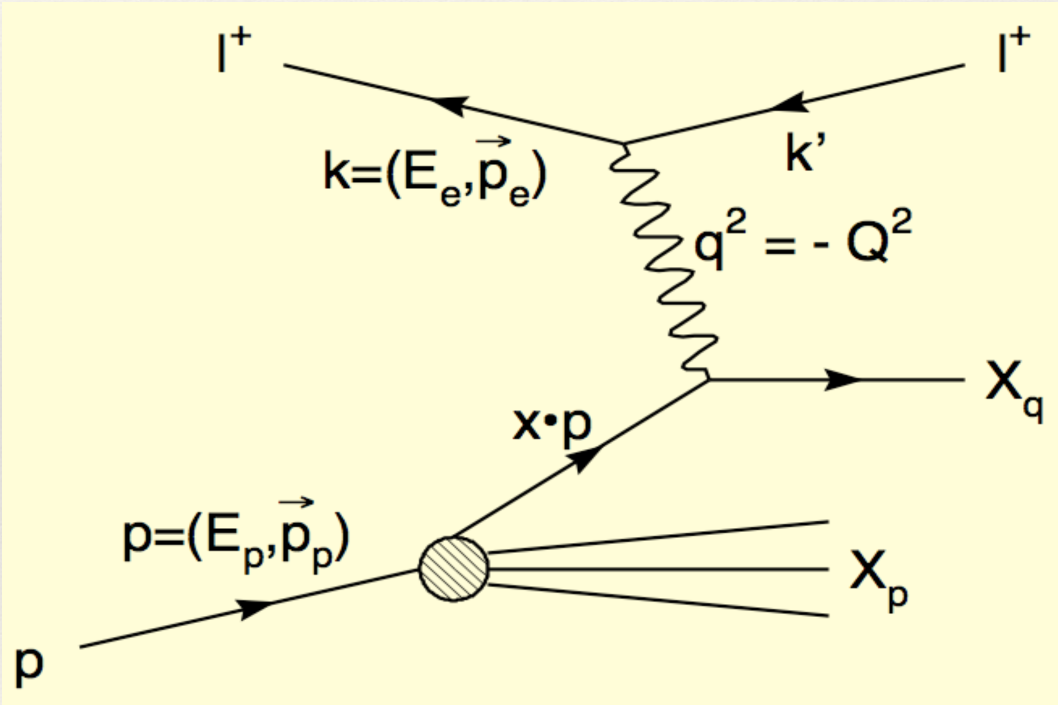
$$y = \frac{p \cdot q}{p \cdot k}$$

  - invariant centre of mass energy:

$$s = (k+p)^2 = \frac{Q^2}{xy}$$

  - Invariant centre of mass energy of the virtual boson-proton system

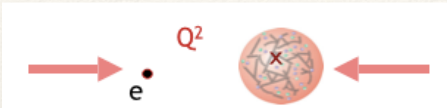
$$W^2 = (P+q)^2 = m_p^2 - Q^2 + 2P \cdot q = ys - Q^2 + m_p^2(1-y).$$



Cross Sections, Structure Functions, PDFs

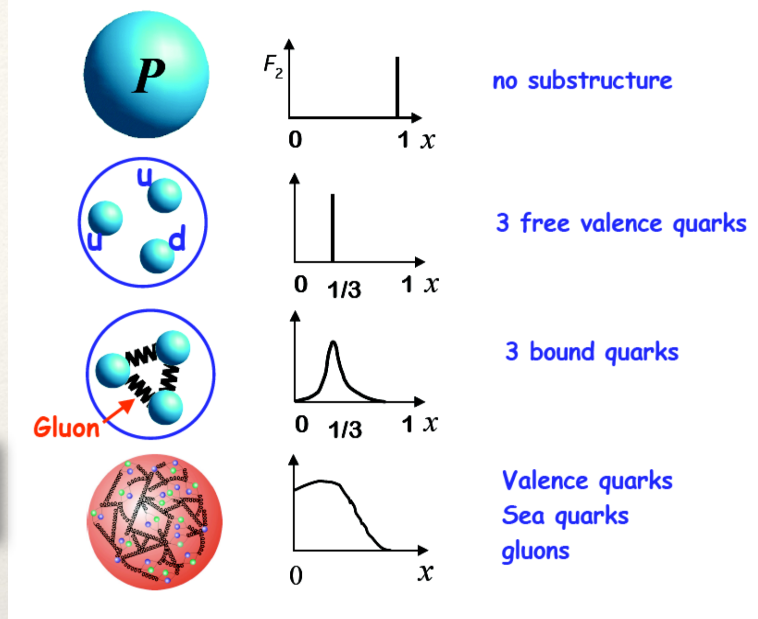
- The proton is a dynamical object and its structure depends on the resolution (Q^2) of the observation:
  - higher the value of Q^2 more detail we examine

- General double differential cross-section for eN scattering:



$$\sigma_r(x,Q^2) = \frac{d^2\sigma(e^\pm p)}{dxdQ^2}\frac{Q^4 x}{2\pi\alpha^2 Y_+} = F_2(x,Q^2) - \frac{y^2}{Y_+}F_L(x,Q^2) \mp \frac{Y_-}{Y_+}xF_3(x,Q^2)$$

$$Y_\pm = 1 \pm (1-y^2)$$



  - F_2, F_L, xF_3 are structure functions (hadronic part) which are related to the momentum distributions of quarks within the nucleon:

Parton Distribution Functions (PDFs):

  - valence quarks: carriers of proton charge
  - sea quarks and gluons: evolved by complex dynamics

At Leading Order (LO):

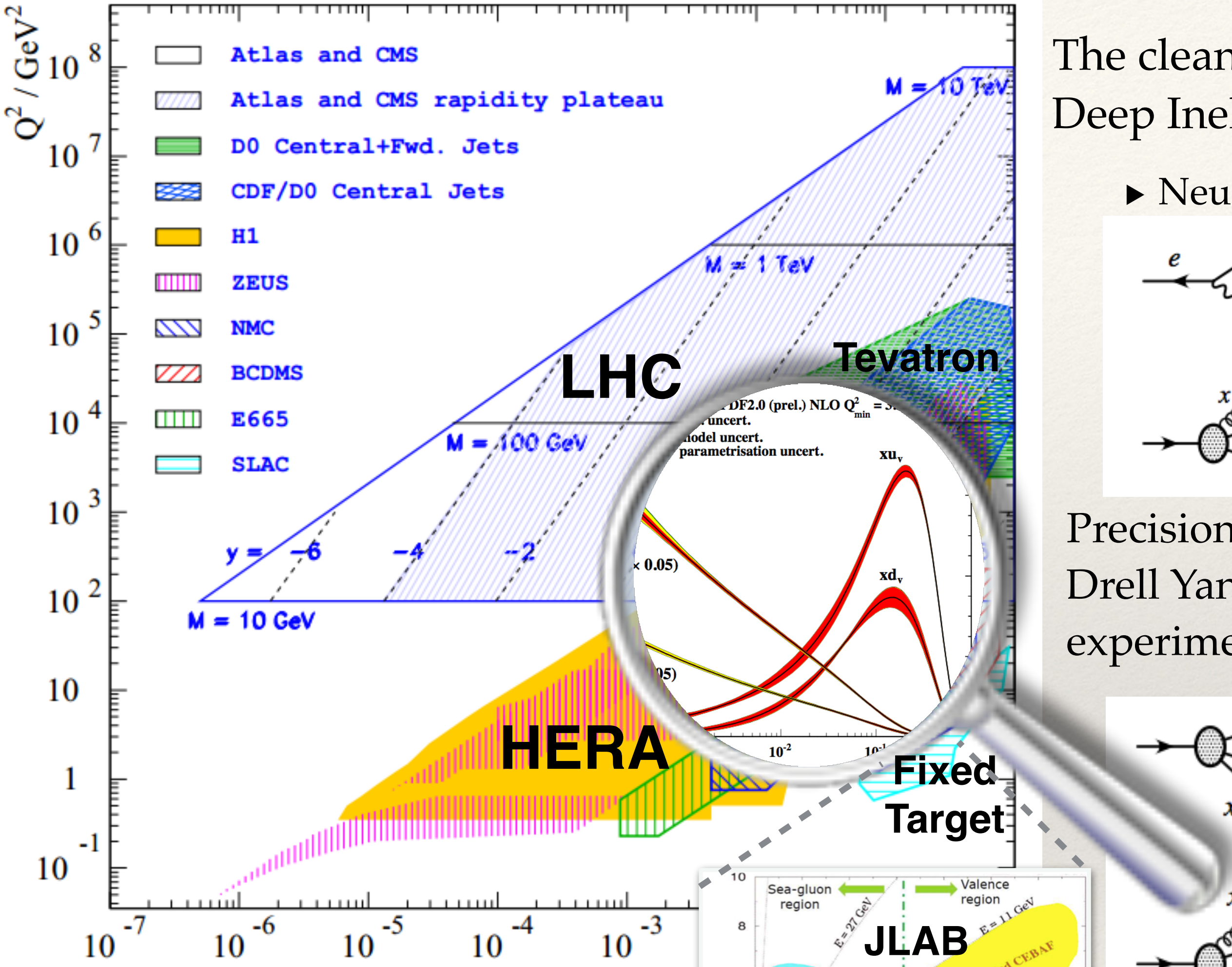
$$F_2 = x\sum e_q^2\left(q(x)+\bar{q}(x)\right)$$

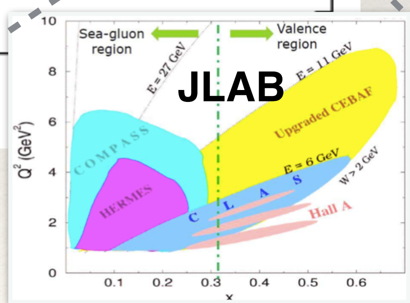
$$xF_3 = x\sum 2e_q a_q\left(q(x)-\bar{q}(x)\right)$$

- F_2 dominates sensitive to all quarks
- xF_3 sensitive to valence quarks
- F_L sensitive to gluons

- Can extract the structure functions experimentally by looking at the x,y,Q^2 dependence of the double differential cross-section

Proton Structure Measurements





$Q^2 = -q^2 = -(k-k')^2$ Photon virtuality

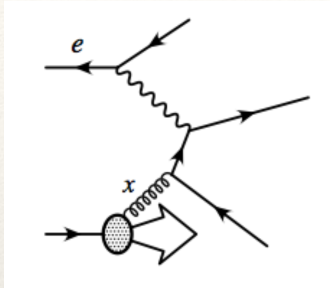
$x = \frac{Q^2}{2p \cdot q}$ Bjorken variable

$y = \frac{p \cdot q}{p \cdot k}$ Inelasticity

X

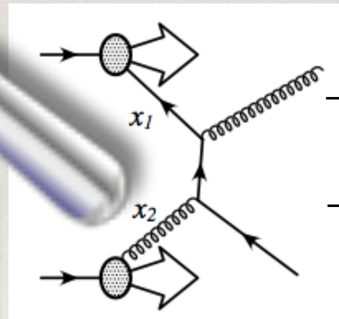
The cleanest way to probe Proton Structure is via Deep Inelastic Scattering [DIS]:

- Neutrinos, muons, electrons



—> probes linear combination of quarks

Precision of PDFs can be complemented by the Drell Yan [DY] processes at the collider experiments - [Tevatron and LHC]



—> can provide flavour separation and more insight into gluons

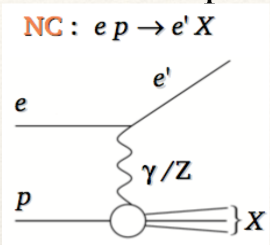
—> probes bilinear combination of quarks

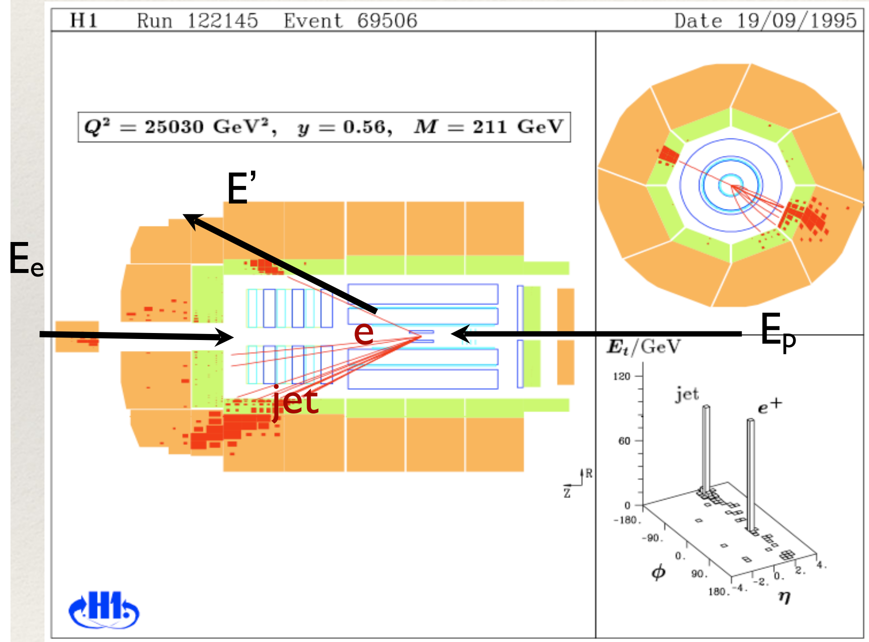
Different data constrain different parton combinations at different x, evolution with the scale is predicted by pQCD:

HERA ep collider (1992-2007) @ DESY

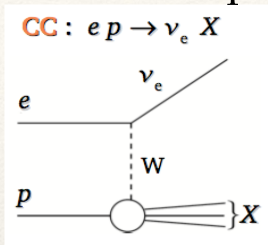
- H1 and ZEUS experiments at HERA collected ~1/fb of data
  - Ep=460/575/820/920 GeV and Ee=27.5 GeV

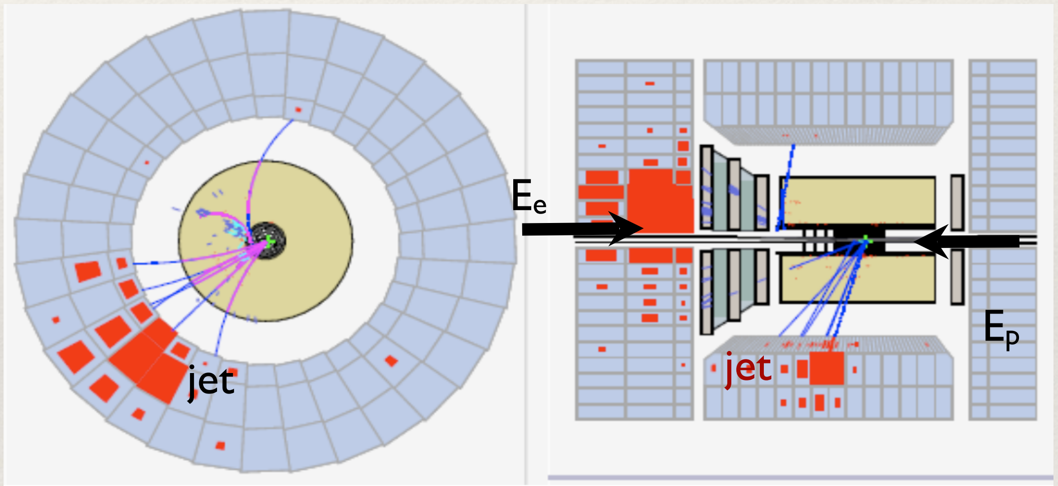
- **Neutral Current** event sample in H1 detector





- **Charged Current** event sample in ZEUS detector





Determination of the Event Kinematics:

- using lepton information (E_e',θ_e)
- **using hadronic final state particles**
- using both lepton and hadronic final state variables

$$s = 4E_eE_p$$
$$Q^2 = E_eE'(1+\cos\theta_e)$$
$$y = 1 - \frac{E'}{E_e}\frac{1}{2}(1-\cos\theta_e)$$
$$x = \frac{Q^2}{sy}$$

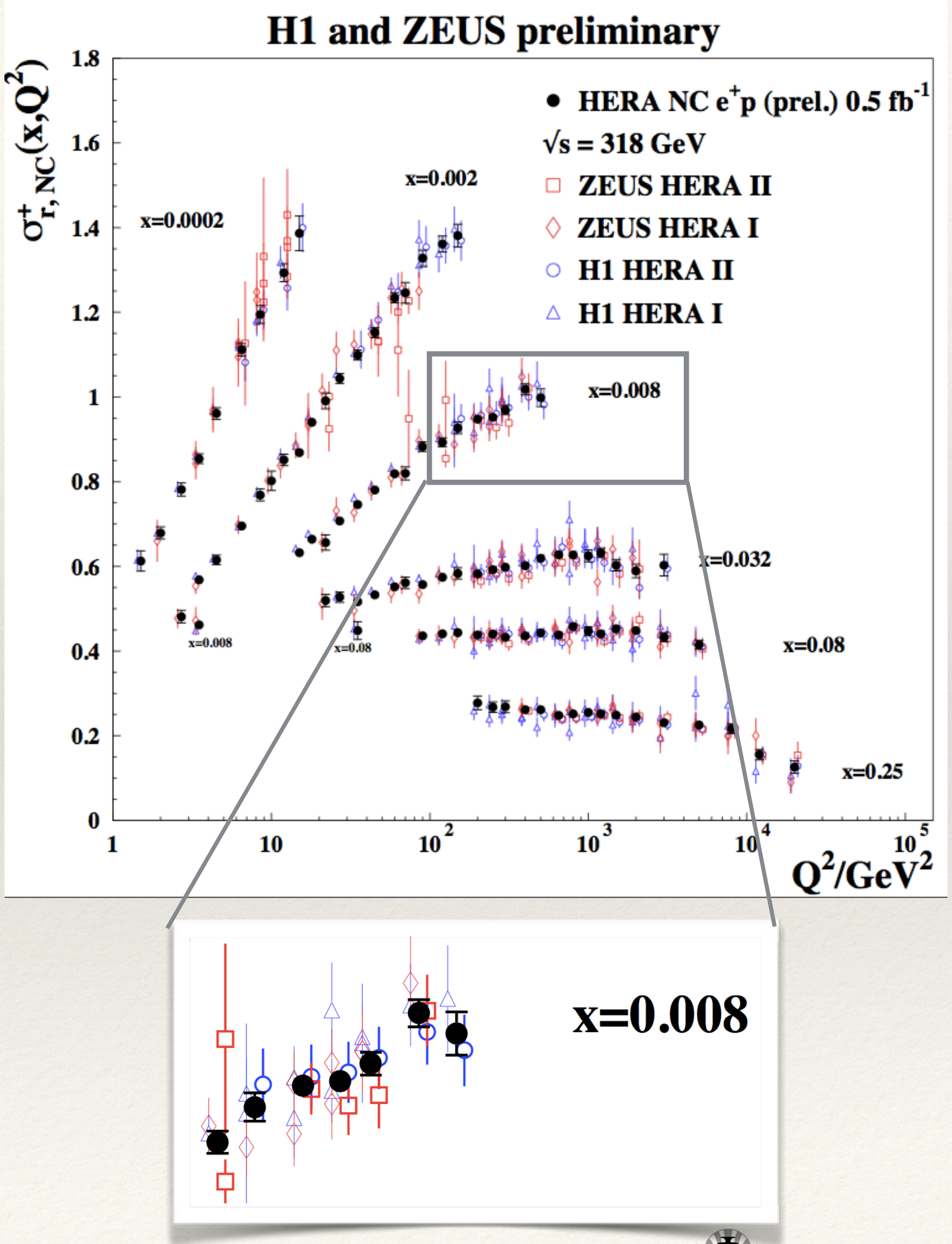
Combination of the H1 and ZEUS Measurements

[JHEP01 (2010) 109], H1 prelim-14-041, ZEUS-prel-14-005

- Ultimate precision is obtained by combining the H1 and ZEUS measurements
- The combination procedure is performed before QCD analysis using χ^2 minimisation

➢ Improvement on Statistical precision:
  - ❖ H1 and ZEUS collected similar amounts of physics data.

➢ Improvement of Systematic precision:
  - ❖ H1 and ZEUS are different detectors and use different analysis techniques;
  - ❖ The H1 and ZEUS cross sections have different sensitivities to similar sources of correlated systematic uncertainty.

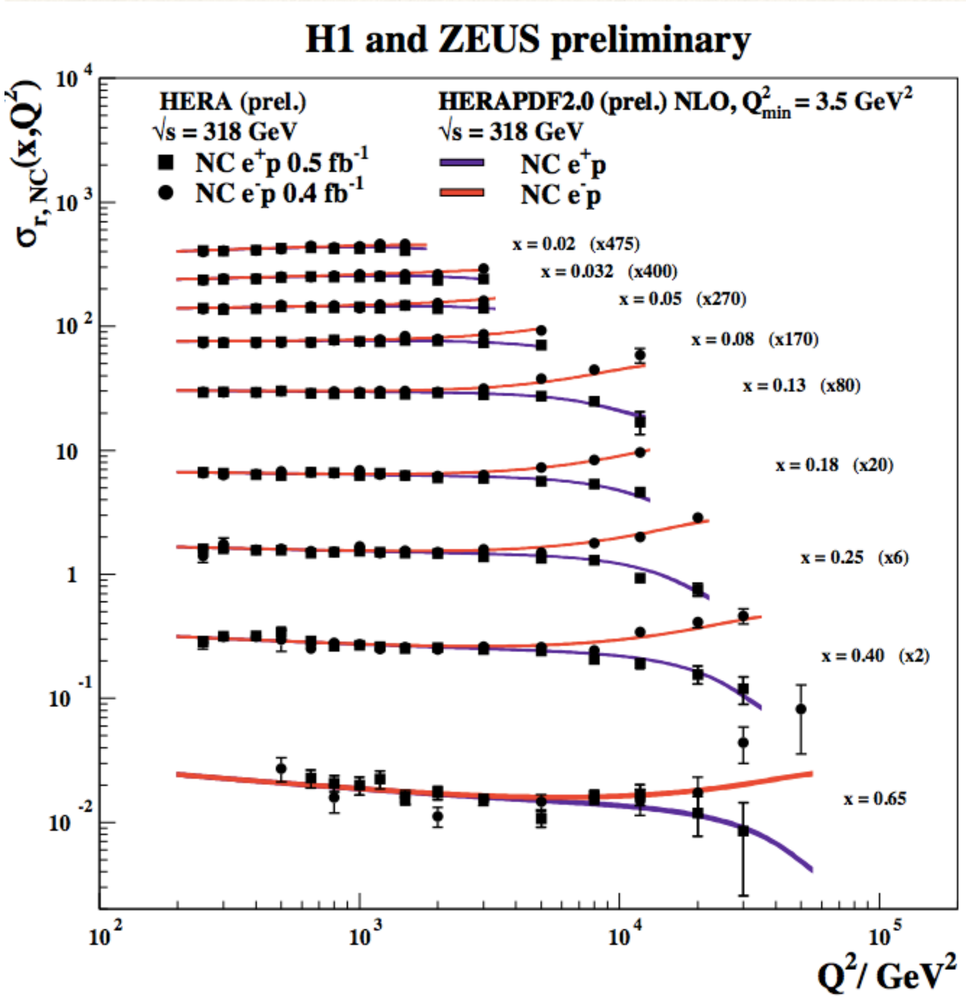


—> total uncertainty < 1.3% for Q2 up to 400 GeV2

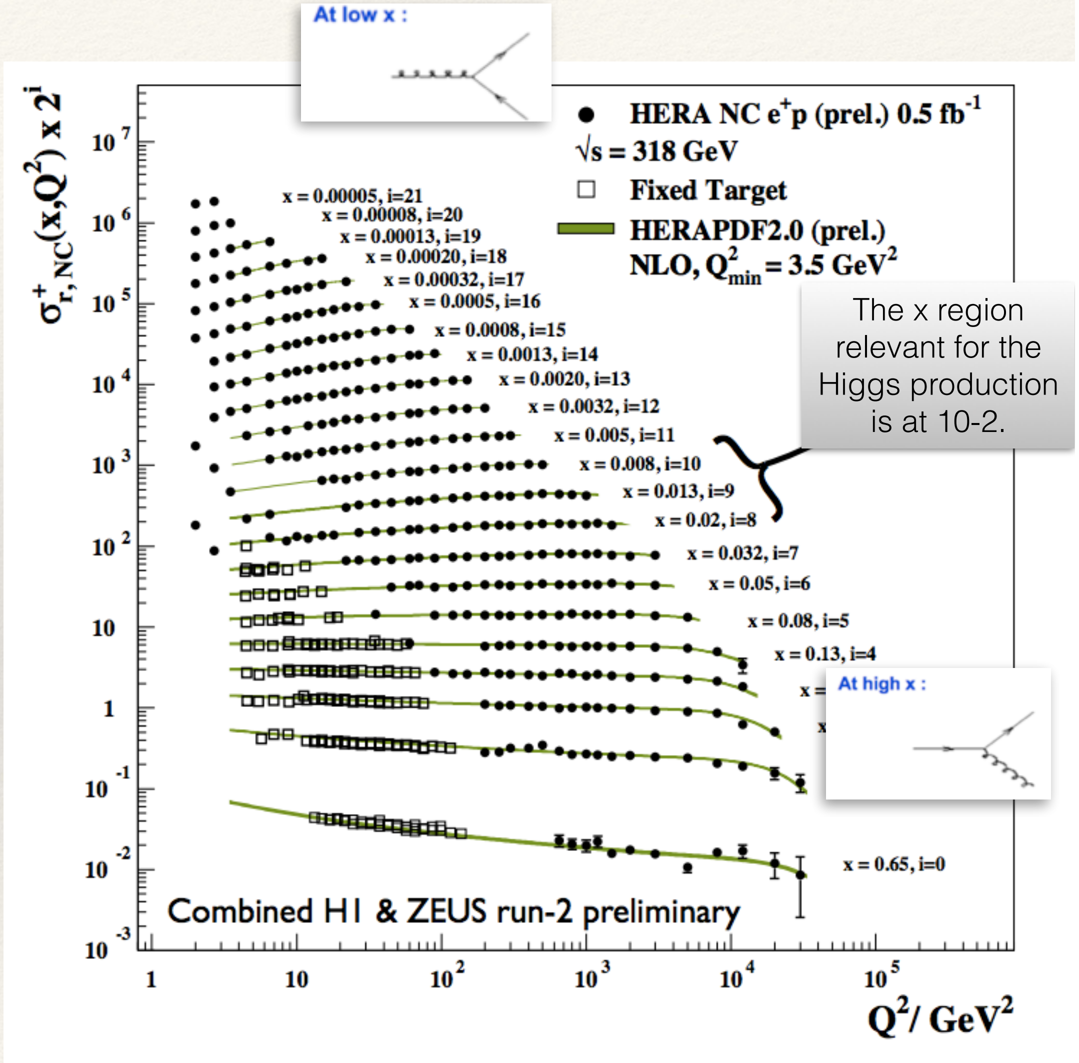
- 41 data sets: 2927 data points are combined to1307 averaged measurements with 162 sources of correlated systematic uncertainties.

QCD scaling and EW effects

- EW effects clearly seen at high Q2:



QCD scaling violations nicely seen:



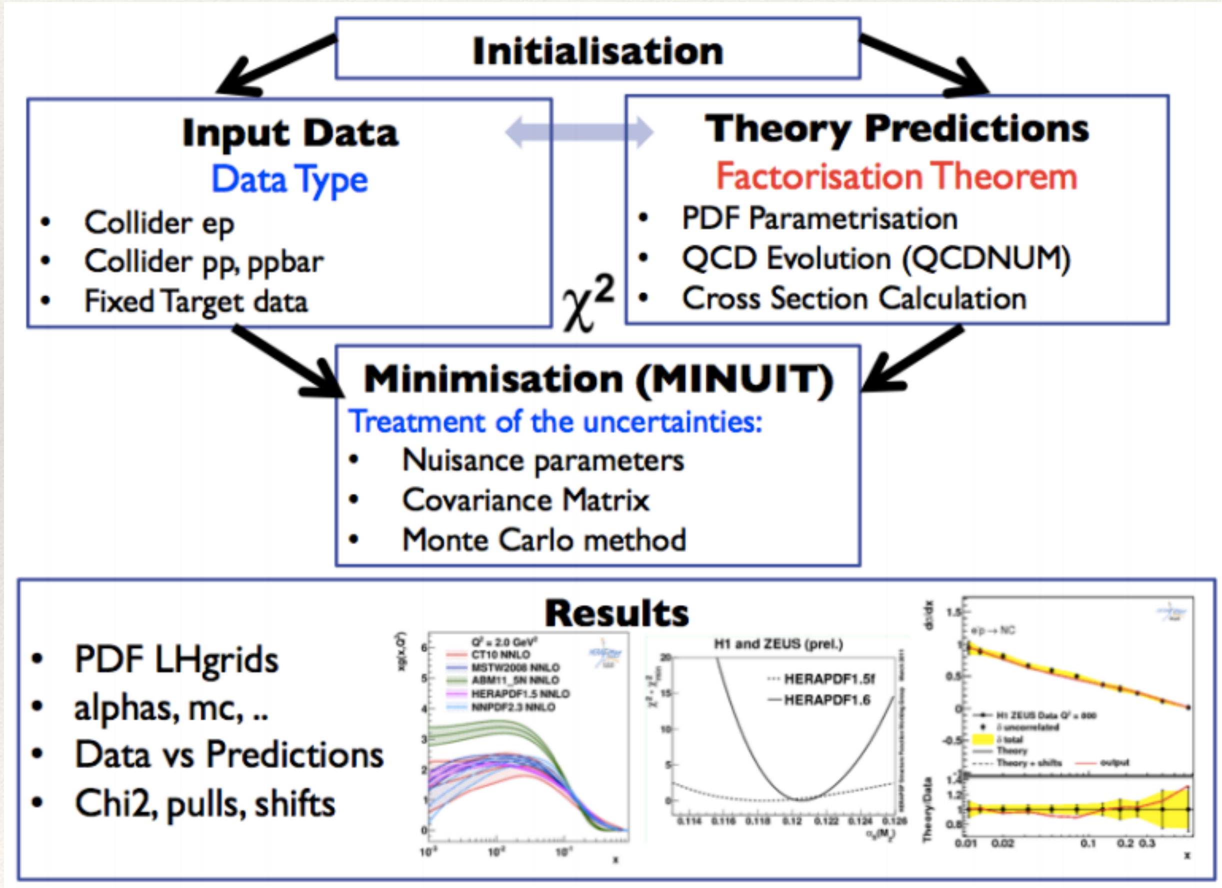
Extraction of PDFs through QCD fits

- Extraction of PDFs relies on the factorisation: $\sigma = \hat{\sigma} \otimes \text{PDF}$

- Typical measurements sensitive to PDFs are precise, with statistical uncertainties < 10%, so they follow normal distribution —> use of χ^2 minimisation for PDF extraction.

Main Steps:

- Parametrise PDFs at a starting scale
- Evolve PDFs to the scale corresponding to data point
- Calculate the cross section
- Compare with data via χ^2
- Minimise χ^2 with respect to PDF parameters



Extraction of PDFs through QCD fits

- Extraction of PDFs relies on the factorisation: $\sigma = \hat{\sigma} \otimes \text{PDF}$

- Typical measurements sensitive to PDFs are precise, with statistical uncertainties < 10%, so they follow normal distribution —> use of χ^2 minimisation for PDF extraction.

Main Steps:

- **Parametrise PDFs at a starting scale**

HERAPDF style has standard polynomial form: $xf_j(x) = A_j x^{B_j}(1-x)^{C_j} P_j(x)$

where $P_j(x) = (1 + \epsilon\sqrt{x} + Dx + Ex^2)$ for HERA PDF style, or

$P_j(x) = e^{a_3 x}(1 + e^{a_4} x + e^{a_5} x^2)$ for CTEQ PDF style

A: overall normalisation
B: small x behavior
C: x→ 1 shape

Bi Log-Normal Distribution style: $xf_j(x) = a_j x^{p_j - b_j \log(x)}(1-x)^{q_j - d_j \log(1-x)}$

Chebyshev is a flexible parametrisation (polynomials with argument log(x)) which can be employed for the gluon and sea distributions
(see A. Glazov, S. Moch, and V. Radescu, Phys. Lett. B 1149 695, 238 (2011), [arXiv:1009.6170])

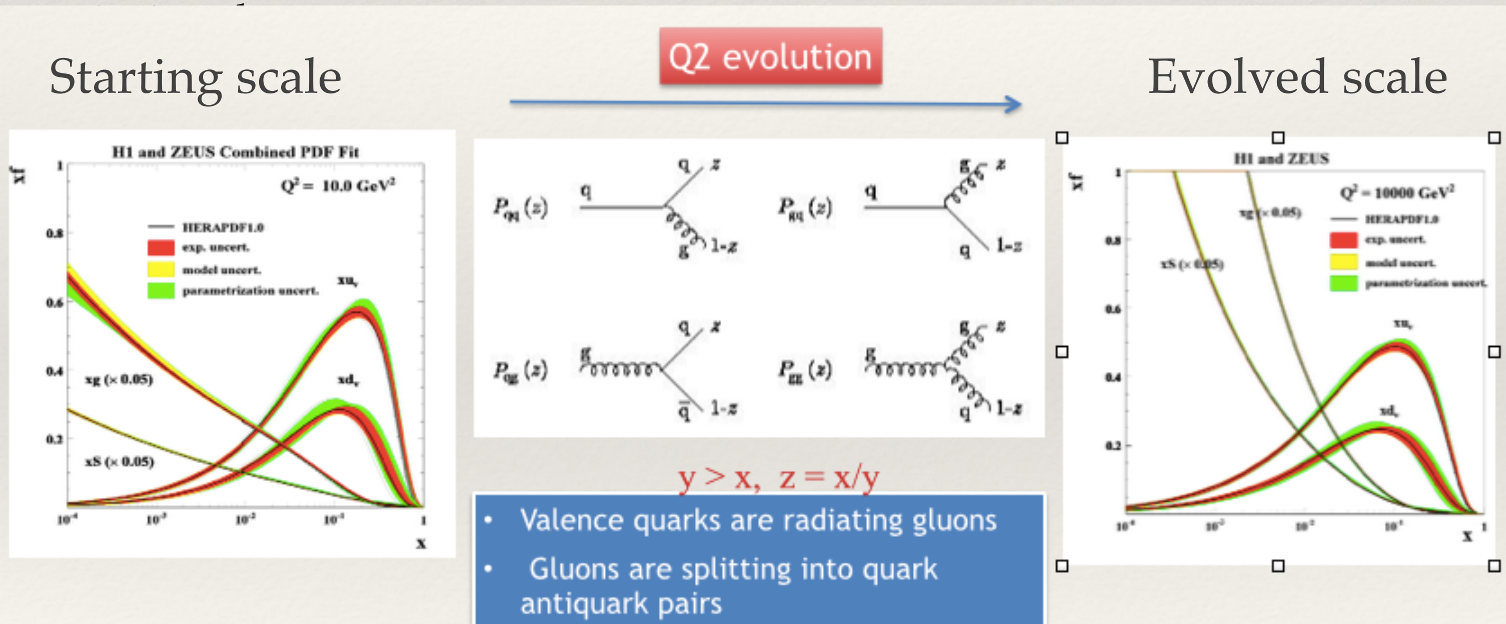
Extraction of PDFs through QCD fits

- Extraction of PDFs relies on the factorisation: $\sigma = \hat{\sigma} \otimes \text{PDF}$

- Typical measurements sensitive to PDFs are precise, with statistical uncertainties < 10%, so they follow normal distribution —> use of χ^2 minimisation for PDF extraction.

Main Steps:

- Parametrise PDFs at a starting scale
- **Evolve PDFs to the scale corresponding to data point**



Extraction of PDFs through QCD fits

- Extraction of PDFs relies on the factorisation: $\sigma = \hat{\sigma} \otimes \text{PDF}$

- Typical measurements sensitive to PDFs are precise, with statistical uncertainties < 10%, so they follow normal distribution —> use of χ^2 minimisation for PDF extraction.

www.herafitter.org

Main Steps:

- Parametrise PDFs at a starting scale
- Evolve PDFs to the scale corresponding to data point
- **Calculate the cross section**
- Compare with data via χ^2
- Minimise χ^2 with respect to PDF parameters

DY integration code:

Simple LO cross section formulae: DY NC: $pp \to Z/\gamma \to e^+e^-$

$$\frac{d\sigma^2_\gamma}{dMdy d\cos\theta^*} = N_c C^2_{q\bar{q}} \frac{8\alpha^2}{3M^3}\tau \times \sum_q e^2_q f_q(x_1, M) f_{\bar{q}}(x_2, M) F_{q\bar{q}}(1+\cos^2\theta^*, \cos\theta^*)$$

DY CC: $pp \to W^\pm \to e^\pm\nu$

$$\frac{d\sigma^3_{W\pm}}{dMdy d\cos\theta^*} = \frac{\pi\alpha^2}{48 s^4_W} M\tau \frac{(1-\cos\theta^*)^2}{(M^2 - M^2_W)^2 + \Gamma^2_W M^2_W} \times \sum_{qq'} V_{qq'} f_q(x_1, M) f_{q'}(x_2, M)$$

Theory:
necessary to have fast tools (APPLGRID, FASTNLO)

Extraction of PDFs through QCD fits

- Extraction of PDFs relies on the factorisation: $\sigma = \hat{\sigma} \otimes \text{PDF}$

- Typical measurements sensitive to PDFs are precise, with statistical uncertainties < 10%, so they follow normal distribution —> use of χ^2 minimisation for PDF extraction.

Main Steps:

www.herafitter.org

- Parametrise PDFs at a starting scale
- Evolve PDFs to the scale corresponding to data point
- Calculate the cross section
- **Compare with data via χ^2**
- **Minimise χ^2 with respect to PDF parameters**

Covariance Matrix Representation:

$$\chi^2(m) = \sum_{i,k} (m_i - \mu_i) C^{-1}_{ik} (m_k - \mu_k)$$

is C^{-1}_{ik}**(stat)**, C^{-1}_{ik}**(syst)** or C^{-1}_{ik}**(tot)**$=C^{-1}_{ik}$**(stat)**$+C^{-1}_{ik}$**(syst)**

Mixed form representation:

$$\chi^2_{\text{exp}}(\boldsymbol{m}, \boldsymbol{b}) = \sum_{ij} \left(m^i - \sum_l \Gamma^i_l(m^i) b_l - \mu^i \right) C^{-1}{}_{ij}(m^i, m^j) \left(m^j - \sum_l \Gamma^j_l(m^j) b_l - \mu^j \right) + \sum_l b_l^2$$

Nuisance parameter case:

data theory

$$\chi^2(m,b) = \sum_i \frac{\left[\mu_i - m_i\left(1 - \sum_j \gamma^i_j b_j\right)\right]^2}{\delta^2_{i,\text{unc}} m_i^2 + \delta^2_{i,\text{stat}}\, \mu_i m_i \left(1 - \sum_j \gamma^i_j b_j\right)} + \sum_j b_j^2 + \log \text{penalty}$$

$\delta_{i,\text{stat}}$ $\delta_{i,\text{uncor}}$ are relative statistical and uncor related systematic uncertainties of the measurement i

γ^i_j quantifies the sensitivity of the measurement to the correlated systematic source j

b_j nuisance parameter

→ for more details see e.g. HERAPDF1.9, JHEP 1001, 109 (2010)

Extraction of PDFs through QCD fits

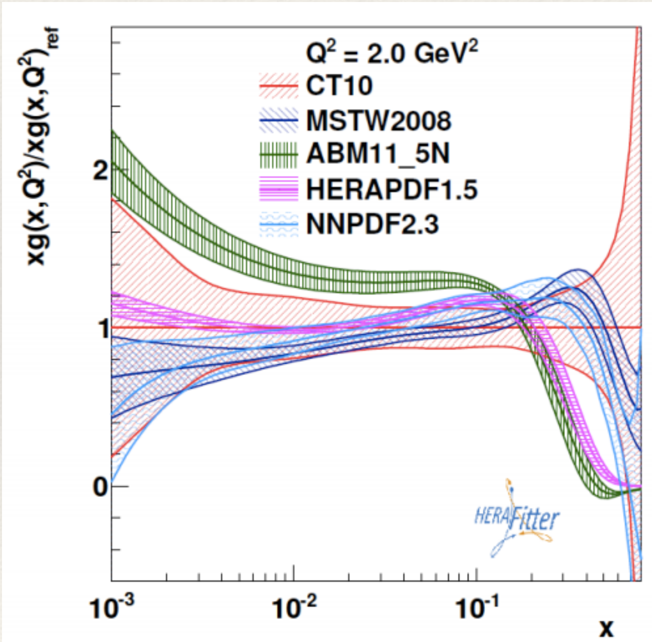
- Extraction of PDFs relies on the factorisation: $\sigma = \hat{\sigma} \otimes \text{PDF}$

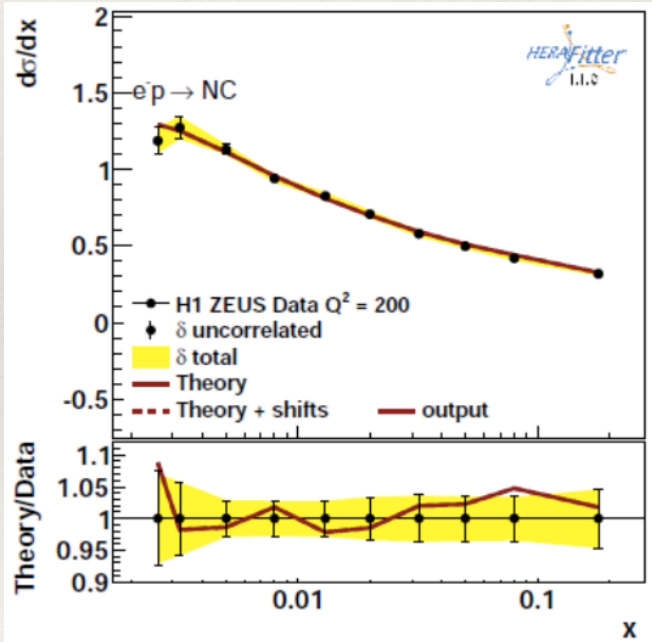
- Typical measurements sensitive to PDFs are precise, with statistical uncertainties < 10%, so they follow normal distribution —> use of χ^2 minimisation for PDF extraction.

Main Steps:

www.herafitter.org

- Parametrise PDFs at a starting scale
- Evolve PDFs to the scale corresponding to data point
- Calculate the cross section
- Compare with data via χ^2
- Minimise χ^2 with respect to PDF parameters
- **extract PDFs, pulls, etc..**





Modern understanding of PDFs

Different types of PDF uncertainties are considered:

- **Experimental:**
  - Hessian method used: MSTW, CT, ABM, JR, HERAPDF, CJ..
    - Consistent data sets → use $\Delta\chi^2$=1 (HERAPDF, ABM, JR ..) or larger tolerance (MSTW, CT)
  - Monte Carlo Method: replicas of data: NNPDF

- Method consists in preparing replicas of data sets allowing the central values of the cross sections to fluctuate within their systematic and statistical uncertainties taking into account all point to point correlations [A.Glazov and VR, HERA-LHC proceedings, arXiV:0901.2504, page 41-42]
  - Shift central values randomly within the uncorrelated errors assuming Gauss distribution of the errors:

$$\sigma_i = \sigma_i(1 + \delta_i^{uncorr} RAND_i)$$

  - Shift central values with the same probability of the corresponding correlated systematic shift assuming Gauss distribution of the errors:

$$\sigma_i = \sigma_i(1 + \delta_i^{uncorr} RAND_i + \sum_j^{N_{sys}} \delta_{ij}^{corr} RAND_j)$$

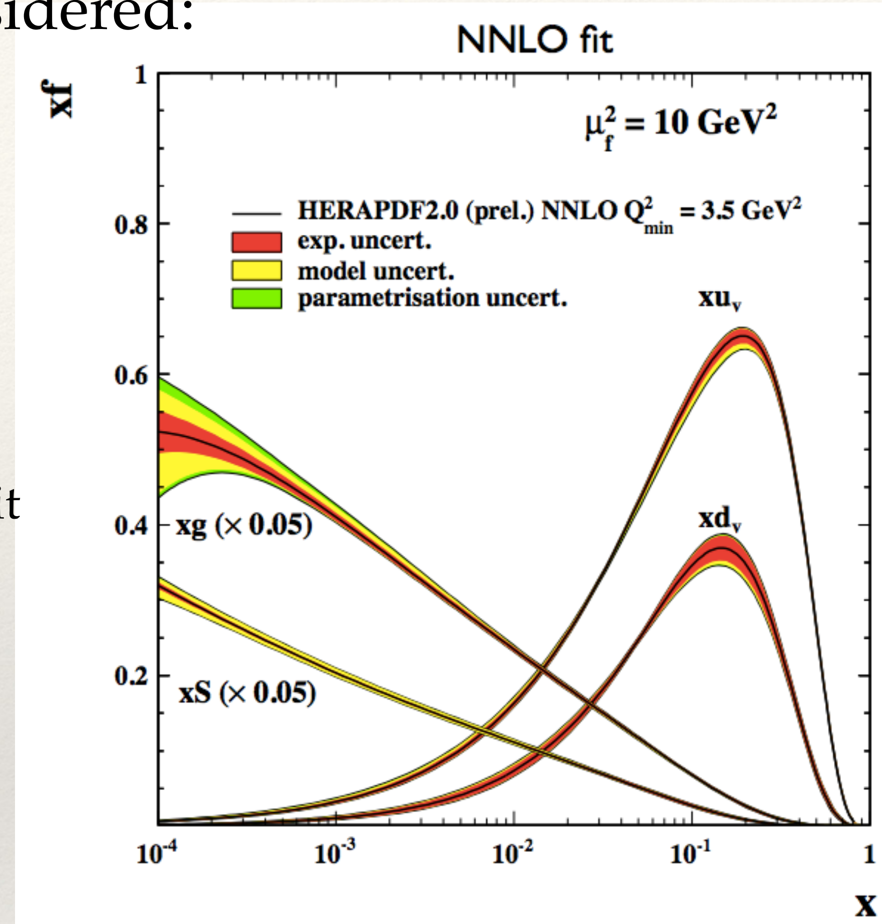
- Preparation of the data is repeated for N times (N>100)
  - For each MC replica, NLO QCD fit is performed to extract the N PDF sets
- Errors on the PDFs are estimated from the RMS of the spread of the N curves corresponding to the N individual extracted PDFs

Modern understanding of PDFs

Different types of PDF uncertainties are considered:

- **Experimental:**
  - Hessian method used: MSTW, CT, …
  - Consistent data sets → use $\Delta\chi^2$=1
  - Monte Carlo Method: replicas of data: NNPDF
- **Model:**
  - variations of all assumed input parameters in the fit

Variation	Standard Value	Lower Limit	Upper Limit
f_s	0.4	0.3	0.5
M_c^{opt}(NLO) [GeV]	1.47	1.41	1.53
M_c^{opt} (NNLO) [GeV]	1.44	1.38	1.50
M_b [GeV]	4.75	4.5	5.0
Q^2_{min} [GeV^2]	10.0	7.5	12.5
Q^2_{min} [GeV^2]	3.5	2.5	5.0
Q^2_0 [GeV^2]	1.9	1.6	2.2



- **Parametrisation:** only HERAPDF includes this as additional unc.
  - An envelope formed from PDF fits using variants of parametrisation form (extra parameter added)
  - NNPDF use neural network approach based on data driven regularisation

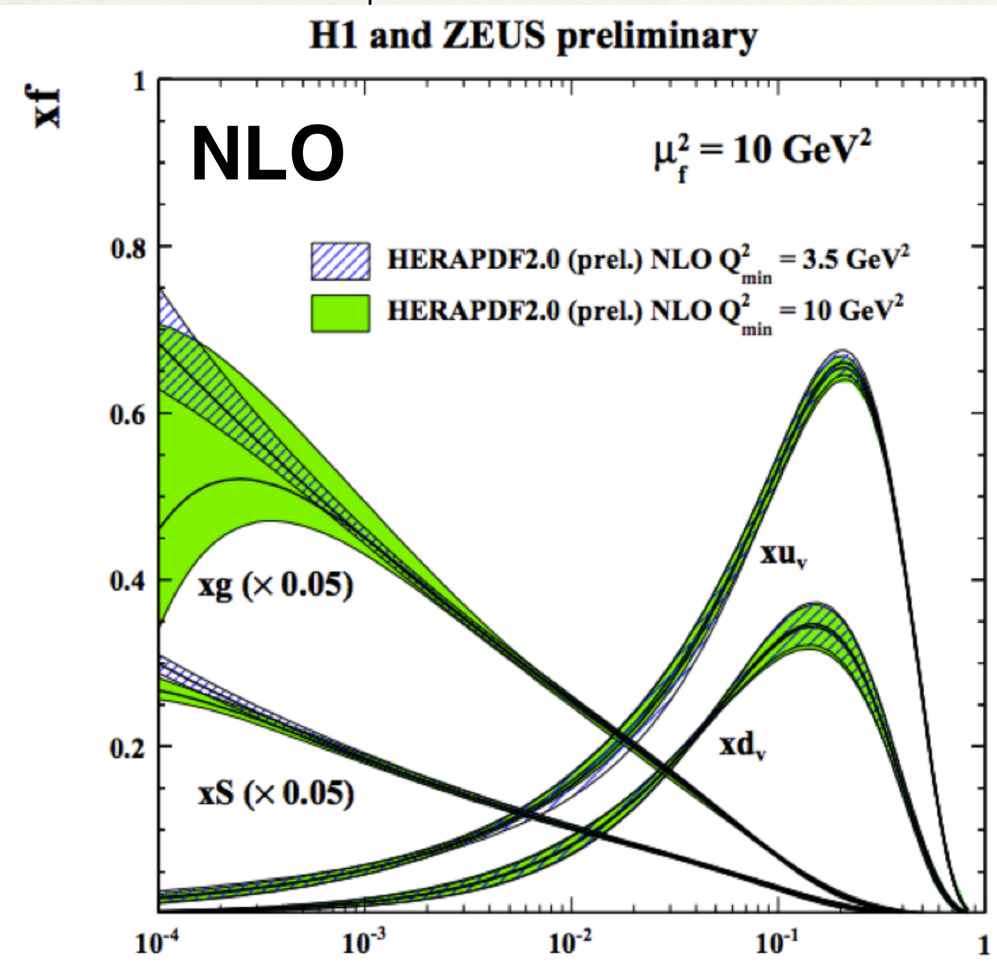
HERAPDF2.0 - fit to HERA data

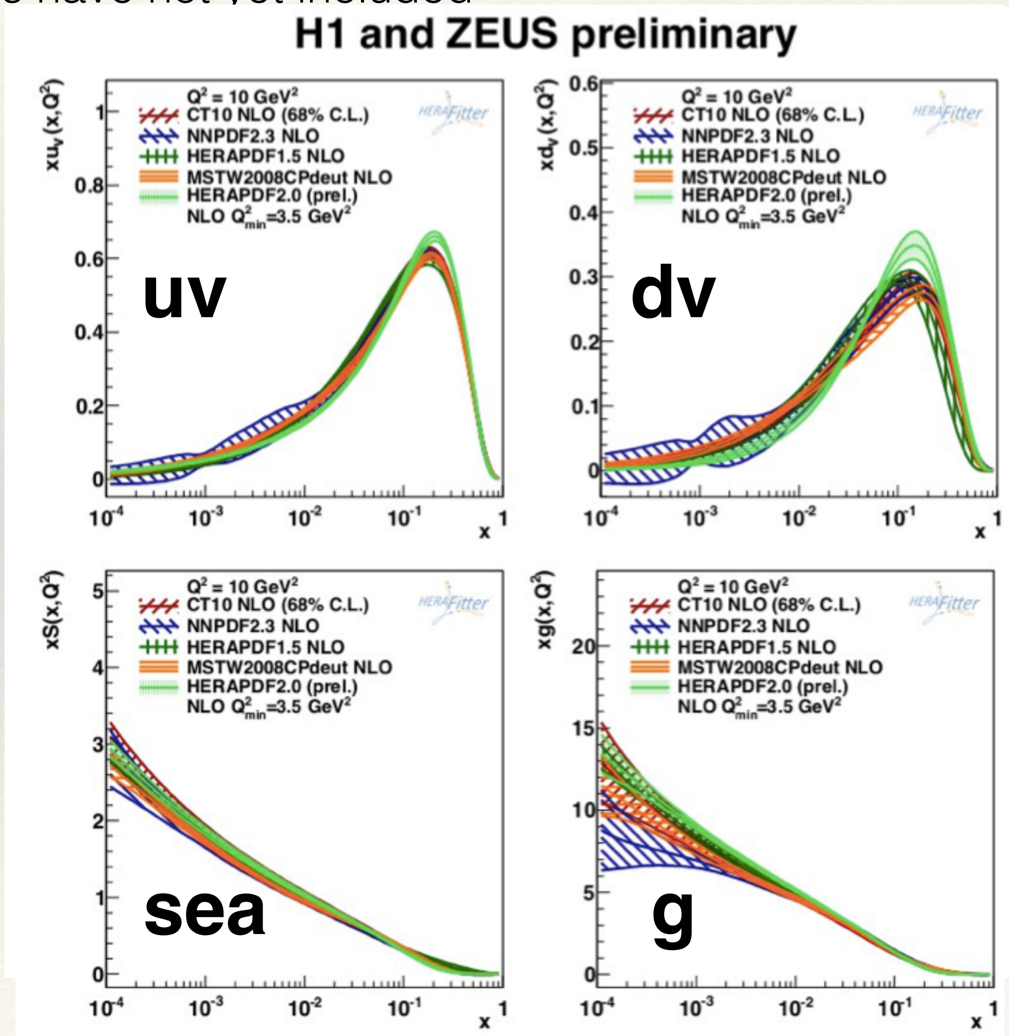
- PDFs with Q2 cut min @ 3.5 GeV2 and @10 GeV2 are shown

  uncertainties are larger for Q2cut=10 GeV2 (more data is cut away) and impact mostly gluon PDF

—> dval extracted from ep data only:

- new CC data provides new info that other PDFs have not yet included





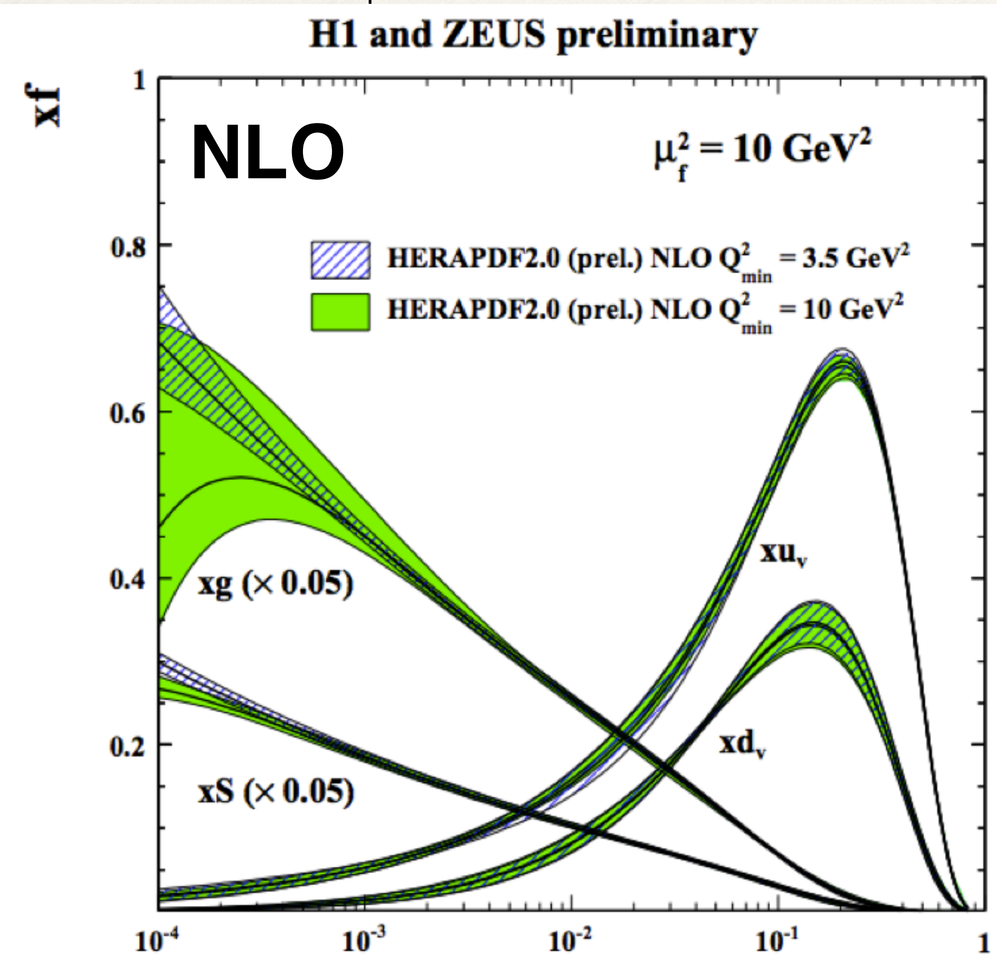
HERAPDF2.0 - fit to HERA data

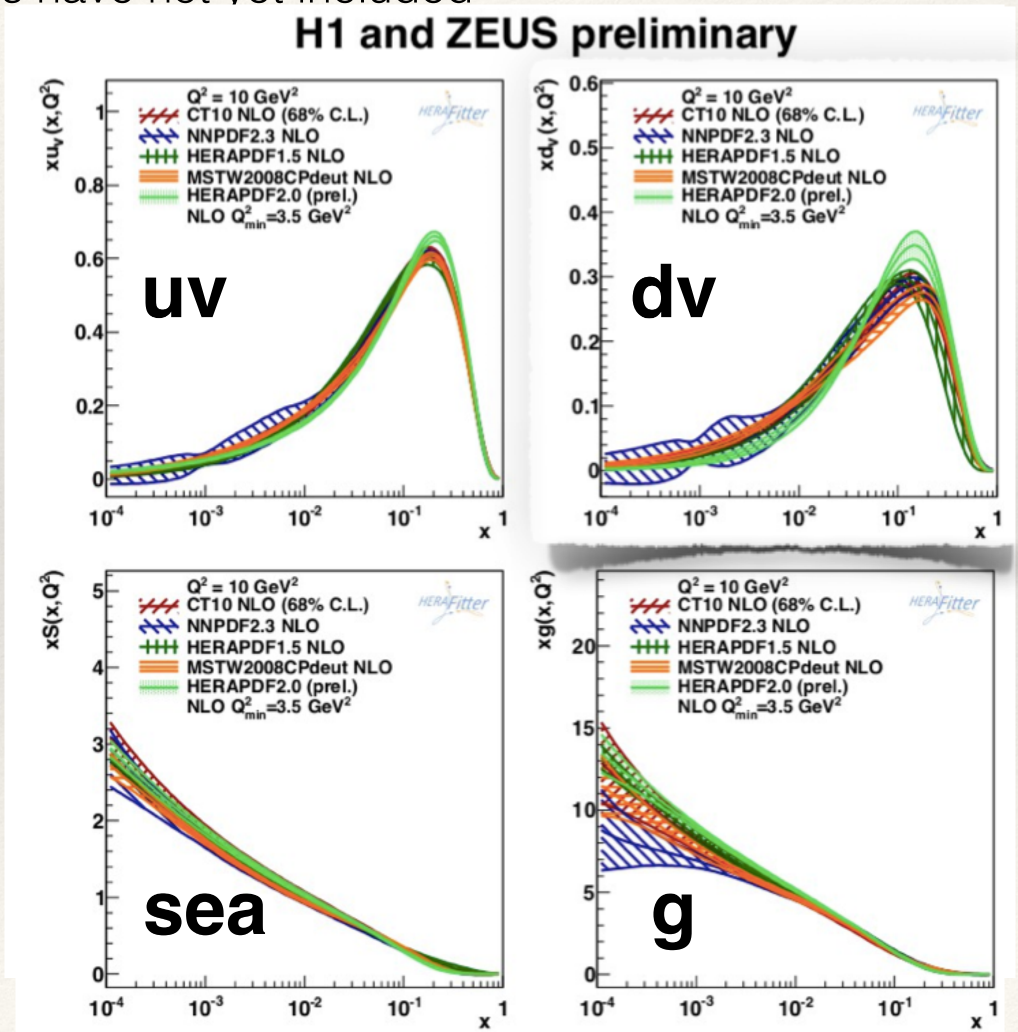
- PDFs with Q2 cut min @ 3.5 GeV2 and @10 GeV2 are shown

  uncertainties are larger for Q2cut=10 GeV2 (more data is cut away) and impact mostly gluon PDF

—> dval extracted from ep data only:

- new CC data provides new info that other PDFs have not yet included





PDF sets in use at the LHC

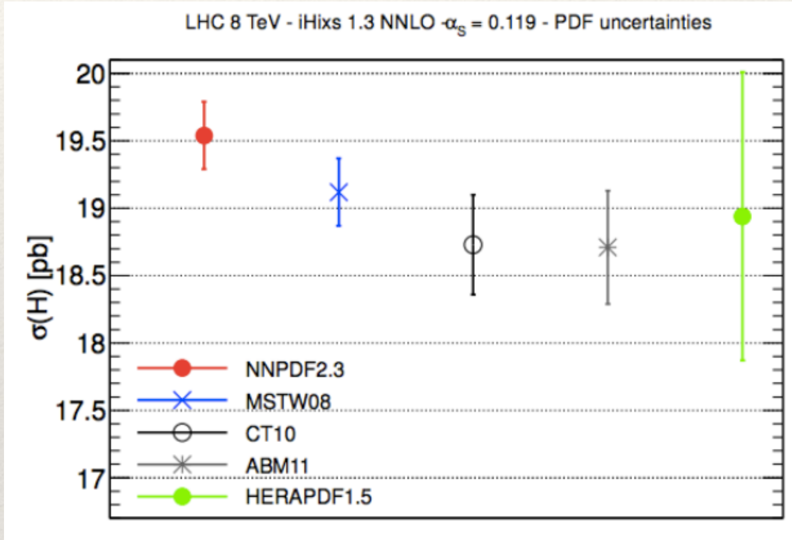
- Data: targeted measurements, detailed information of sources of systematic uncertainties, addressing the importance of correlation information
- Theory: state of the art methods, advancement in computational powers that allowed fo higher order calculations to be available

Global proton PDF groups used at the LHC

Note:

here are also JLAB PDF sets (CJ) to address the high x phenomenology

August 2014	CT10(w)	MSTW2008	NNPDF2.3	ABM12	HERAPDF15
Fixed Target DIS	✔	✔	✔	✔	✘
HERA	✔	✔	✔	✔	✔
Fixed Target DY	✔	✔	✔	✔	✘
Tevatron W,Z	✔	✔	✔	✘	✘
Tevatron jets	✔	✔	✔	✘	✘
LHC data	✘	✘	✔	✔	✘
Stat. treatment	Hessian $\Delta\chi^2 = 100$	Hessian $\Delta\chi^2$ dynamical	Monte Carlo	Hessian $\Delta\chi^2=1$	Hessian $\Delta\chi^2=1$
Parametrization	Pol. (26 pars)	Pol. (20 pars)	NN (259 pars)	Pol. (14 pars)	Pol. (14 pars)
HQ scheme	ACOT-χ	TR'	FONLL	FFN	TR'
α_S	Varied	Fitted+varied	Varied	Fitted	Varied



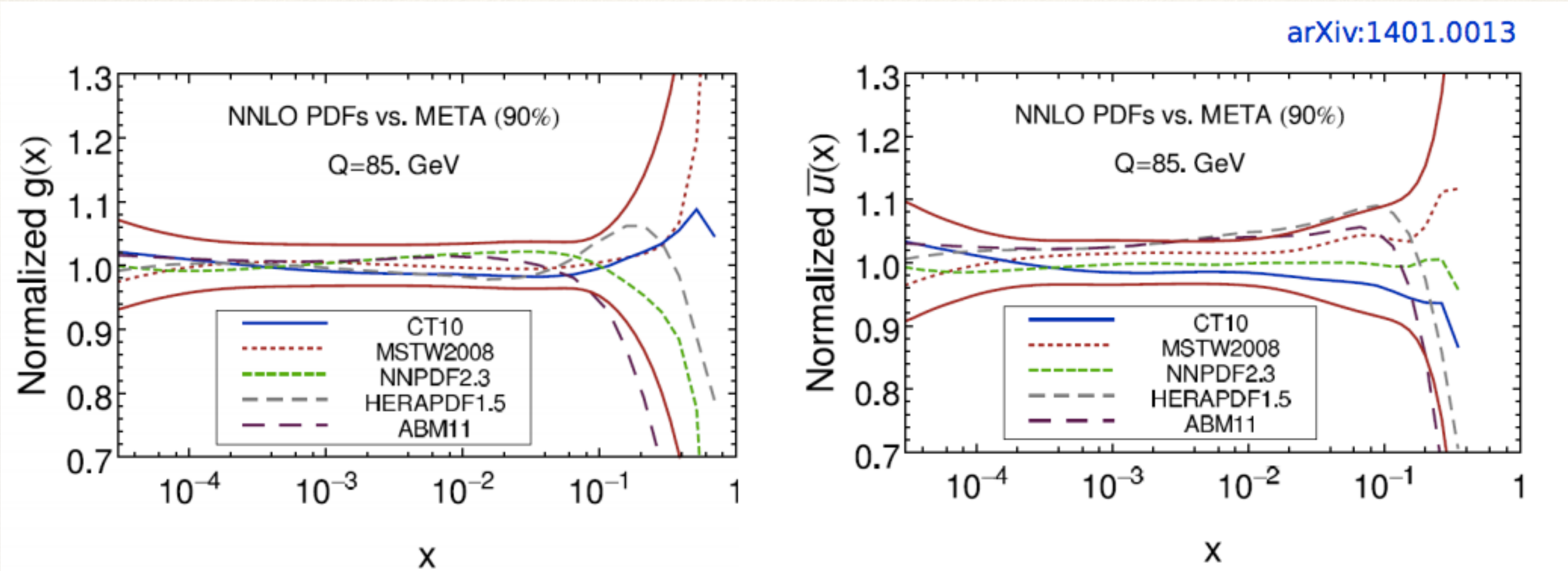
Dedicated studies to address this difference, PDF4LHC, http://arxiv.org/pdf/1405.1067.pdf

The analyses differ in many areas:

- different treatment of heavy quarks
- inclusion of various data sets and account for possible tensions
- different alphas assumption

METAPDF proposal

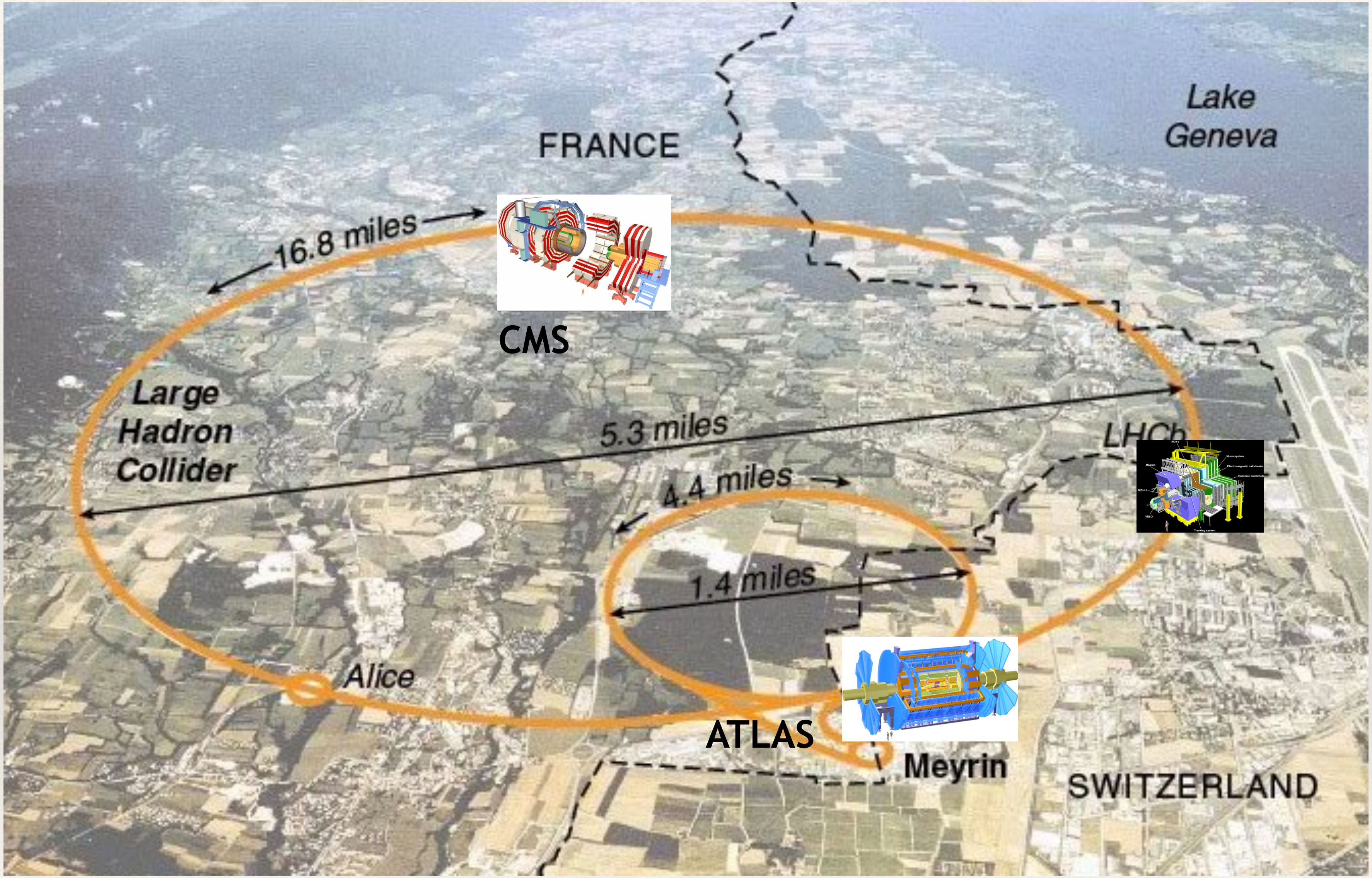
To simplify the life of analysers, instead of 3-5 PDF sets with their uncertainties, a proposal is to provide already an averaged PDFs set, METAPDF set: based on the meta data from various PDF fitting groups



-> discussed at PDF4LHC as one of the possible recommendation case

https://indico.cern.ch/event/343303/

… the LHC


FRANCE
Lake Geneva
16.8 miles
CMS
Large Hadron Collider
5.3 miles
LHCb
4.4 miles
1.4 miles
Alice
ATLAS
Meyrin
SWITZERLAND

Why do we still need to care about PDFs?

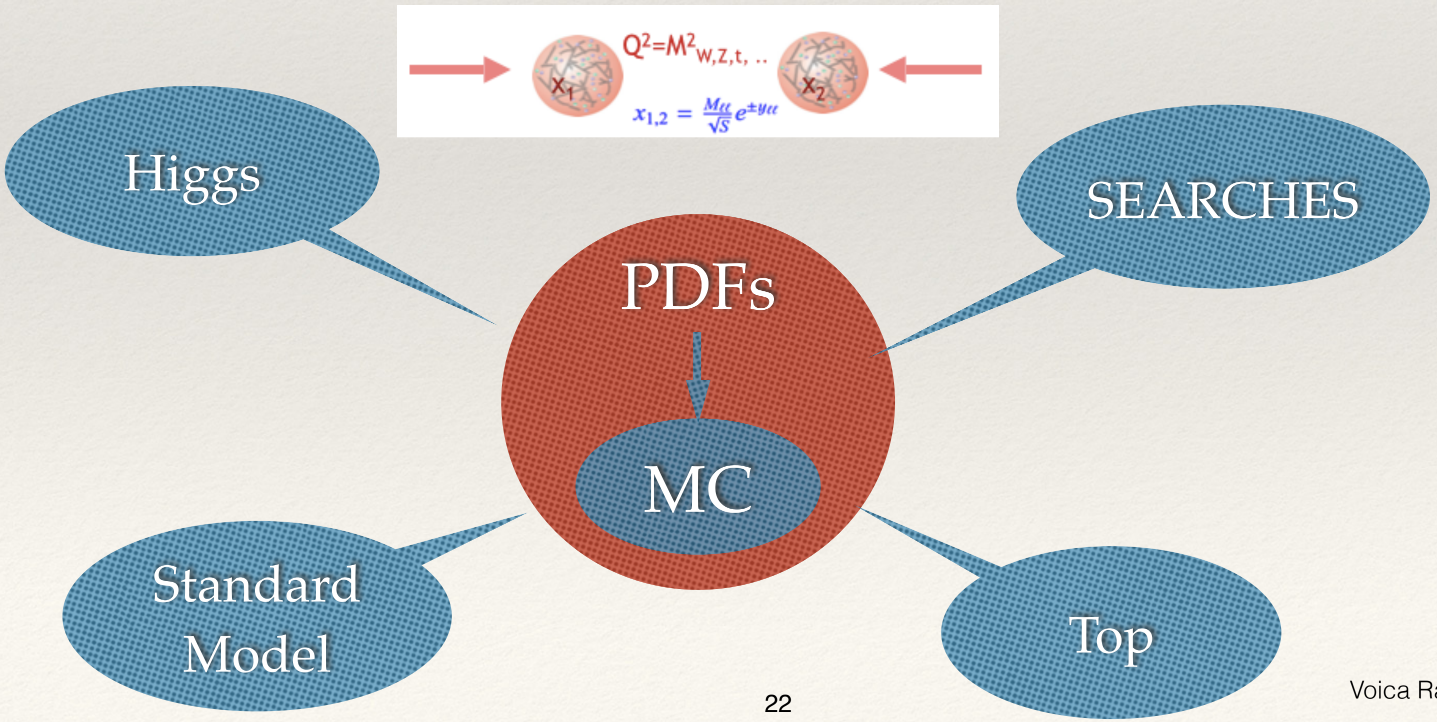
- Discovery of new exciting physics relies on precise knowledge of proton structure.

  - PDFs are one of the main theory uncertainties in Mw measurement
  - PDFs are one of main theory uncertainties in Higgs production.

- **Factorisation theorem:**
  - Cross section can be calculated by convoluting short distance partonic reactions (calculable in pQCD) with PDFs:

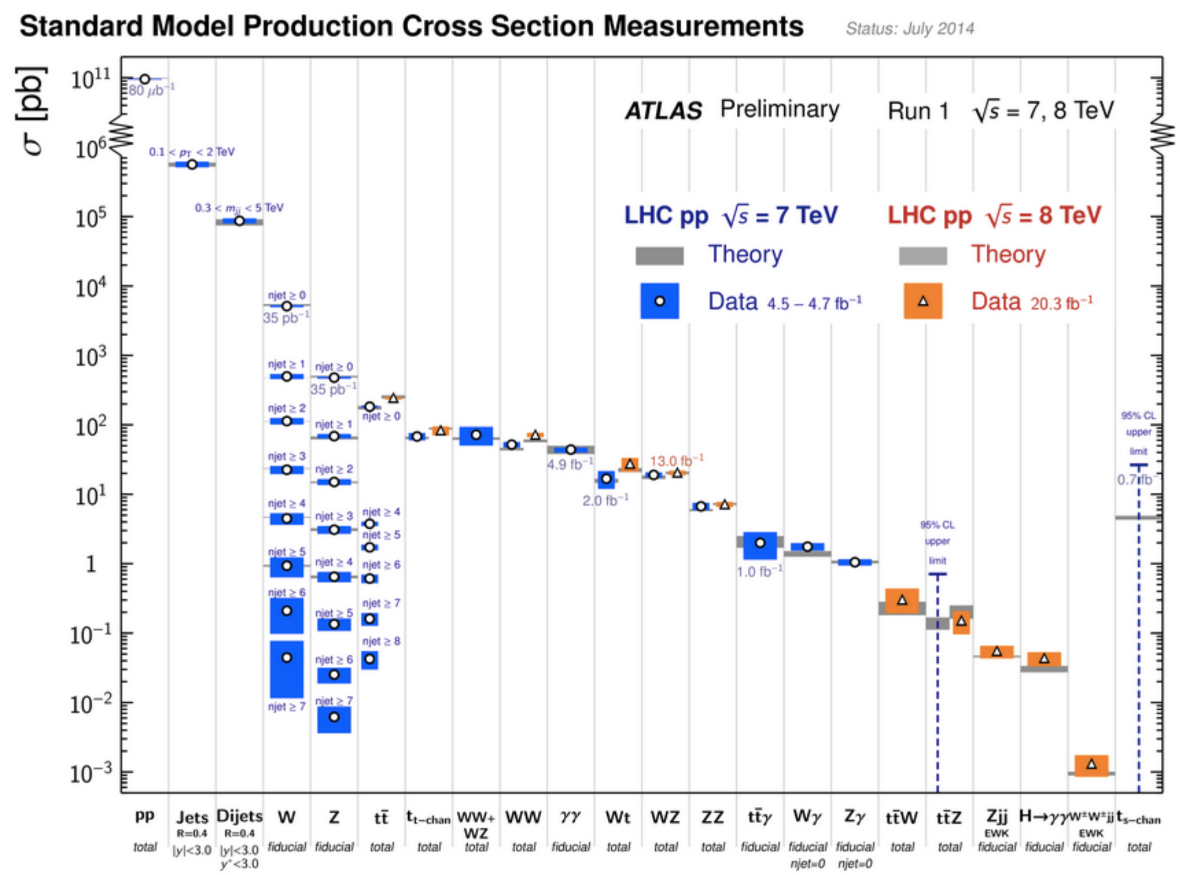
$$d\sigma(\mathrm{h}_1\mathrm{h}_2 \to cd) = \int_0^1 dx_1 dx_2 \sum_{a,b} f_{a/\mathrm{h}_1}(x_1, \mu_F^2) f_{b/\mathrm{h}_2}(x_2, \mu_F^2) d\hat{\sigma}^{(ab\to cd)}(Q^2, \mu_F^2)$$

  - PDFs cannot be calculated in perturbative QCD, however they are process independent (universal) and their evolution with the scale is predicted by pQCD:



The LHC measurements from RUN1

- Successful run in 2010 - 2012 at the LHC confirmed and tested SM



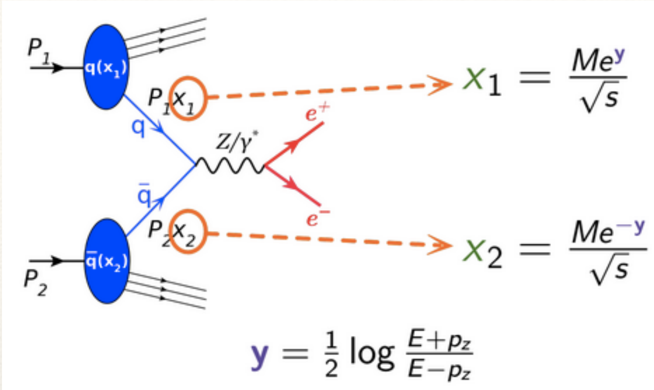
LHC can provide with its multitude of new measurements:

• PDF discrimination by confronting theory with data
• PDF improvement by using LHC data in QCD fit

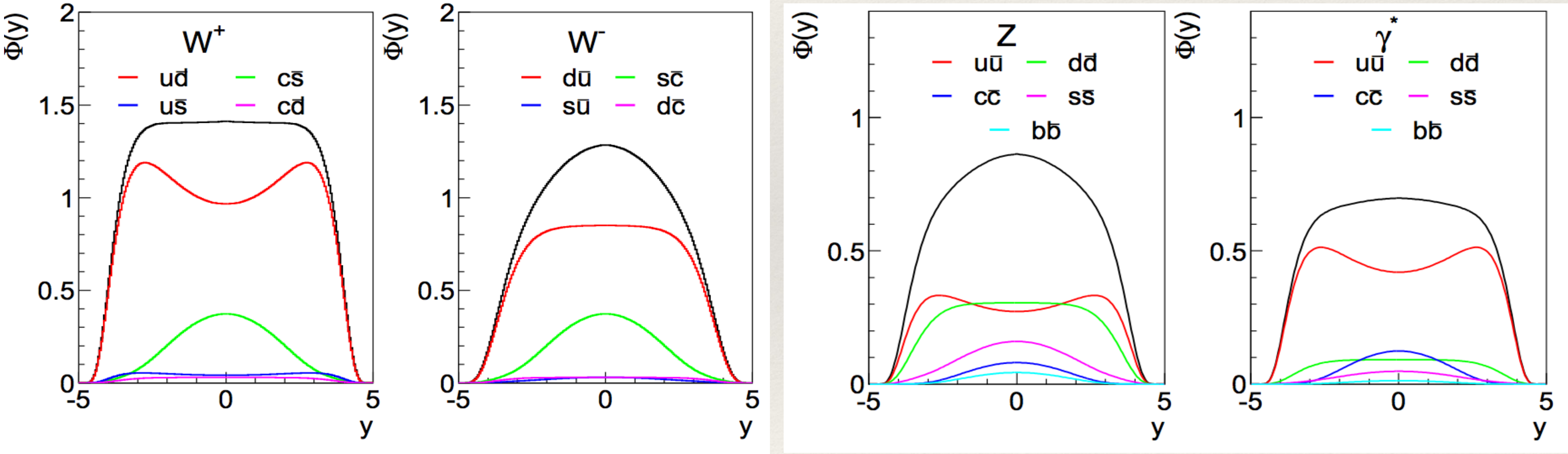
1. W and Z production —> valence, light sea quarks
2. W+c production —> strange
3. Drell-Yan: low and high invariant mass —> sea quarks at high-x, test evolution formalism at low x
4. Inclusive Jet, Di-Jet and Tri-jet production —> gluon and alphas
5. Prompt Photon + Jets —> gluon
6. Top, ttbar —> gluon and alphas
7. W,Z +jets or ZpT —> gluon

Flavour decomposition at LHC (EW bosons)

Additional constraints on PDFs come from DY and jet data at the LHC probe a bi-linear combination of quarks

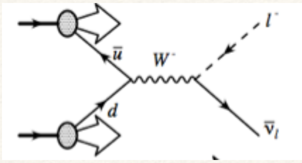


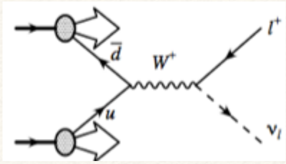
Different flavour decomposition for the W, Z bosons and photon exchange.



Measurements of W, Z production differentially in y_Z and η_ℓ provide information on light sea decomposition

W charge asymmetry

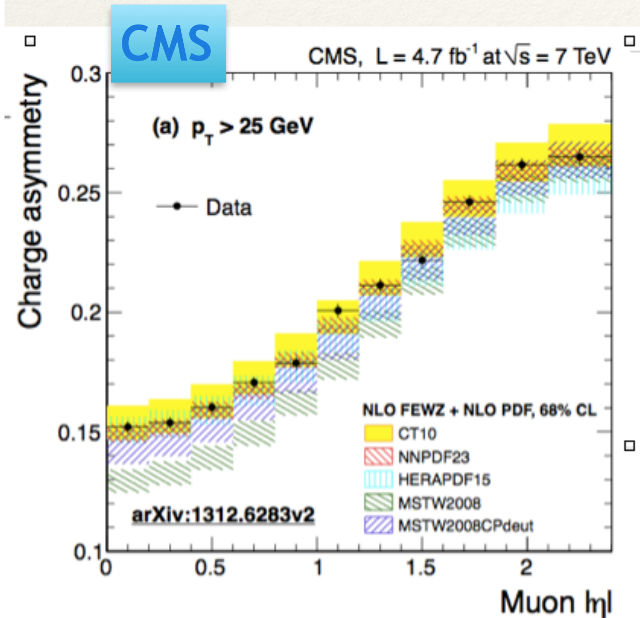




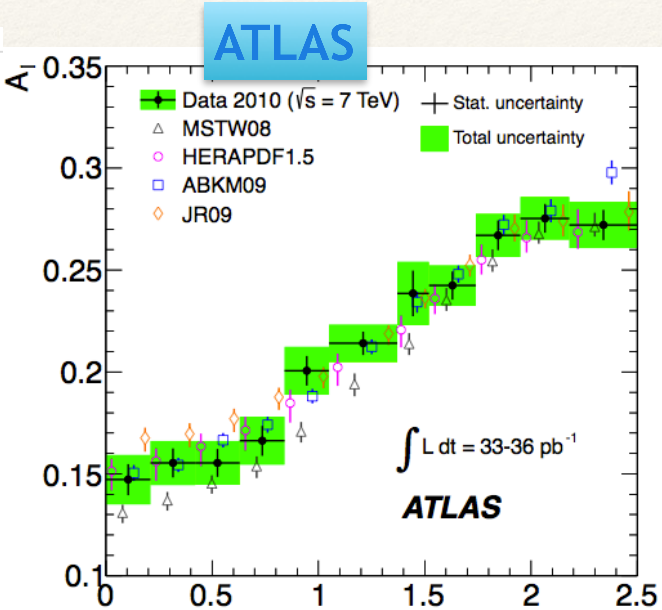
- The interplay between the flavour asymmetries can be enhanced via ratio measurements:
  - W-asymmetry measurement
    - sensitive to uv, dv

$$A_W^l = \frac{d\sigma_{W^+}/d\eta_{l^+} - d\sigma_{W^-}/d\eta_{l^-}}{d\sigma_{W^+}/d\eta_{l^+} + d\sigma_{W^-}/d\eta_{l^-}}$$

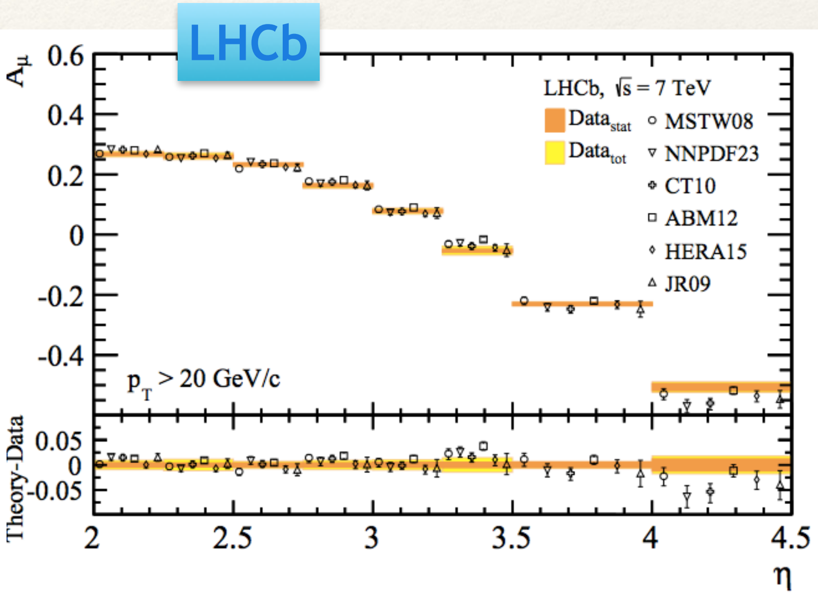
$$A_W \approx \frac{u_v - d_v}{u + d}$$



PRD 90 (2014) 034004



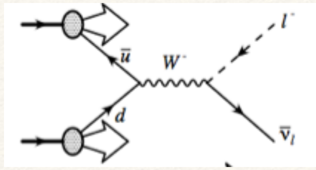
Phys Rev D 85 (2012) 072004

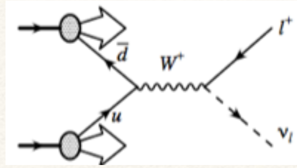


- CMS measures directly the electron asymmetry data from 2011
- ATLAS differential measurements of W^+ and W^- (combined muon and electron) based on 2010 data translated into charge asymmetry Al:
  - proper treatment of correlations are accounted for.
- LHCb extends the measurement to forward region

Selection criteria are optimized for each experiment—> a challenge for data combination

W charge asymmetry





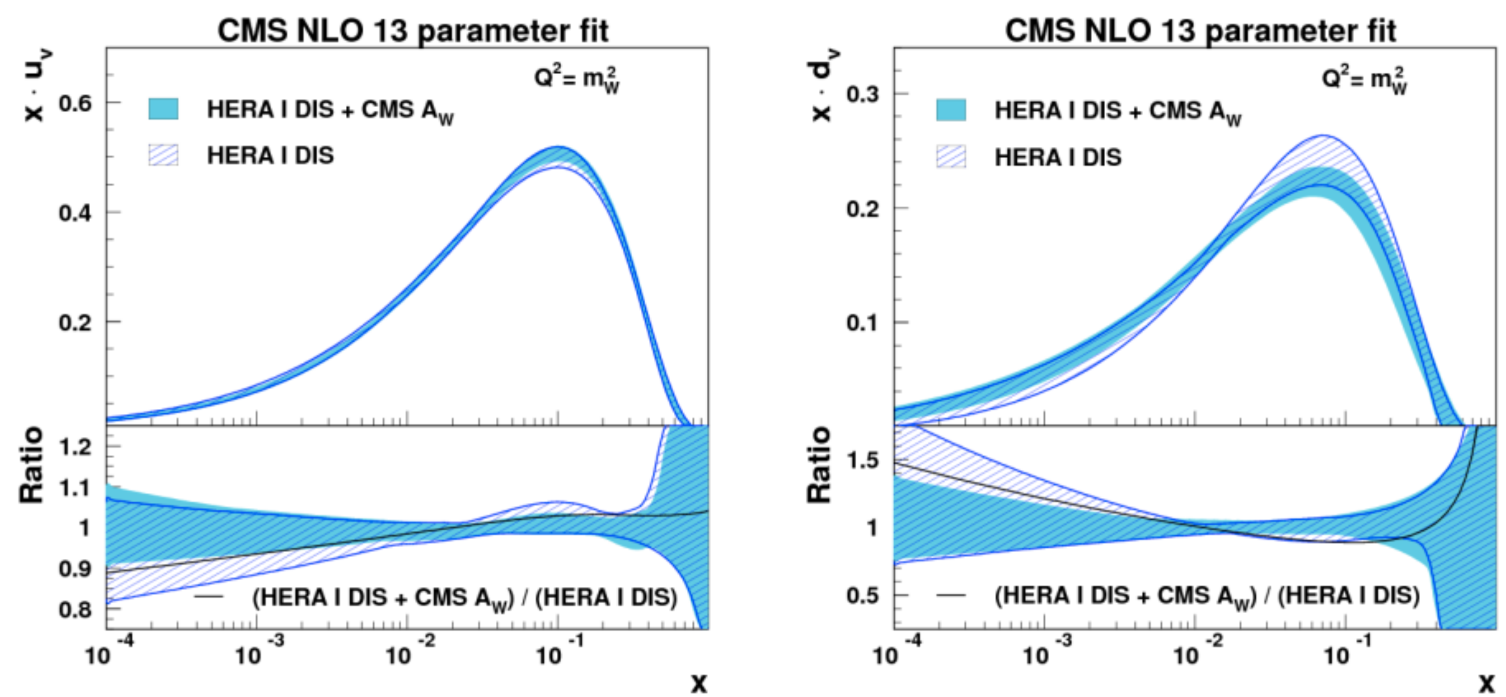
- The interplay between the flavour asymmetries can be enhanced via ratio measurements:
  - W-asymmetry measurement
    - sensitive to uv, dv

$$A_W^l = \frac{d\sigma_{W^+}/d\eta_{l^+} - d\sigma_{W^-}/d\eta_{l^-}}{d\sigma_{W^+}/d\eta_{l^+} + d\sigma_{W^-}/d\eta_{l^-}} \qquad A_W \approx \frac{u_v - d_v}{u + d}$$

- A PDF fit of these CMS muon asymmetry data together with the combined HERA-I inclusive deep inelastic scattering (DIS) data shows the potential of the LHC data to constrain valence quarks

(JHEP 1001 -109)

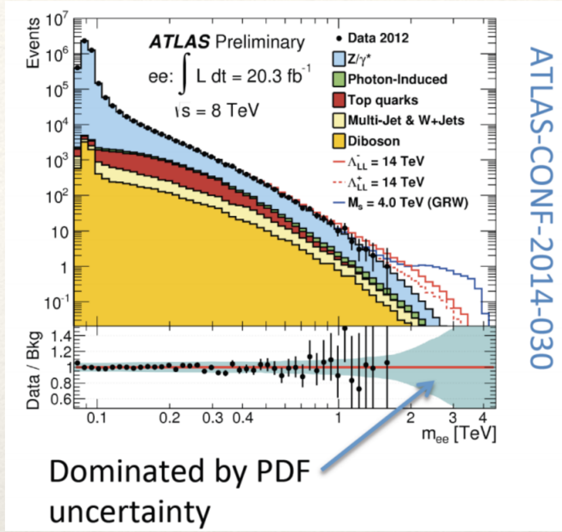
The largest effect is on the $u_{valance}$ and $d_{valance}$ PDFs ($0.001 < x < 0.1$).



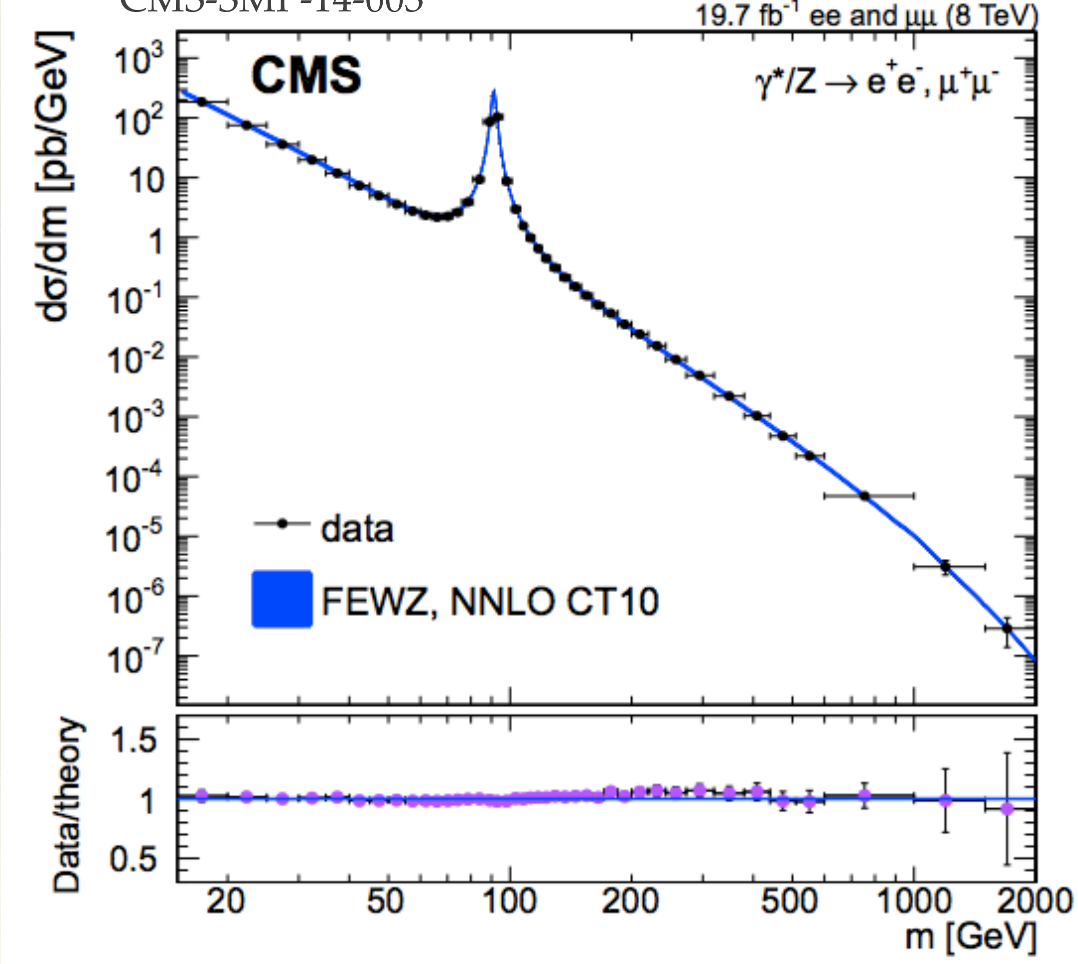
Change of PDF shape, improved constraints on the valence distributions

Neutral Current Drell Yan di-lepton measurements

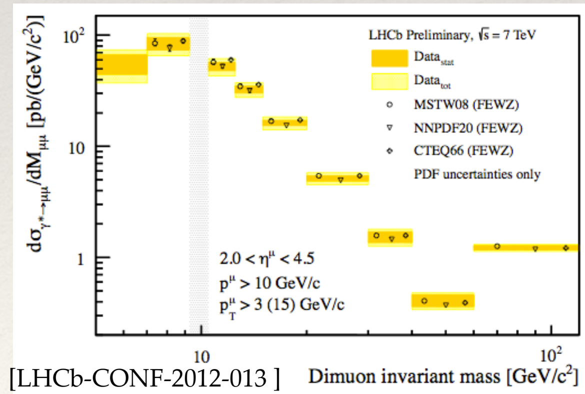
- Drell Yan Mass Spectra data are sensitive to new physics at high-scale and can give information on sea quark PDFs.
  - High Mass: sensitive to sea quarks at high x
  - Low Mass: sensitive to low x region (test of DGLAP?)



CMS-SMP-14-003



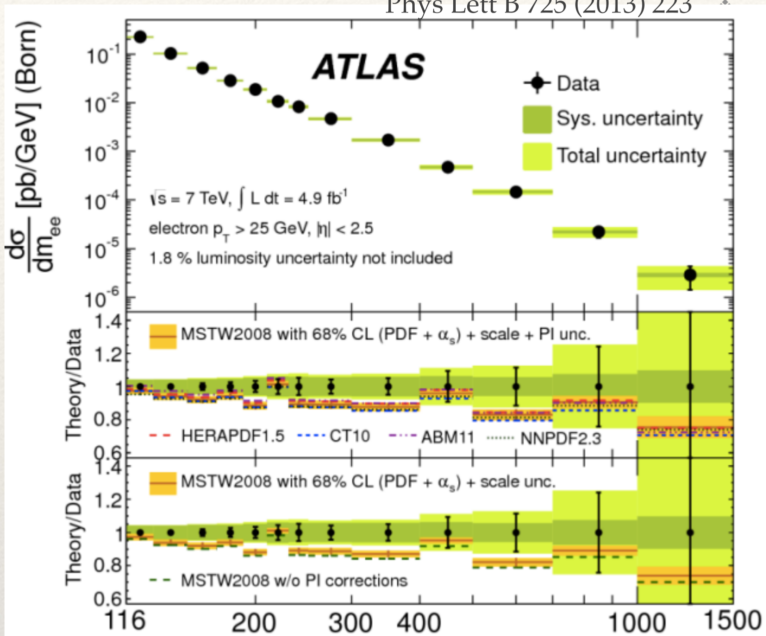
Measurement complemented by LHCb down to 5 GeV



Z Differential Cross Section (off resonance region)

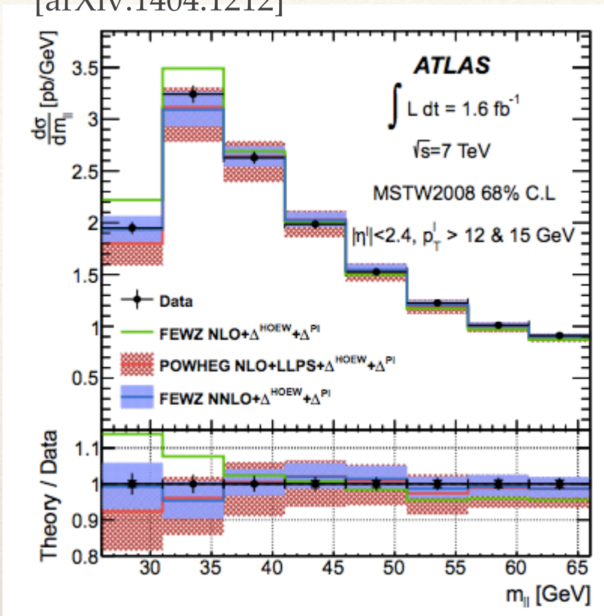
❖ High Mass region — Low Mass Region

Phys Lett B 725 (2013) 223



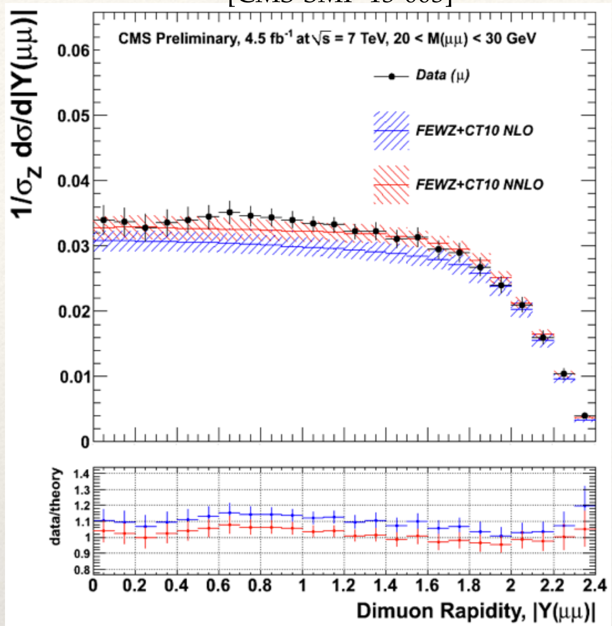
ATLAS: $116<M_{ee}<1500$

[arXiv:1404.1212]



ATLAS: $20<M_{ll}<65$

[CMS-SMP-13-003]



CMS: $20<M_{\mu\mu}<30$

The Drell Yan invariant mass spectrum in the off resonance region:

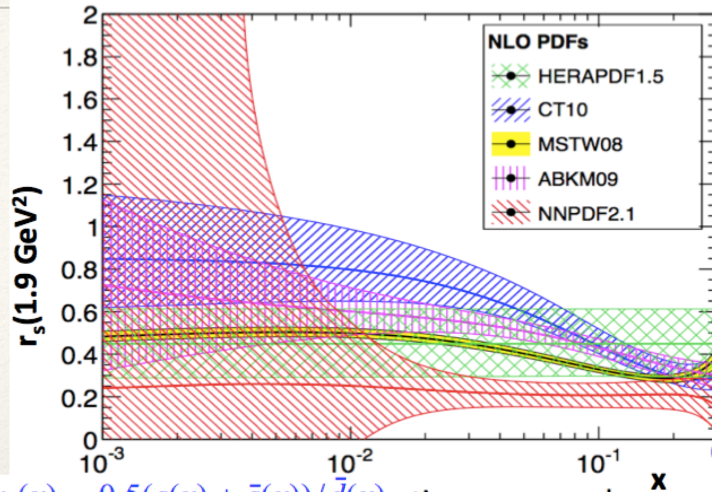
- ATLAS in the dilectron channel for high mass region, combined in low mass region
- normalized to the Z resonance region, function of dimuon rapidity for CMS in selected $M_{\mu\mu}$ bins

Data is confronted with NNLO predictions corrected for NLO EW effects

- Currently all PDFs shown give a good description

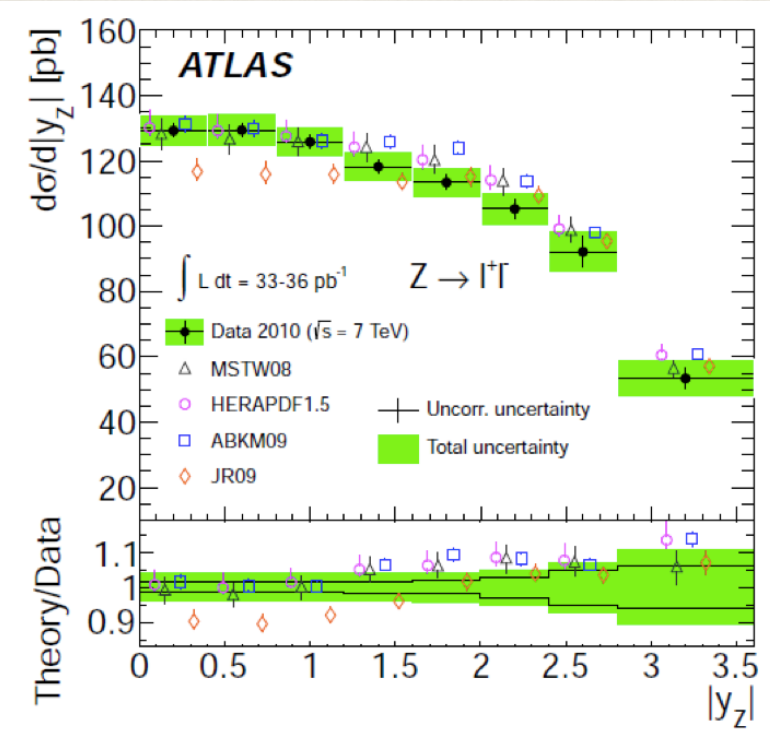
Strange quark from W, Z measurements at LHC

- Strange quark is not so well constrained —>
  - Neutrino dimuon data provides constraints:
    - prefers rather strongly suppressed strange (s/d)
    - data suffer from corrections: i.e. nuclear target
- In 2010, at LHC the EW boson data was used to constrain strange quark through a QCD fit analysis
  - Impact comes mainly from Z rapidity distribution:



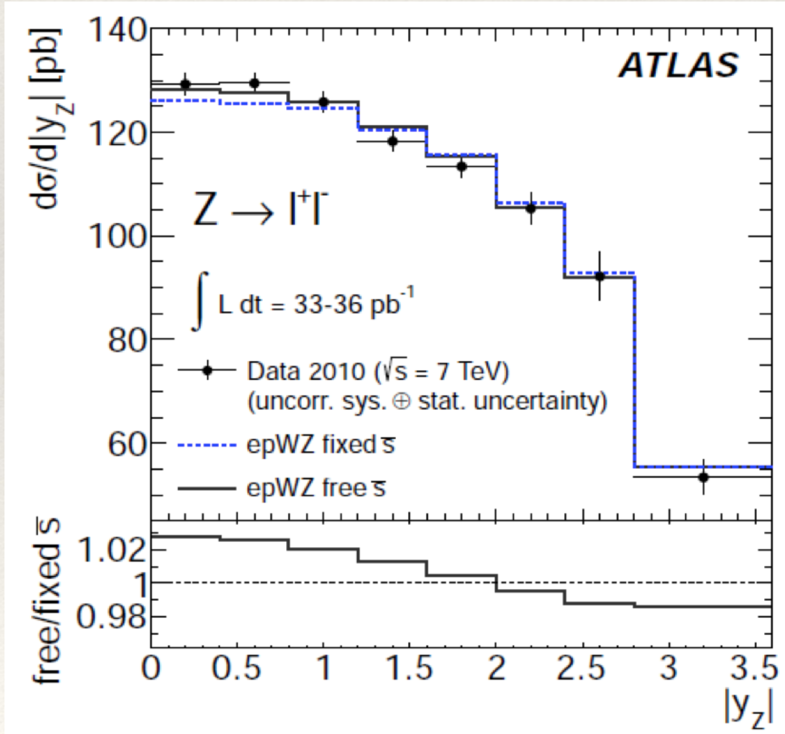
$r_s(x) = 0.5(s(x) + \bar{s}(x))/\bar{d}(x)$

G.Watt



—> discrimination among PDFs

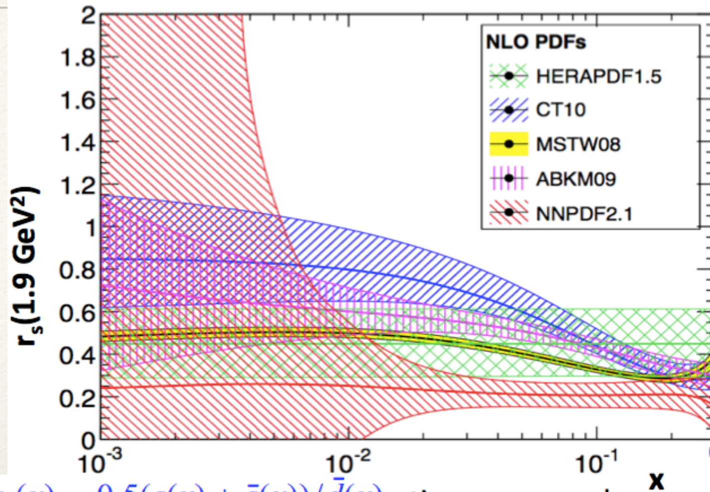
enhanced strange <— is preferred by data



Phys Rev Lett 109 (2012) 012001

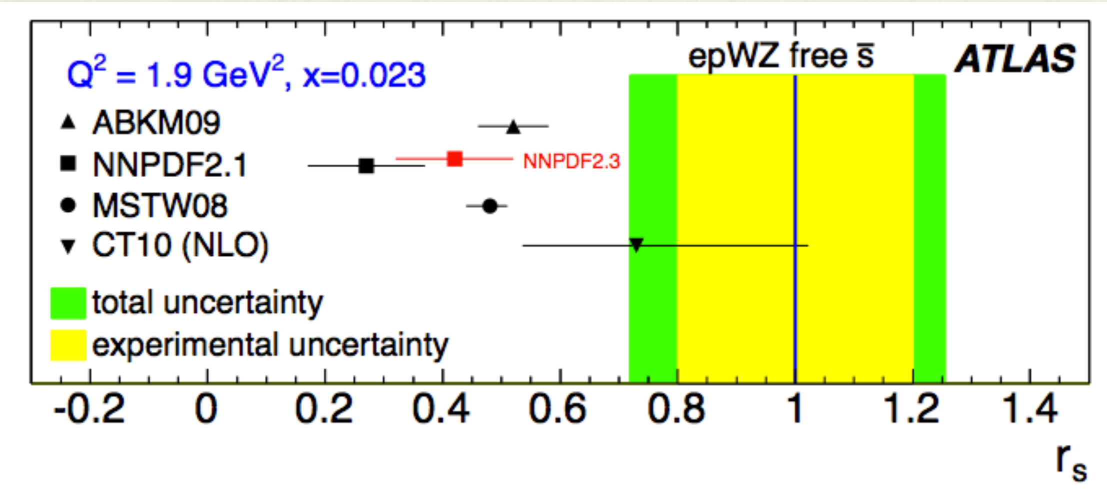
Strange quark from W, Z measurements at LHC

- Strange quark is not so well constrained —>
  - Neutrino dimuon data provides constraints:
    - prefers rather strongly suppressed strange (s/d)
    - data suffer from corrections: i.e. nuclear target
- In 2010, at LHC the EW boson data was used to constrain strange quark through a QCD fit analysis
  - Impact comes mainly from Z rapidity distribution:

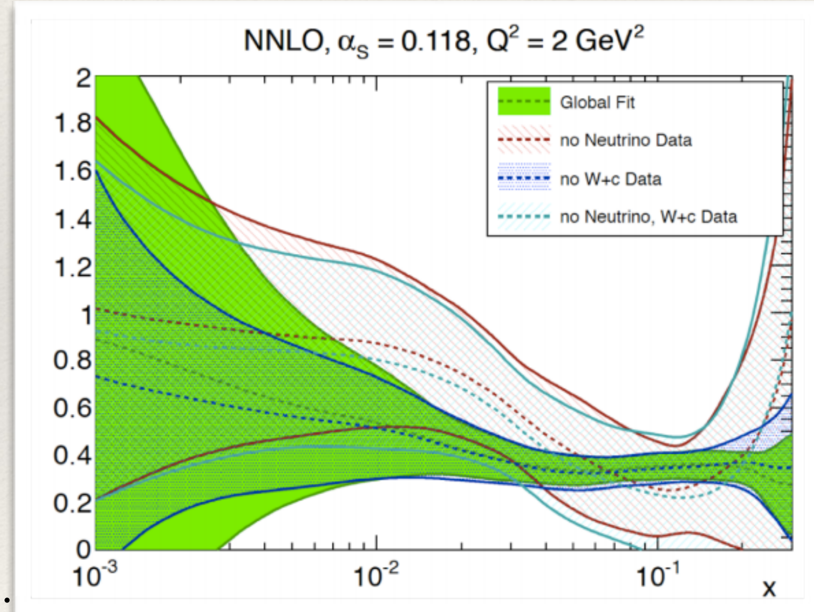


$r_s(x) = 0.5(s(x) + \bar{s}(x))/\bar{d}(x)$

G.Watt



Phys Rev Lett 109 (2012) 012001



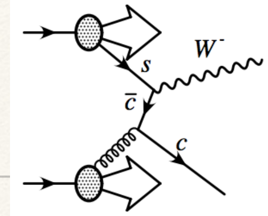
NNPDF3.0

rs=(s+sb)/(ub+db)

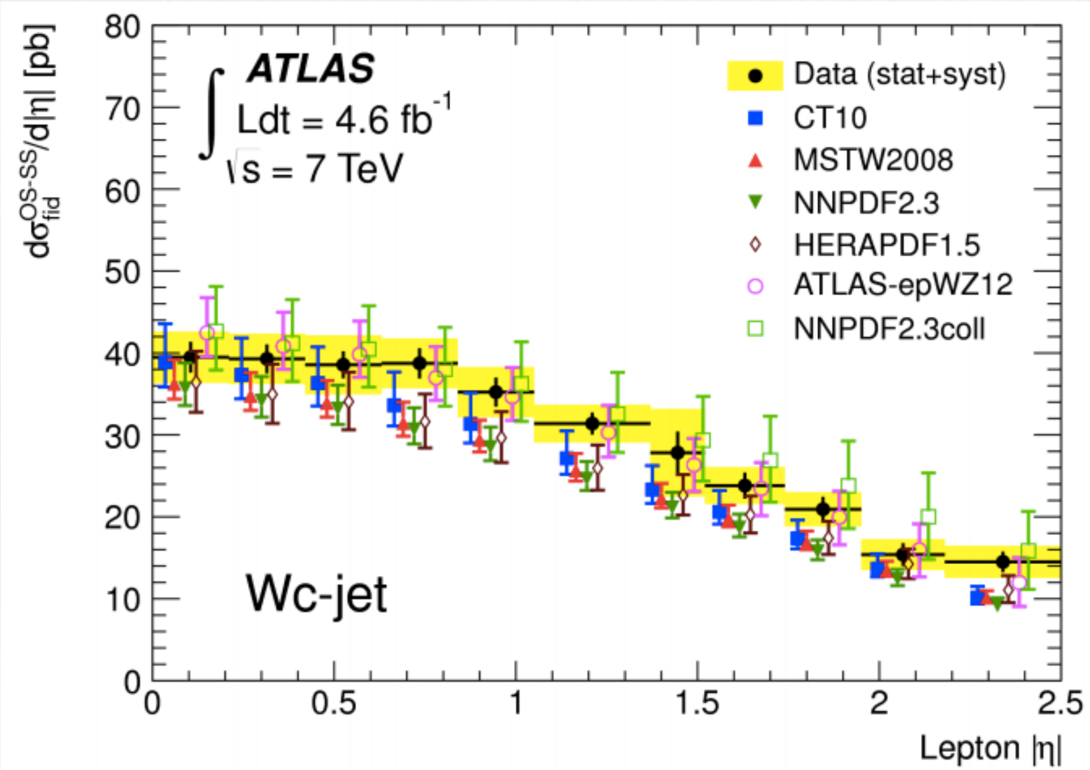
Since then, new measurements and QCD analyses were performed:

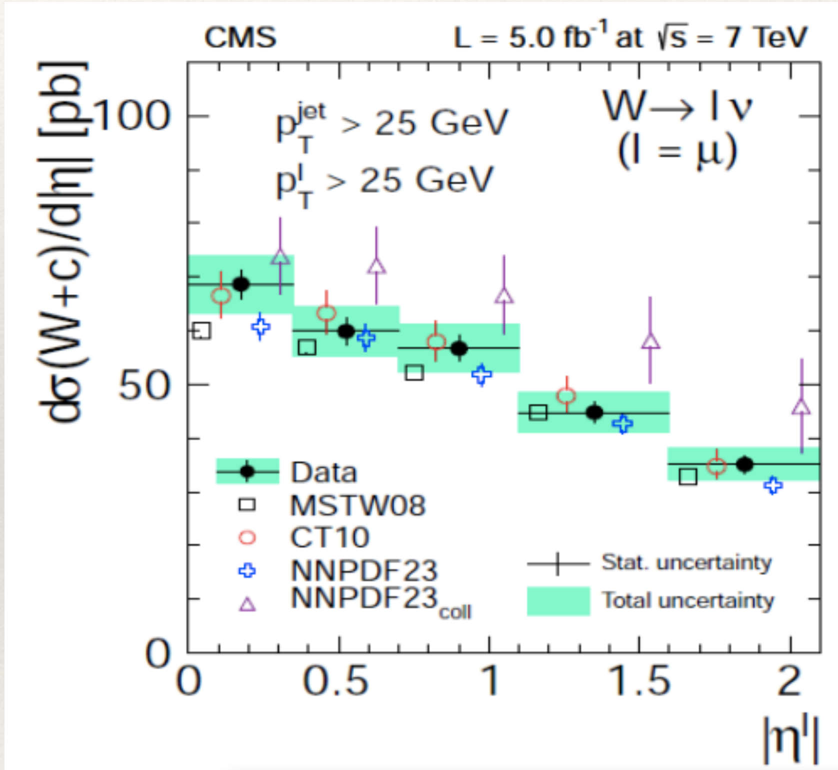
- W+charm from ATLAS and CMS
- QCD fits to W asymmetry + W+charm data @ CMS

W+c sensitivity to strange



- W + charm data is directly sensitive to the strange quark density
- Both ATLAS and CMS have performed dedicated measurements:
  - Measure fully reconstructed D* mesons or soft leptons within a jet
    - ATLAS @ particle level [arXiv:1402.6263v1] CMS @parton level [arXiv:1310.1138]

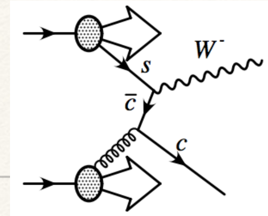




—>consistent with ATLAS-epWZ12 (PDF set from ATLAS Z, W inclusive)

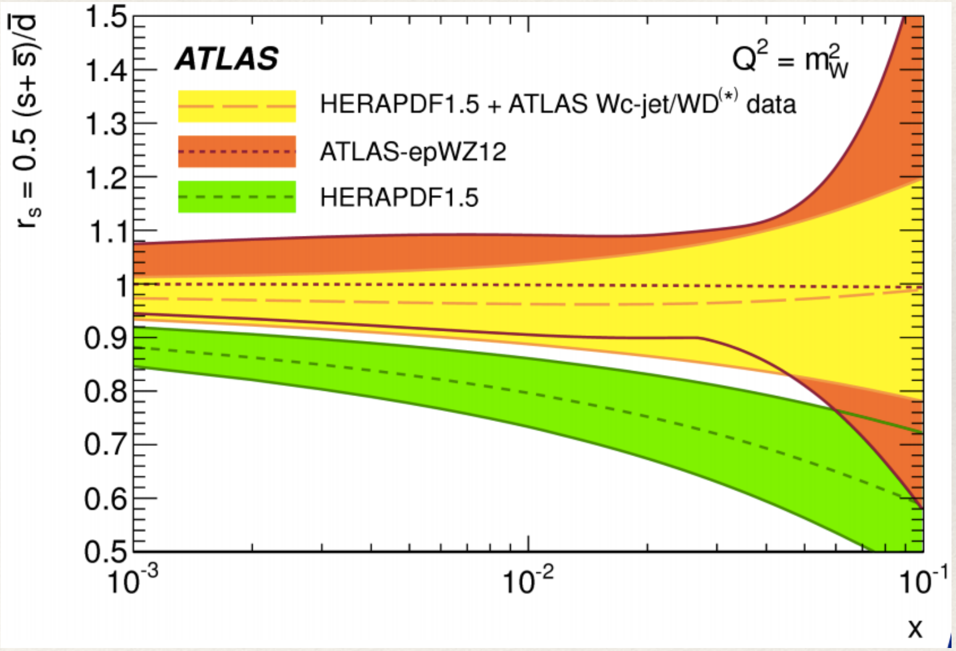
—>consistent with CT10

W+c sensitivity to strange

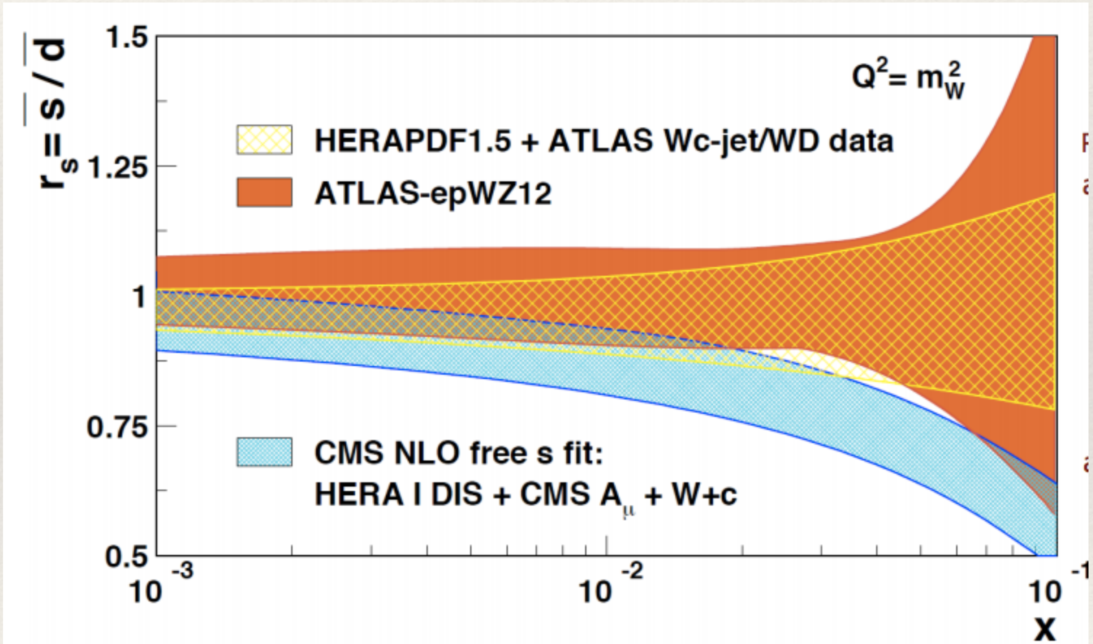


- W + charm data is directly sensitive to the strange quark density
- Both ATLAS and CMS have performed dedicated measurements:
  - Measure fully reconstructed D* mesons or soft leptons within a jet
    - ATLAS @ particle level [arXiv:1402.6263v1] CMS @parton level [arXiv:1310.1138]

PDF eigenvector analysis



QCD Fit

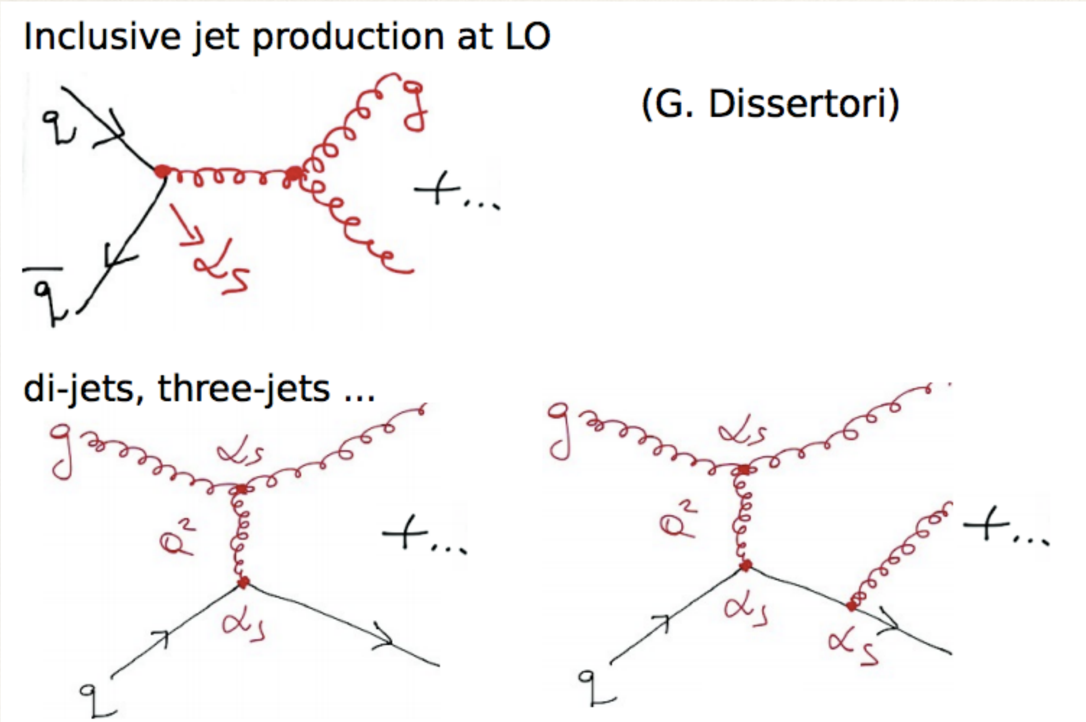


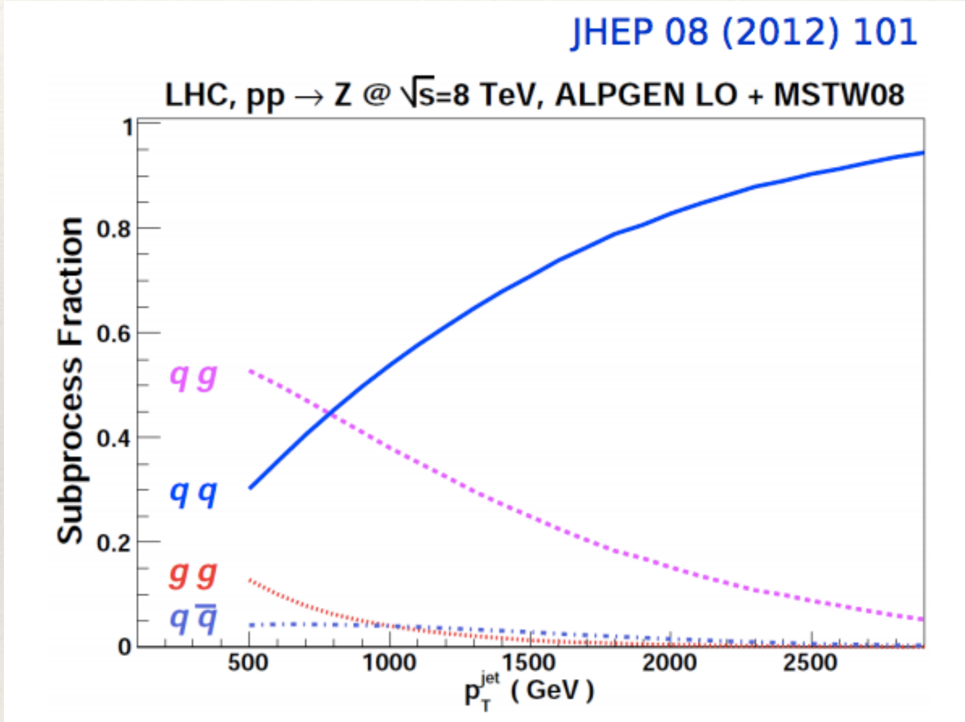
—> consistent with previous ATLAS result

—> mild differences between CMS and ATLAS

Jet Production at the LHC

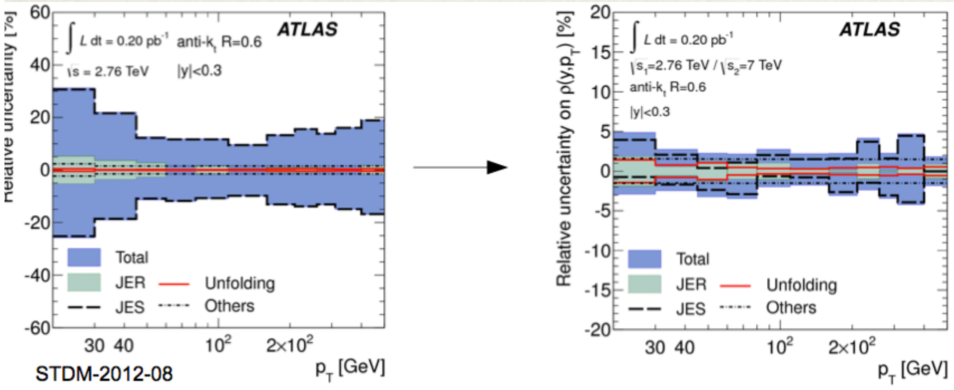
- Jet production at the highest scales may reveal new physics and the reliability of the predictions depends on how well we know the high-x gluon PDF
- Jet production at LHC provides information about hard QCD, PDFs , strong coupling:
  —> PDFs and alphas depend on scale of the process (Pt of the jet)



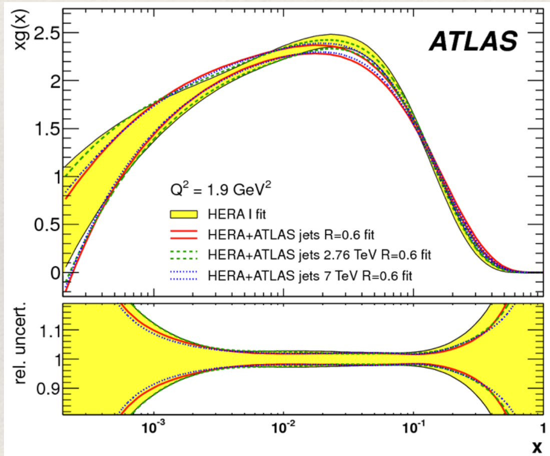


Inclusive Jet Production at the LHC

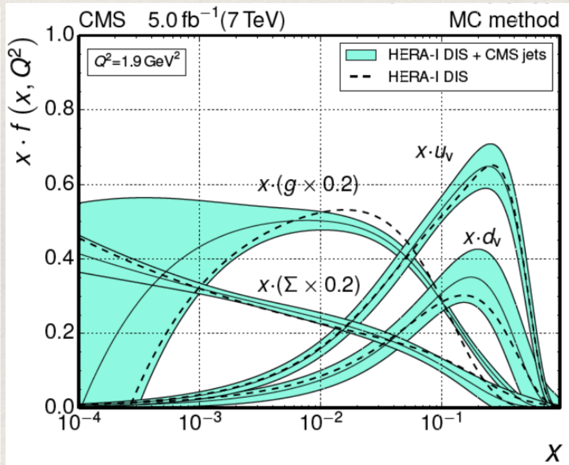
- ATLAS 2010 7 TeV jet data [PRD86(2012)014022] shows no sign of new physics and are fitted well by NLO.
- To enhance the impact of jet data, it's smart to consider ratios:
  - the major experimental systematic - the Jet Energy Scale - cancels out (i.e. 2.76 vs 7 TeV data)
- The impact of the LHC 7 TeV inclusive jet data on proton PDFs is investigated by including the jet cross section measurement in a combined fit with the HERA-I inclusive DIS cross sections.



EPJC (2013) 73 2509



arXiv:1410.6765

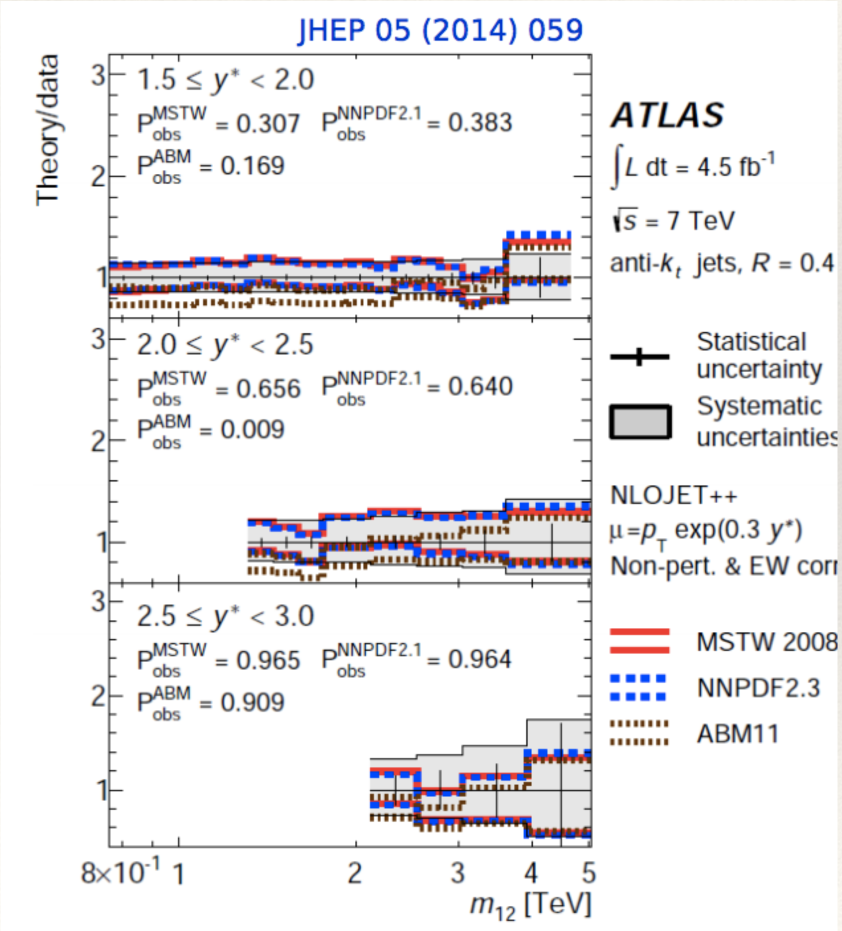


- jet data can help to improve gluon distribution function in high-x region
- Data can be used in extracting strong coupling, however one limitation is that jet calculations are still only available to NLO and there is thus still a substantial scale dependence on predictions

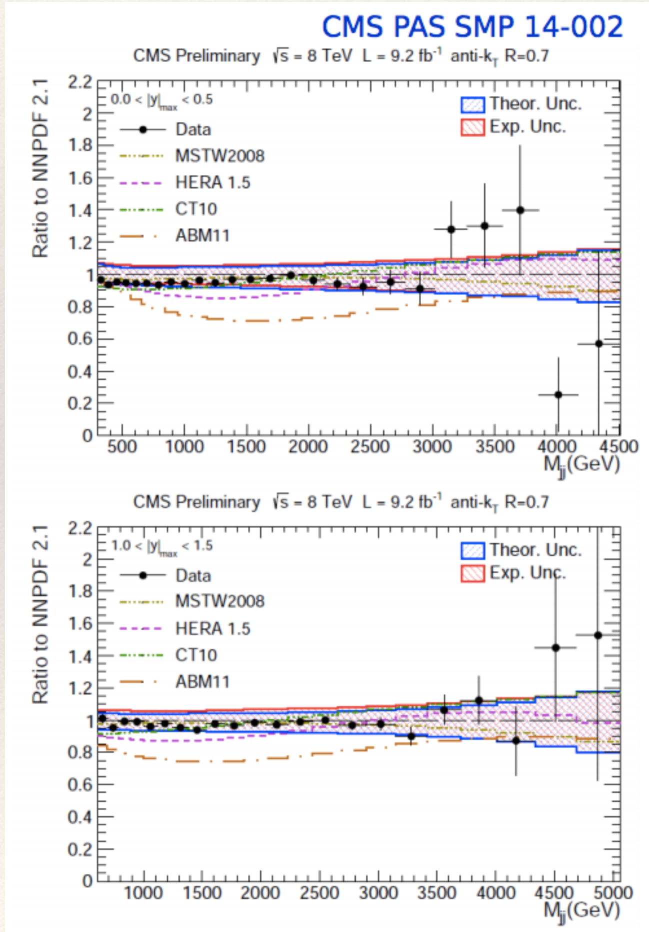
$$\alpha_S(M_Z) = 0.1185 \pm 0.0019(\mathrm{exp}) \pm 0.0028(\mathrm{PDF})\,^{+0.0055}_{-0.0022}(\mathrm{scale})$$

Di-Jet, Tri-Jet Production at the LHC

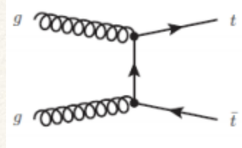
- More data on jets from 2011 7 TeV running (4.5/fb of data) from CMS and ATLAS
- a detailed statistical analysis is performed when comparing data to theory
  -> lowest probability found for ABM



Different rapidity and mass ranges are considered.

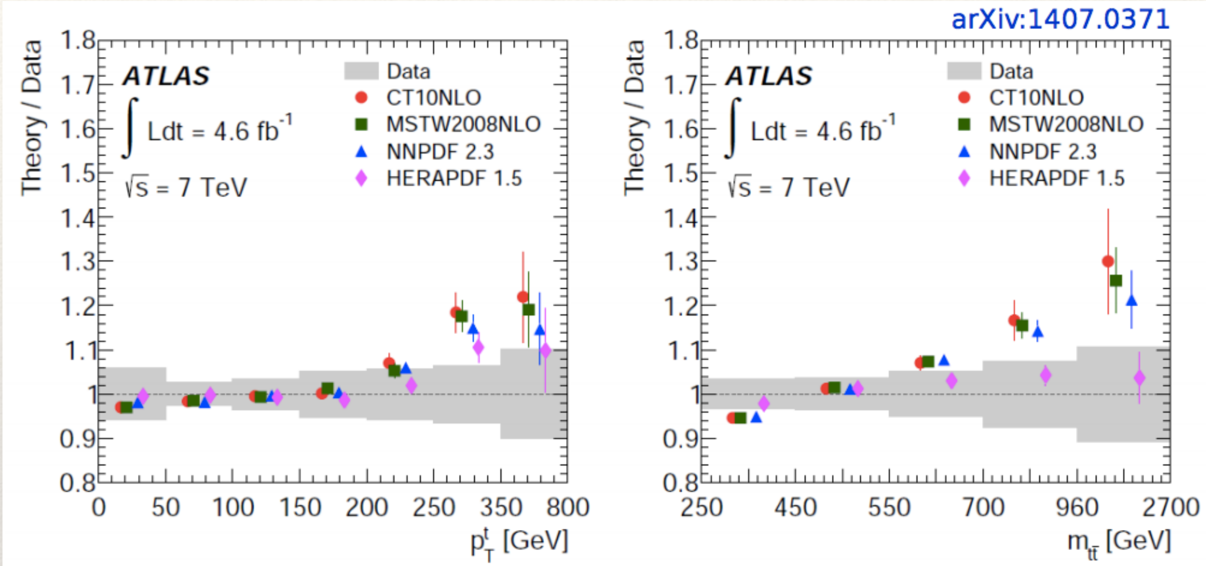


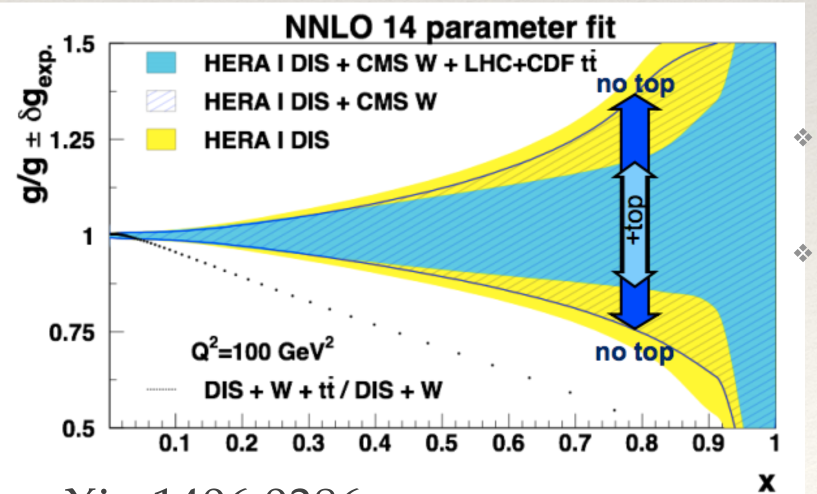
Sensitivity to PDFs from Top Production



- Top-quark pair production at the LHC probes high-x gluon ($x \approx 0.1$):
  —> there is a strong correlation between g(x), αs and the top-quark mass mt
- Precise measurements of the total and differential (normalised and absolute) cross section of ttbar pair production can constrain and de-correlate αs, gluon, mt

  - compared with theory (NLO) using different PDFs
  - NNLO theory calculations are becoming available ...





arXiv:1406.0386

- QCD analysis with ATLAS and CMS ttbar data (together with HERA, Tevatron and W production data at LHC)
- moderate improvement of the uncertainty on the gluon distribution for $x > 0.1$ and significant change of the shape of the gluon distribution observed

Impact of LHC data on PDFs

- Abundant LHC data with possible novel constraints on PDFs are investigated:

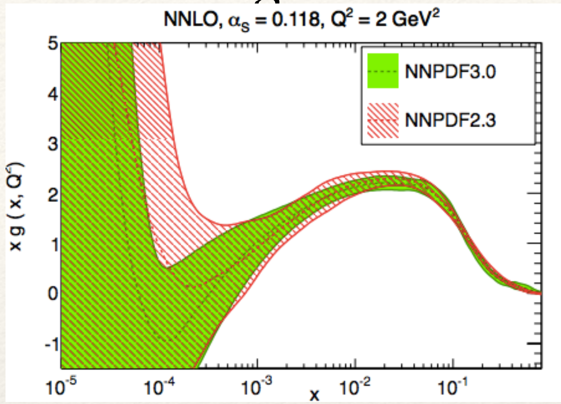
GLUON
- Inclusive jets and dijets **(medium/large x)**
- Isolated photon and γ+jets **(medium/large x)**
- Top pair production **(large x)**
- High p_T Z(+jets) distribution **(small/medium x)**

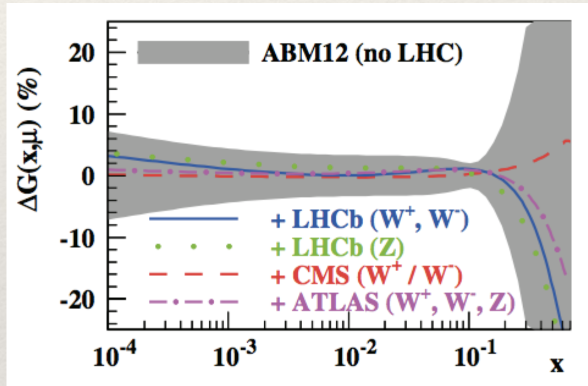
QUARKS
- High p_T W(+jets) ratios **(medium/large x)**
- W and Z rapidity distns **(medium x)**
- Low and high mass Drell-Yan **(small and large x)**
- Wc **(strangeness at medium x)**

PHOTON
- Low and high mass Drell-Yan
- WW production

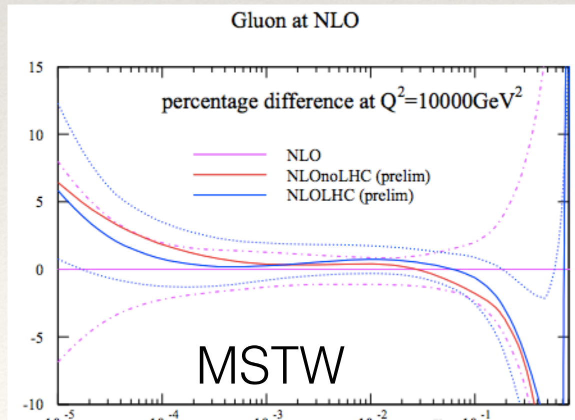
Intense activity of global PDF groups to include these measurements in the new PDF releases in time for Run2 data.

PDF4LHC, QCD@LHC, 2014



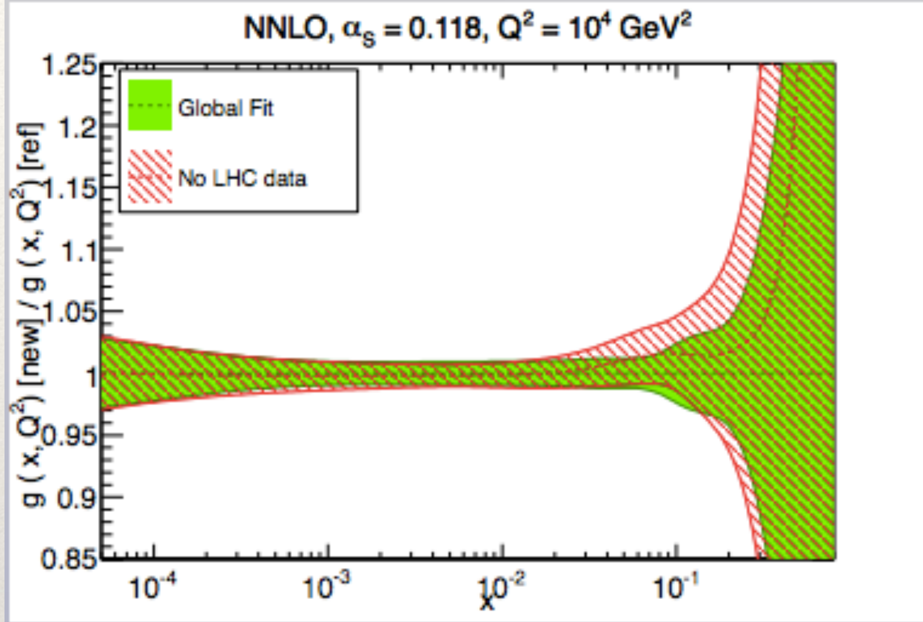


NNPDF3.0 is in LHAPDF, announced updates from: MSTW, HERA, CT, ABM

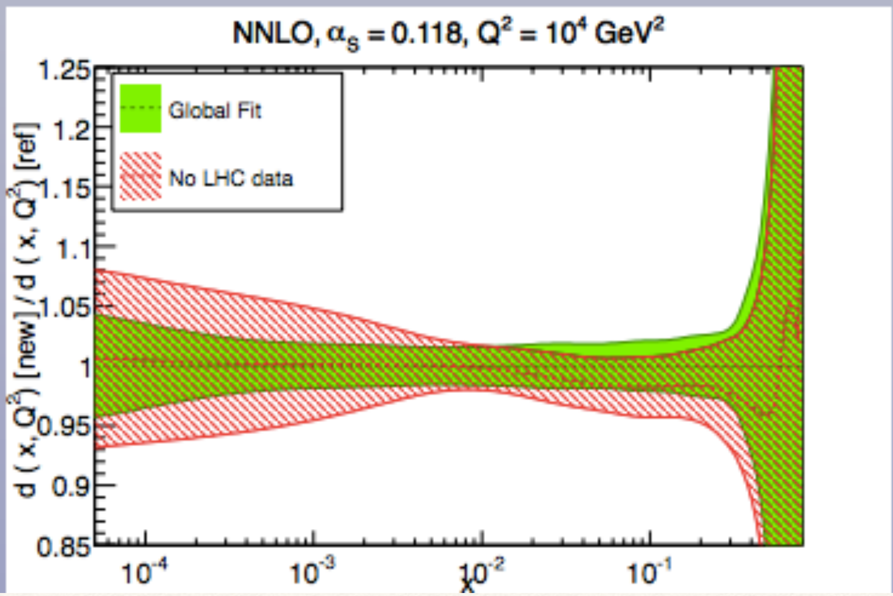


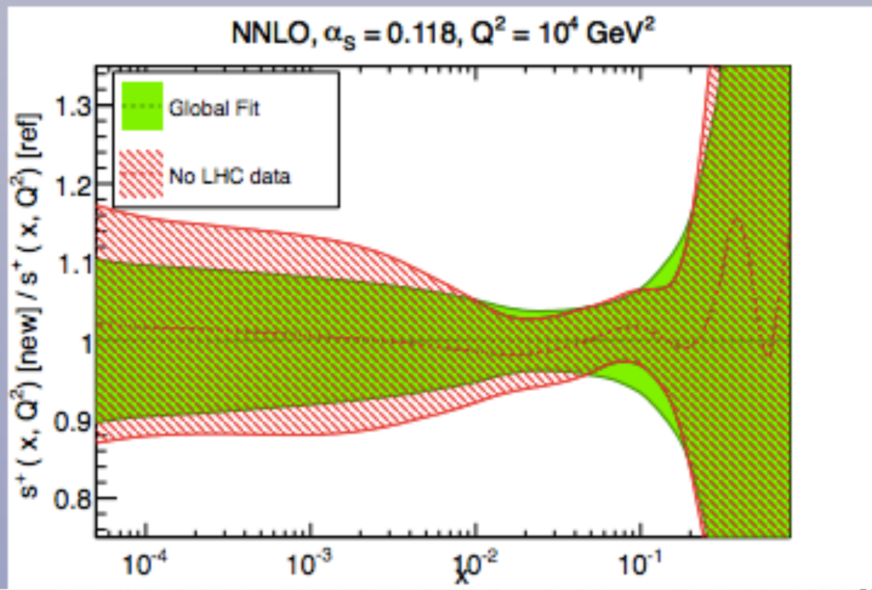
Impact of LHC data on PDFs

- Abundant LHC data with constraints on PDFs are investigated: J. Rojo [NNPDF3.0]



- Compare global NNPDF3.0 fit with a fit **without LHC data**
- PDF uncertainties on **large-x gluon** reduced due to **top quark and jet data**
- PDF uncertainties on **light quarks** reduced from the **Drell-Yan** and **W+charm data**
- The **description of all new LHC data**, already good in NNPDF2.3, is further improved in NNPDF3.0

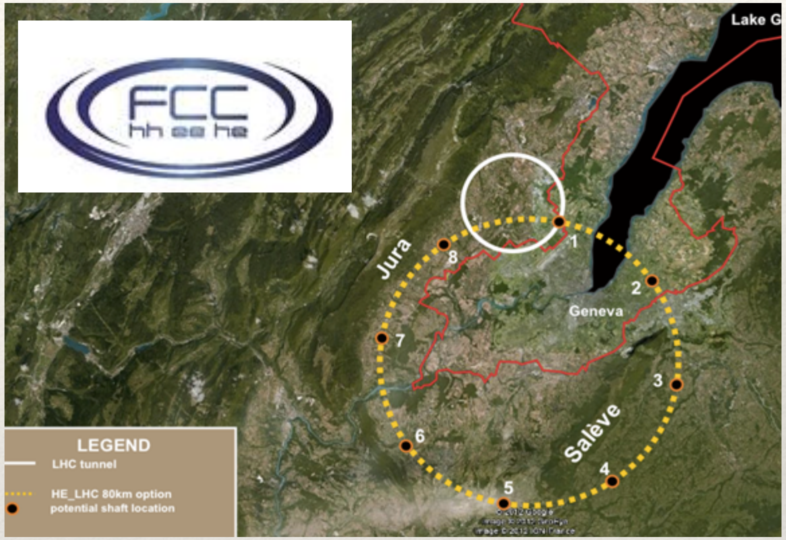


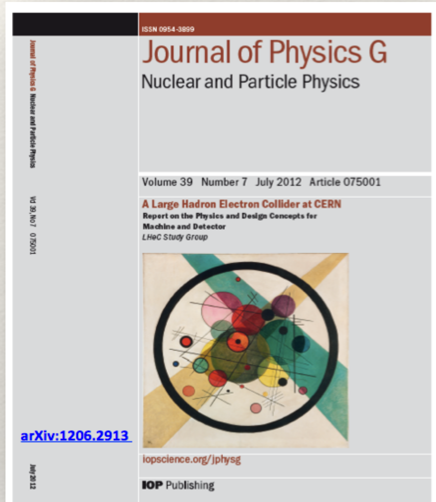


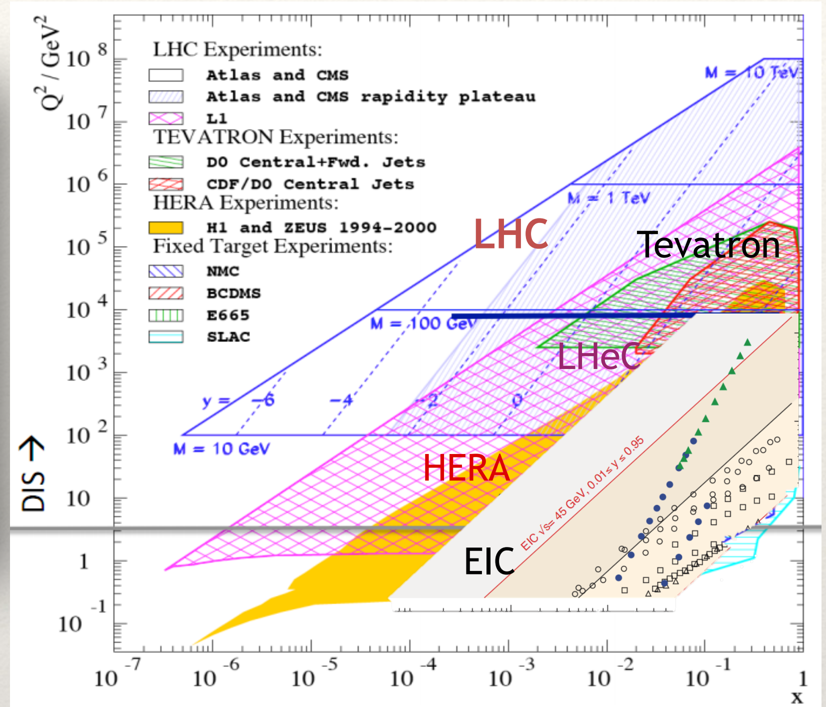
Future prospects for better PDFs?

- Future Circular Collider project: ee, eh, hh (100 TeV proton)

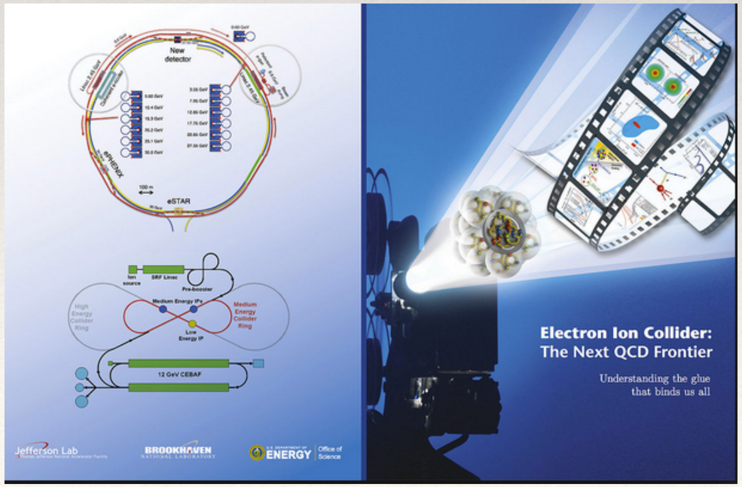
https://espace2013.cern.ch/fcc/Pages/default.aspx

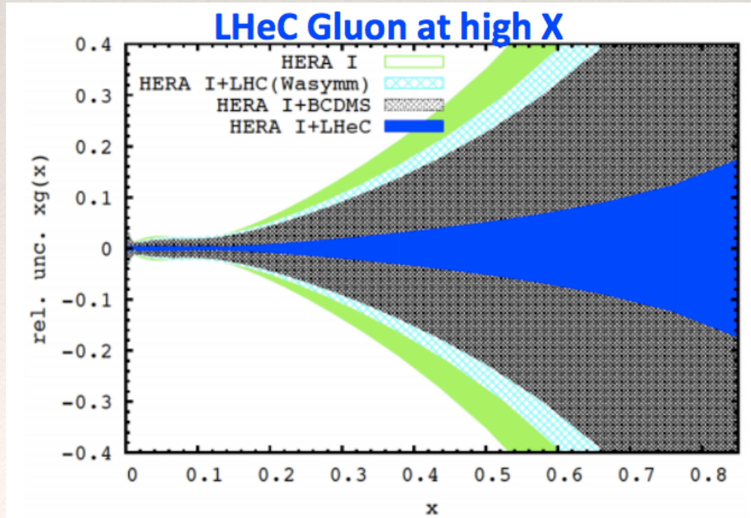




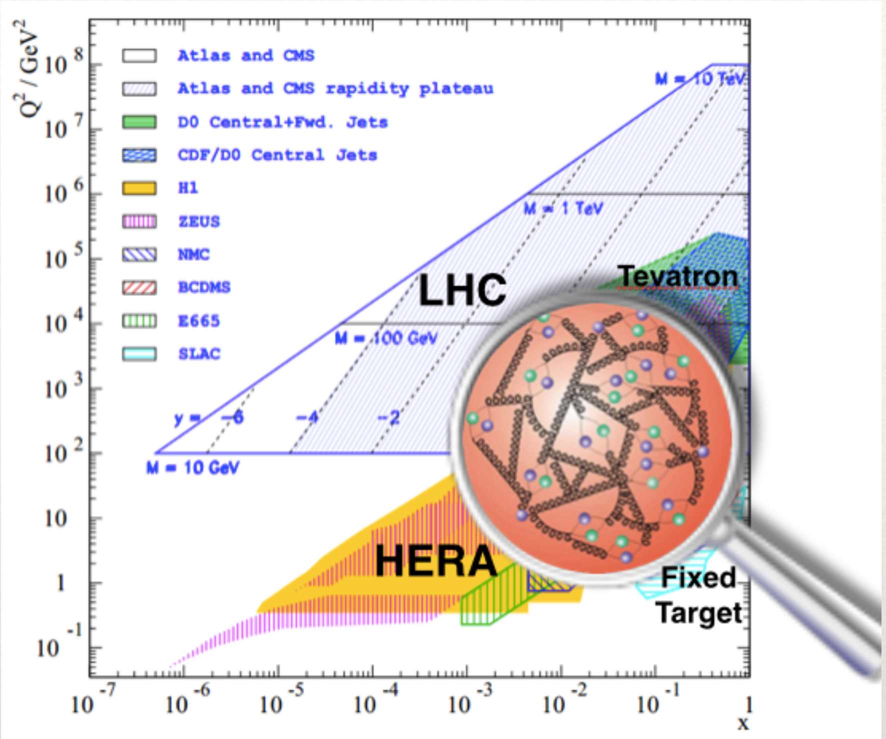


- EIC project (JLAB, Brookhaven)





Summary



PDFs are very important as they still limit our knowledge of cross sections whether SM or BSM.

- HERA has finalised its separate measurements relevant to PDFs and ongoing efforts on combining final measurements to reach its ultimate precision:
  - PDFs, mc, mb, alphas …
- JLAB has a dedicated program on improving PDFs at high x

- Standard Model LHC measurements can themselves contribute to PDF discrimination and PDF improvement:

… Many more valuable measurements are already available, but not covered in this talk …

- More precision measurements from LHC to come from Run I and in future from Run 2
- Intense activity of PDF groups to include constraining information in new releases
- Future Facilities to further push our limits are being considered …

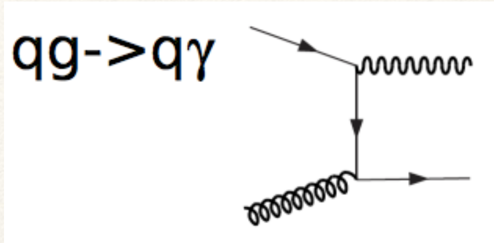
Many Thanks!

back-up slides

not necessarily useful …

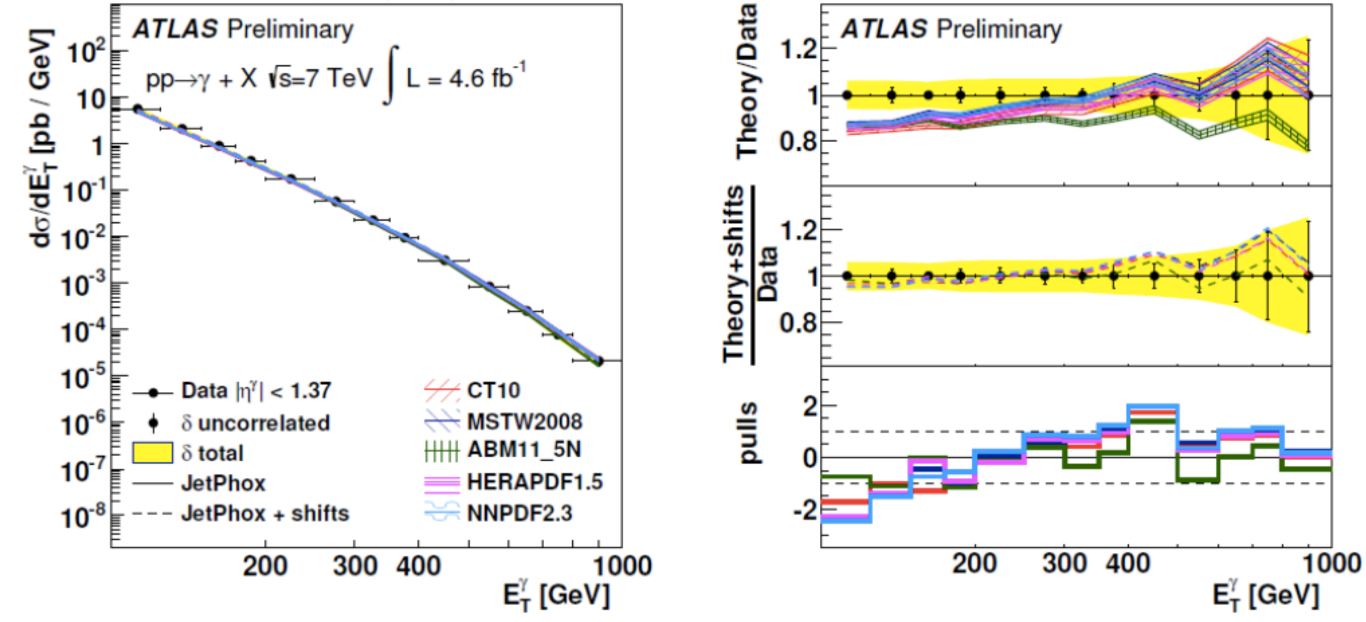
Prompt Photon production at LHC

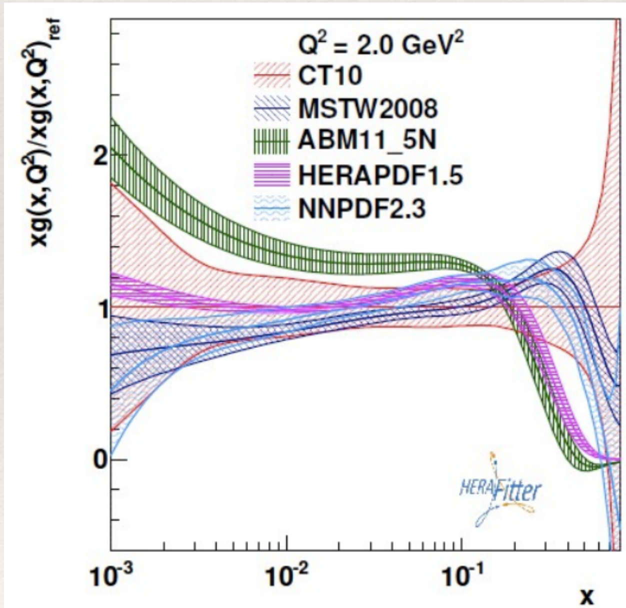
- Prompt photon data at LHC is sensitive to gluon content at high x
  - Dominantly via Compton-like process



- ATLAS study of the inclusive photon data sensitivity to parton distributions
  - data show potential to improve gluon distribution
  - currently limited by scale uncertainty

ATL-PHYS-PUB-2013-018

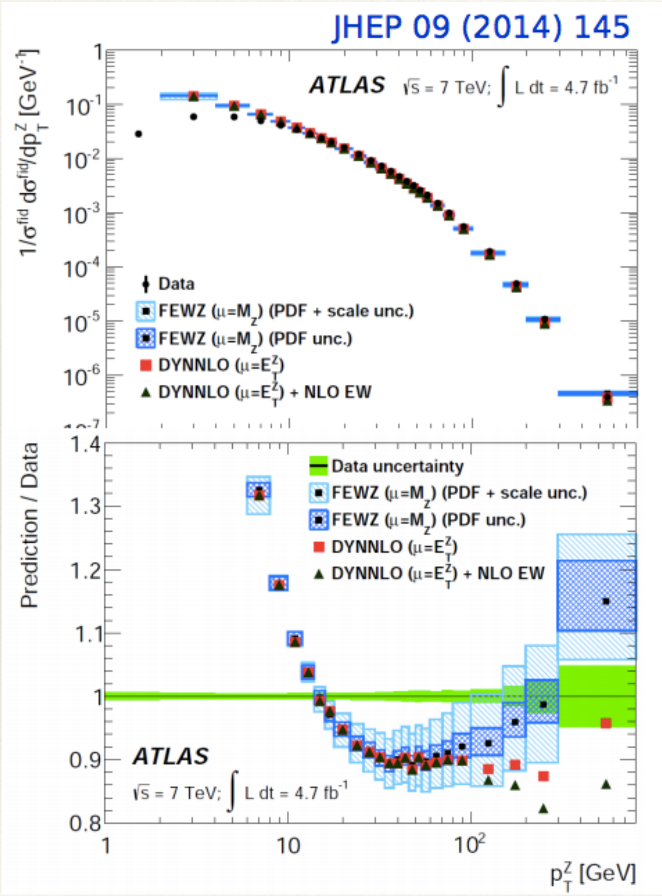


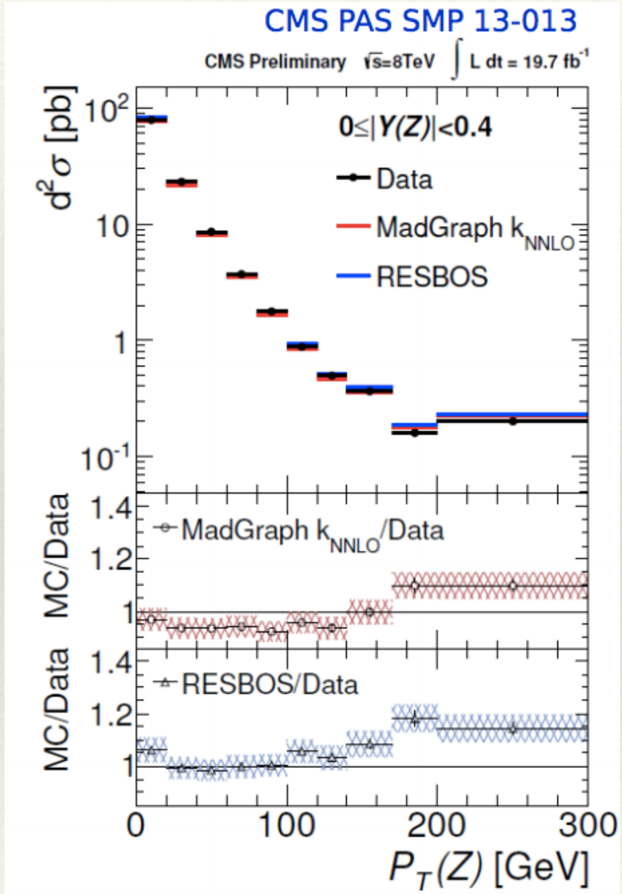


  - large differences observed with theory (NLO) using different PDFs

Vector boson Pt spectrum at the LHC

- ATLAS and CMS both studied the Pt spectrum in rapidity bins
  - low Pt region: dominated by the emission of soft partons (resummation and shower models, fixed order calculations don't work)
  - high Pt region: quark-gluon scattering (PDFs)



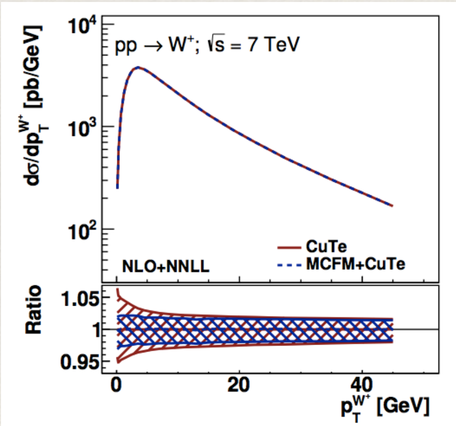


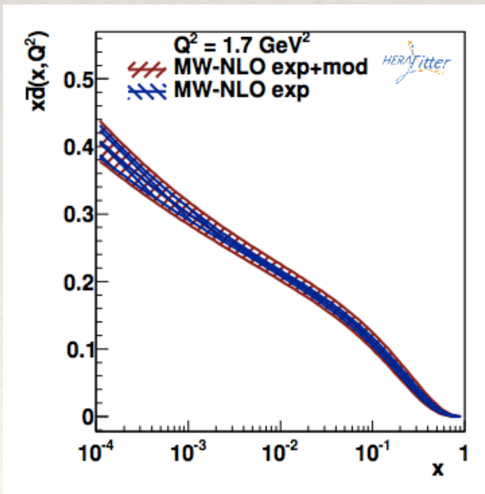
- sensitive data for W mass measurement, PDFs at high x
- currently, limited by precision in theory (needs NNLO and EW corrections)

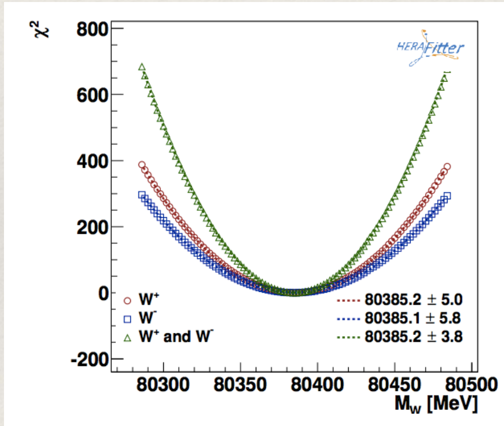
Studies of theoretical uncertainties of Mw mass at the LHC

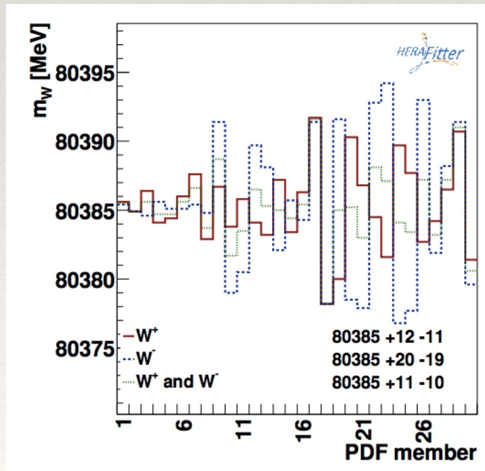
ATL-PHYS-PUB-2014-015

- The measurement of the mass of the W boson provides a stringent test of the SM
- At the LHC, the best experimental precision on Mw might be achieved from the pT distribution of the charged electron/muon from leptonic decay of W:
- A quantitative study of the theoretical uncertainties due to the incomplete knowledge of the quark PDF, and to the uncertainties on the modelling of the low-pT region of W/Z bosons, was performed using HERAFitter platform.
  - Theoretical predictions is based on MCFM and CuTe (interfaced to APPLGRID)
  - A PDF set is generated using simply HERA I data to study the model variations (mc, strange) and propagated via chi2 profiling method to study the effect of PDF uncertainties





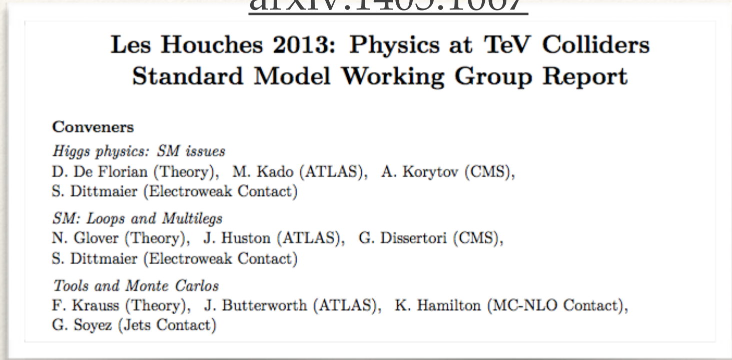




ggH benchmark studies

- Efforts in reducing the PDF uncertainties arising from discrepancy between PDF groups:
  - Benchmark comparisons of NNLO neutral current DIS cross sections (Exercise on HERA-I only data)

arxiv:1405.1067


Les Houches 2013: Physics at TeV Colliders
Standard Model Working Group Report

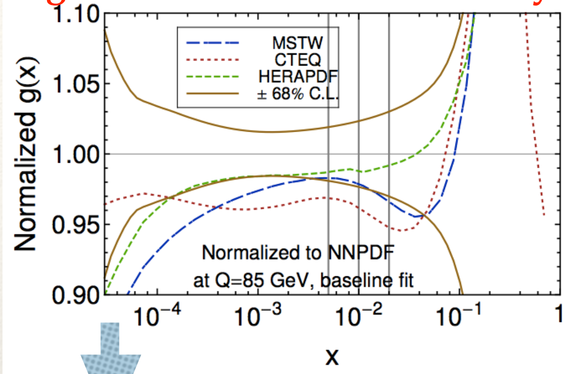
Conveners

Higgs physics: SM issues
D. De Florian (Theory), M. Kado (ATLAS), A. Korytov (CMS),
S. Dittmaier (Electroweak Contact)

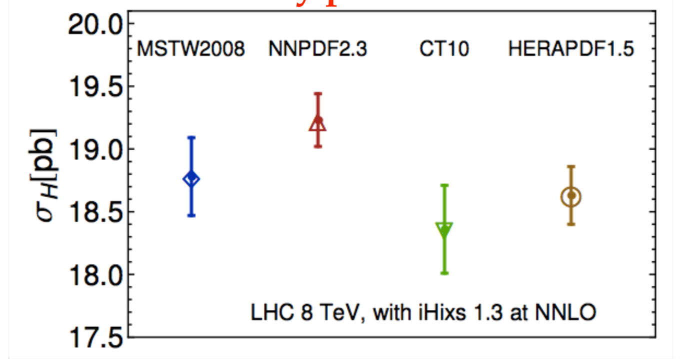
SM: Loops and Multilegs
N. Glover (Theory), J. Huston (ATLAS), G. Dissertori (CMS),
S. Dittmaier (Electroweak Contact)

Tools and Monte Carlos
F. Krauss (Theory), J. Butterworth (ATLAS), K. Hamilton (MC-NLO Contact),
G. Soyez (Jets Contact)

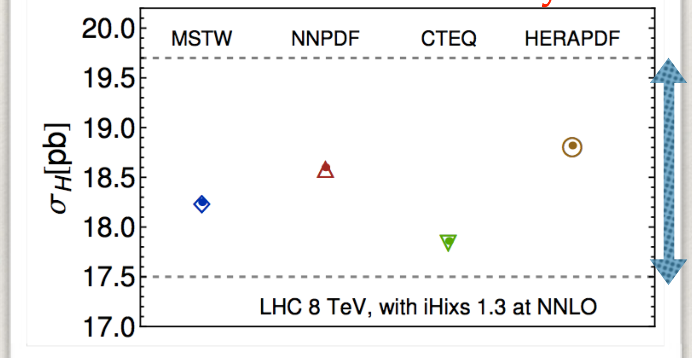
gluon PDF from HERA-I only fits



based on recently published NNLO PDFs.



based on HERA-I only fits



PDF uncert

  - predictions from MSTW, CT, NNPDF and HERAPDF all consistent within PDF uncertainties
  - however the tendency among NNPDF, MSTW and CT is maintained
- **Next step:**
  - continue this exercise by adding additional experimental data sets into the PDF fits sequentially:
  - benchmarking the theoretical predictions used by each group for the different observables -
  - **==> HERAFitter will continue to participate in these studies.**

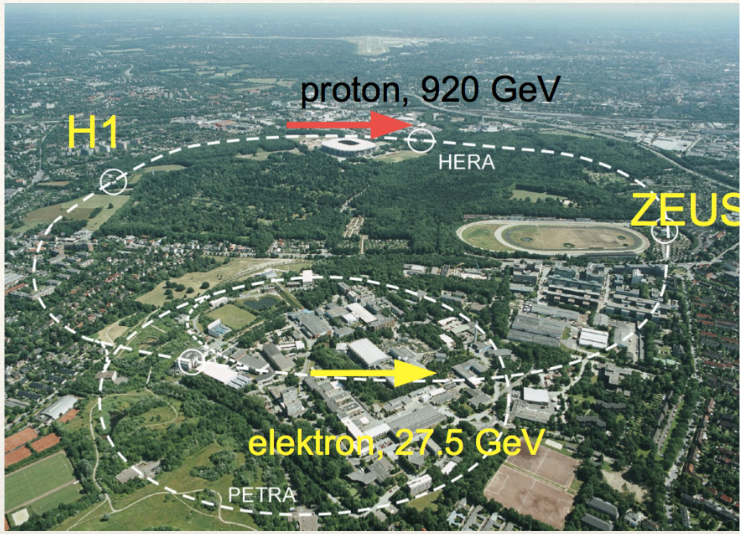
HERAFitter Program at glance

- HERAFitter code is a combination of C++ and Fortran 77 libraries with minimal dependencies and modular structure with interface to external packages:
    - QCDNUM for evolution of PDFs
- DIS inclusive processes in ep and fixed target
    - Different schemes of heavy quark treatment
        - VFNS, FFNS:
            - OPENQCDRAD (ABM)
            - TR' (MSTW)
            - ACOT (CT)
    - Diffractive PDFs
    - Dipole Models
    - Unintegrated PDFs (TMDs)
- Jet production (ep, pp, ppbar)
    - FastNLO and APPLGRID techniques
- Drell-Yan processes (pp, ppbar)
    - LO calculation x NLO k-factors
    - APPLGRID technique
- Top pair production
    - total inclusive ttbar cross sections (HATHOR)
    - differential (DiffTop approx NNLO via fastNLO grids)

```
--enable-openmp         enable openmp support
--enable-trapFPE        Stop of floating point errors (default=no)
--enable-checkBounds    add -fbounds-check flag for compilation (default=no)
--enable-nnpdfWeight    use NNPDF weighting (default=no)
--enable-lhapdf         use lhapdf (default=no)
--enable-applgrid       use applgrid for fast pdf convolutions (default=no)
--enable-genetic        use genetic for general minimia search (defaults=no)
--enable-hathor         use hathor for ttbar cross section predictions
                        (default=no)
--enable-updf           use uPDF evolution (default=no)
--enable-doc            Build documentation (default=no)
```

Experimental Data	Process	Reaction	Theory schemes calculations
HERA, Fixed Target	DIS NC	$ep \to eX$ $\mu p \to \mu X$	TR', ACOT, ZM (QCDNUM), FFN (OPENQCDRAD, QCDNUM), TMD (uPDFevolv)
HERA	DIS CC	$ep \to \nu_e X$	ACOT, ZM (QCDNUM), FFN (OPENQCDRAD)
	DIS jets	$ep \to e$ jets$X$	NLOJet++ (fastNLO)
	DIS heavy quarks	$ep \to ec\bar{c}X$, $ep \to eb\bar{b}X$	TR', ACOT, ZM (QCDNUM), FFN (OPENQCDRAD, QCDNUM)
Tevatron, LHC	Drell-Yan	$pp(\bar{p}) \to l\bar{l}X$, $pp(\bar{p}) \to l\nu X$	MCFM (APPLGRID)
	top pair	$pp(\bar{p}) \to t\bar{t}X$	MCFM (APPLGRID), HATHOR, DiffTop
	single top	$pp(\bar{p}) \to tl\nu X$, $pp(\bar{p}) \to tX$, $pp(\bar{p}) \to tWX$	MCFM (APPLGRID)
	jets	$pp(\bar{p}) \to$ jets$X$	NLOJet++ (APPLGRID), NLOJet++ (fastNLO)
LHC	DY heavy quarks	$pp \to VhX$	MCFM (APPLGRID)

HERA ep collider (1992-2007) @ DESY



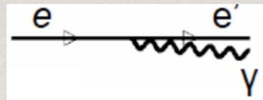
- HERA: unique lepton-proton collider
  - Operational:
    - 1992-2000 (HERA I)
    - 2003-2007 (HERA II)
  - Ep=460-920 GeV, Ee = 27.6GeV
- H1 and ZEUS collected 0.5/fb per experiment
- Rich Physics Program:
  - proton structure, EW, QCD, diffraction, BSM searches,…

Kinematic variables

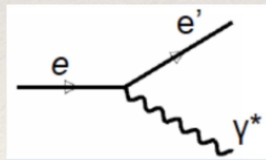
$Q^2 = -q^2 = -(k - k')^2$	Photon virtuality
$x = \dfrac{Q^2}{2p\cdot q}$	Bjorken variable
$y = \dfrac{p\cdot q}{p\cdot k}$	Inelasticity

Two kinematic regimes:

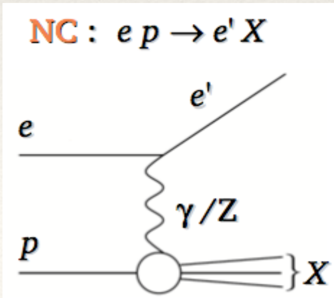
- **Photo-production (PHP)**: $Q^2 < 1\ GeV^2$

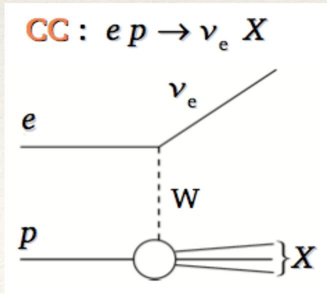


- **Deep Inelastic Scattering (DIS)**: $Q^2 > 1\ GeV^2$



4 processes are available at HERA:





Longitudinal Structure Function

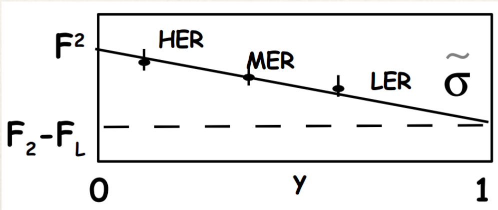
Longitudinal structure function FL is a pure QCD effect:

—> an independent way to probe sensitivity to gluon

$$F_L = \frac{\alpha_s}{4\pi} x^2 \int_x^1 \frac{dz}{z^3}\left[\frac{16}{3} F_2 + 8\sum_q e_q^2 (1-\frac{x}{z}) z g(z)\right]$$

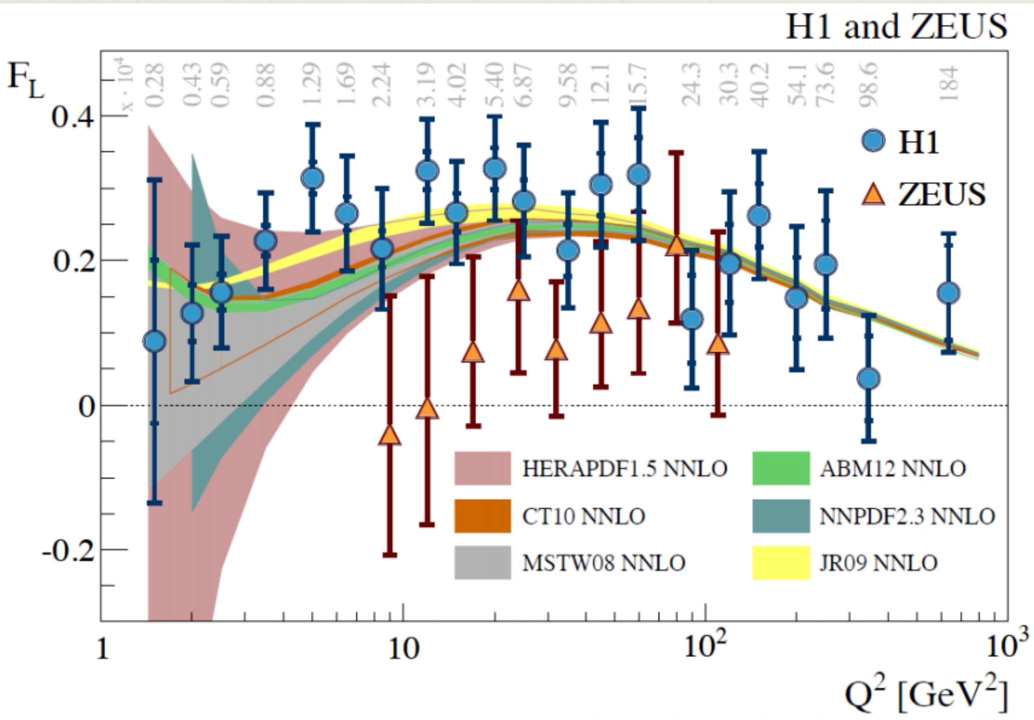
quarks radiating a gluon — gluons splitting into quarks

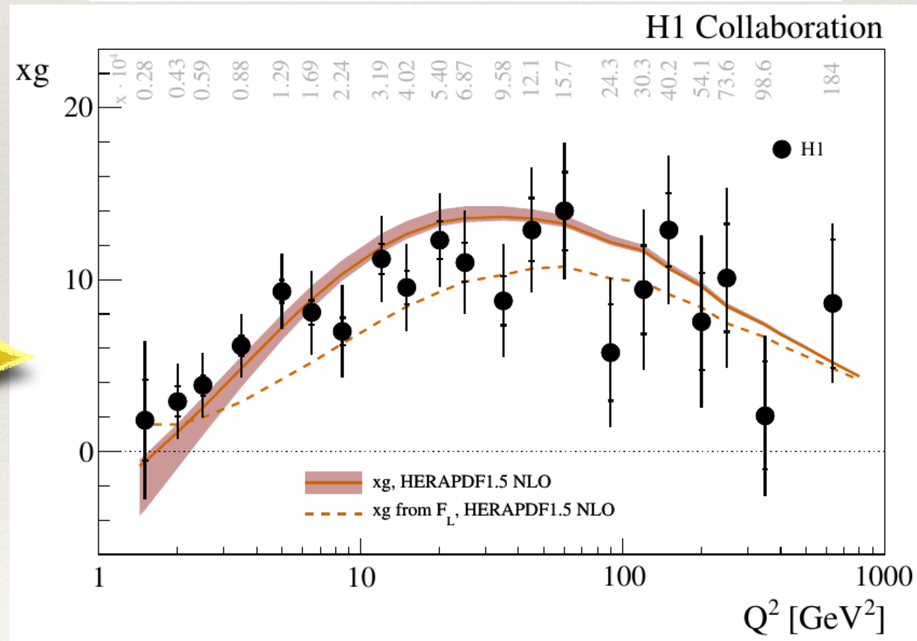
Direct measurement of FL at HERA required differential cross sections at same x and Q^2 but different y —> different beam energies: Ep= 460, 575, 920 GeV



$$\sigma_{NC}(x,Q^2,y) \propto F_2(x,Q^2) - \frac{y^2}{1+(1-y)^2} F_L(x,Q^2)$$

$$xg(x,Q^2) \approx 1.77 \frac{3\pi}{2\alpha_S(Q^2)} F_L(ax,Q^2)$$





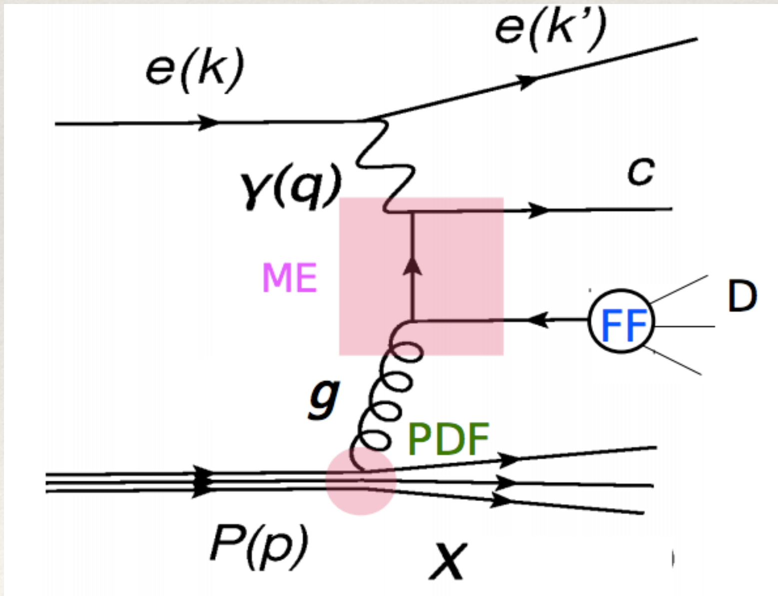
Eur. Phys. J. C 74 (2014) 2814 [arXiv:1312.4821]

DESY-14-053, submitted to Phys. Lett. B [arXiv:1404.6376]

Heavy Flavour Production at HERA

- Heavy Flavour (HF) production: multi-hard scales pose a challenge for pQCD
  - m_c, m_b, p_T, Q^2 —> several calculations (schemes) exist
    - Zero-Mass Variable Flavour Number Scheme (ZMVFNS) — massless scheme
    - Fixed Flavour Number Scheme (FFNS) — massive scheme
    - General-Mass Variable Flavour Number Scheme (GM-VFNS) — matched scheme

- Main process of heavy quark production at HERA is Boson Gluon Fusion



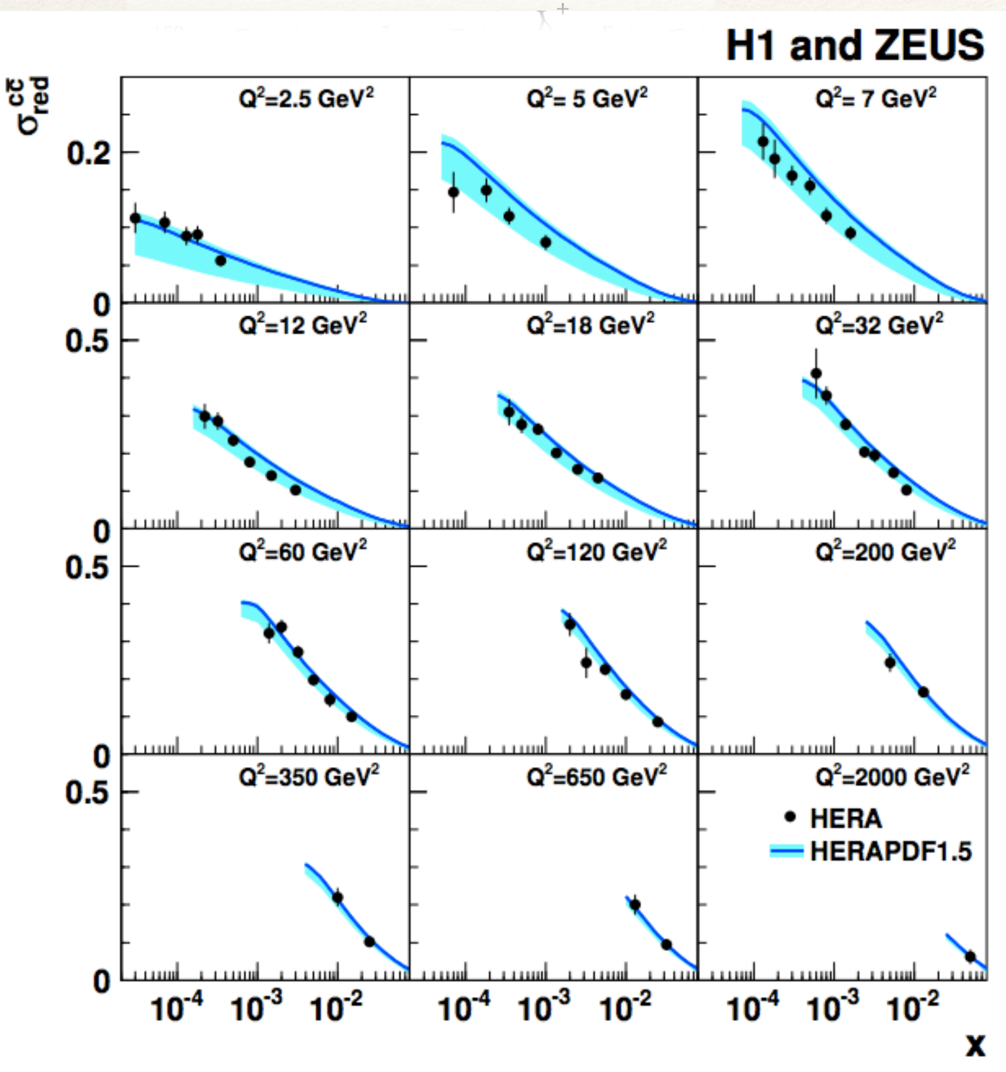
- Measurements of heavy quarks:
  - are sensitive to the gluon PDF
  - are sensitive to the masses of the heavy quarks
  - are sensitive to the fragmentation process of heavy flavour hardons

F2 charm Structure Function

EPJC 73 (2013) 2311

- Rates at HERA in DIS regime σ(b) : σ(c) ≈ O(1%) : O(20%) of σ_{TOT}
- Charm data combination is performed at charm cross sections level:
  - they are obtained from xsec in visible phase space and extrapolated to full space

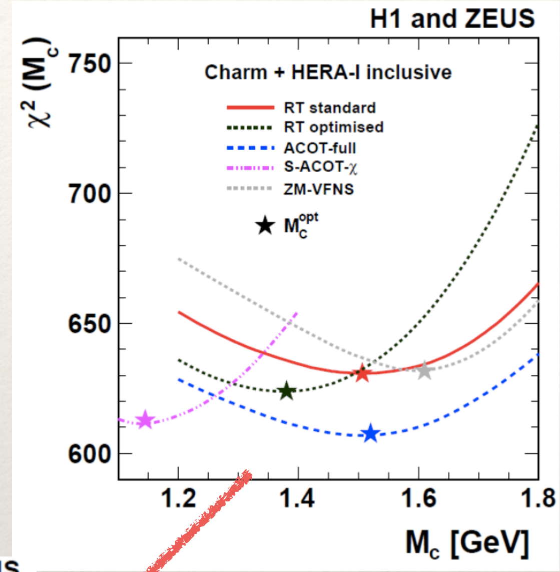
$$\sigma_{red}^{c\bar{c}}(x,Q^2,s) = F_2^{c\bar{c}}(x,Q^2) - \frac{y^2}{Y_+} F_L^{c\bar{c}}(x,Q^2)$$

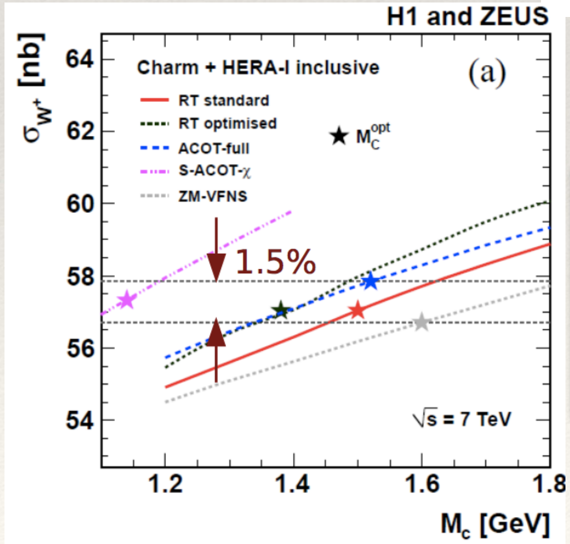


QCD Fits

HERA I+charm

Different calculation schemes prefer different Mc





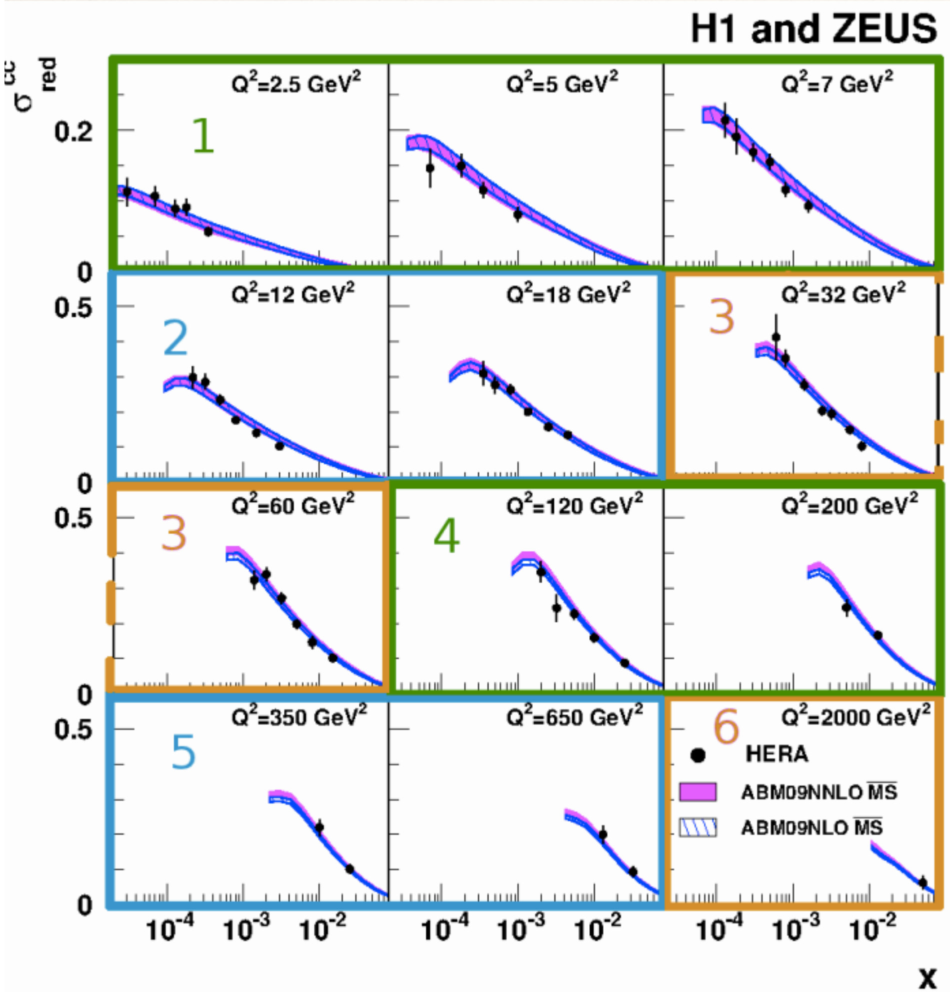
measurements help reduce uncertainties of predictions for the LHC

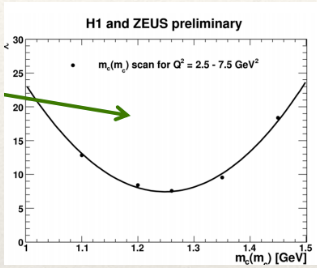
New Measurement of Charm Mass Running

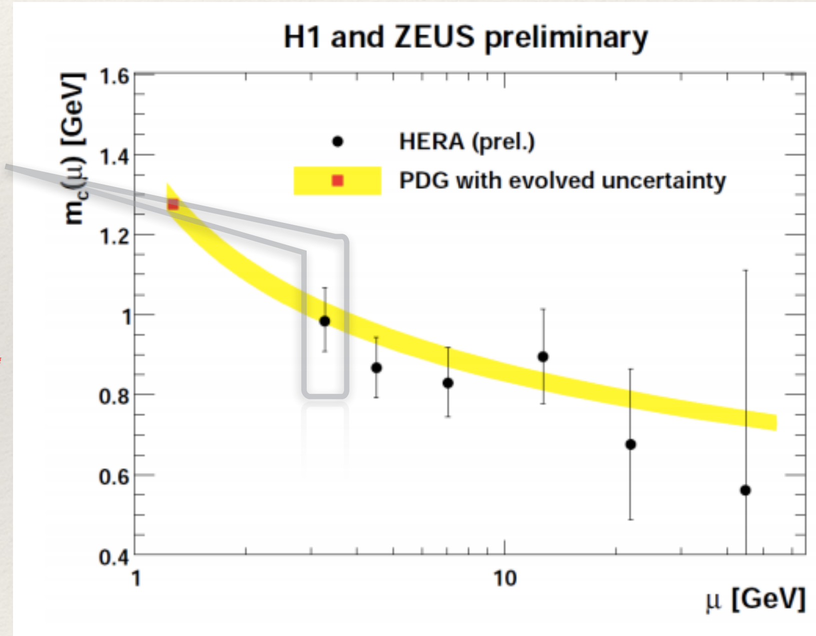
H1-prelim-14-071 ZEUS-prel-14-006 and S. Moch

The running of the charm mass in the $\overline{MS}$ scheme is measured for the first time from the same HERA combined charm data:

- Extract $m_c(m_c)$ in 6 separate kinematic regions
- Translate back to $m_c(\mu)$ [with $\mu=\sqrt{Q^2+4m_c^2}$] using OpenQCDrad [S.Alekhin's code].



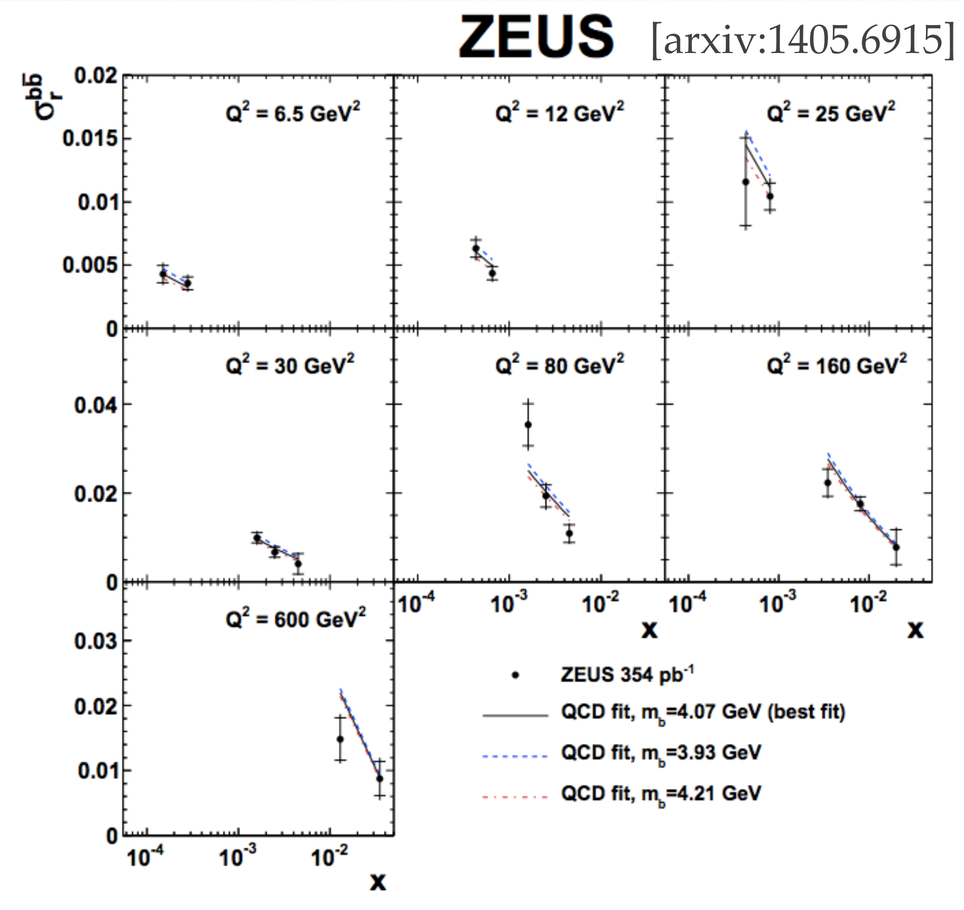




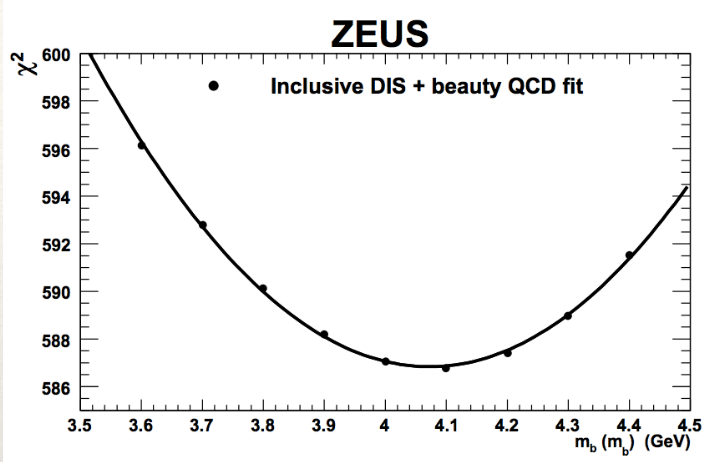
The scale dependence of the mass is consistent with QCD expectations

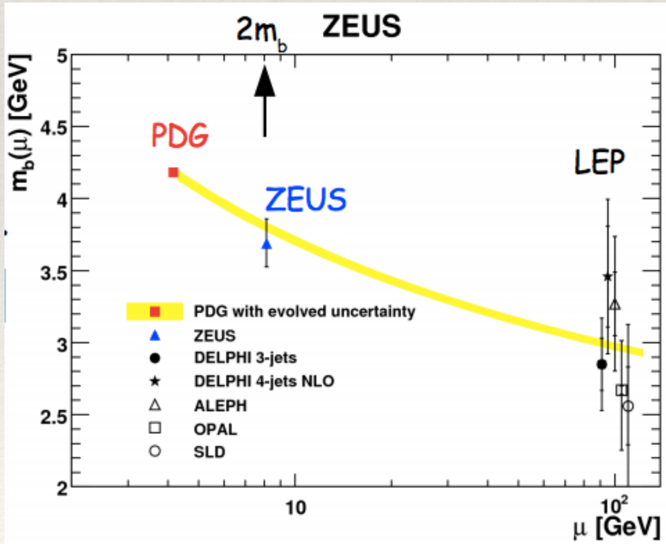
Running beauty mass from F2b

- The value of the running beauty mass is obtained using HERAFitter (via OPENQCDRAD):
  - chi2 scan method from QCD fits in FFN scheme to the combined HERA I inclusive data + beauty measurements, beauty-quark mass is defined in the $\overline{\text{MS}}$ scheme.



QCD Fits
HERA I+beauty





The extracted $\overline{\text{MS}}$ beauty-quark mass is in agreement with PDG average and LEP results.

Transverse Momentum Distributions

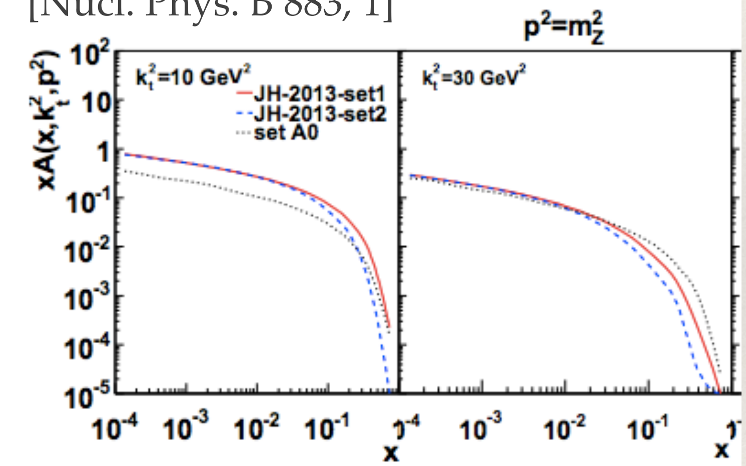
- QCD applications to multiple-scale scattering problems and complex final-state observables require in general formulations of factorisation which involve transverse-momentum dependent (TMD) - or known also as unintegrated PDFs.

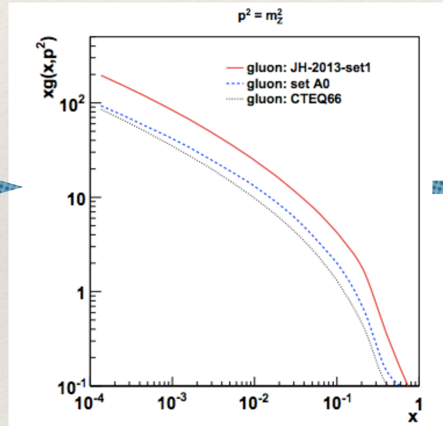
$$\sigma_j(x,Q^2) = \int_x^1 dz \int d^2k_t \; \hat{\sigma}_j(x,Q^2,z,k_t) \; \mathcal{A}(z,k_t,\mu)$$

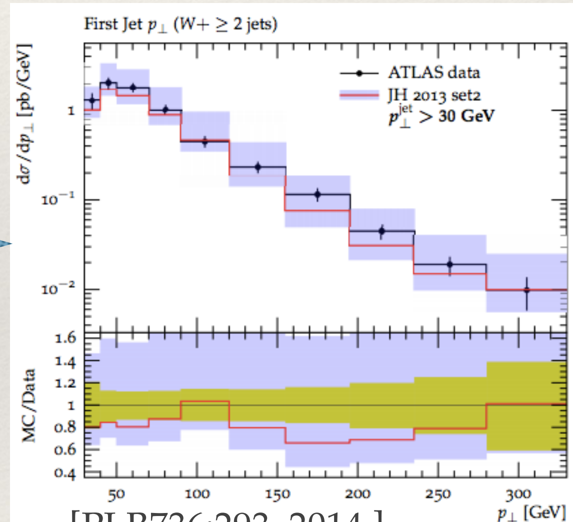
→ a convolution in both longitudinal and transverse momenta of TMD with off-shell partonic matrix elements

- Fits to combined measurements of proton's structure functions from HERA using transverse momentum dependent QCD factorisation and CCFM evolution is performed using HERAFitter platform

[Nucl. Phys. B 883, 1]







[PLB736:293, 2014.]

- The extracted gluon TMD with experimental and theory uncertainty [JH-2013-set1] is then used as prediction to vector boson+jet production process at the LHC [Phys. Rev. D 85 (2012) 092002.]
  - This process is important both for SM physics and for new physics searches at the LHC
  - Results compare well with the measurements of jet multiplicities and transverse momentum spectra within the pdf uncertainties

QCD Settings for HERAPDF2.0

The QCD settings are optimised for HERA measurements of proton structure functions:

PDFs are parametrised at the starting scale Q_0^2=1.9 GeV^2 as follows:

$$
\begin{aligned}
xg(x) &= A_g x^{B_g}(1-x)^{C_g} - A'_g x^{B'_g}(1-x)^{C'_g}, \\
xu_v(x) &= A_{u_v} x^{B_{u_v}}(1-x)^{C_{u_v}}\left(1 + D_{u_v}x + E_{u_v}x^2\right), \\
xd_v(x) &= A_{d_v} x^{B_{d_v}}(1-x)^{C_{d_v}}, \\
x\bar{U}(x) &= A_{\bar{U}} x^{B_{\bar{U}}}(1-x)^{C_{\bar{U}}}(1 + D_{\bar{U}}x), \\
x\bar{D}(x) &= A_{\bar{D}} x^{B_{\bar{D}}}(1-x)^{C_{\bar{D}}}.
\end{aligned}
$$

fixed or constrained by sum-rules

parameters set equal but free

NC structure functions

$$F_2 = \frac{4}{9}\left(xU + x\bar{U}\right) + \frac{1}{9}\left(xD + x\bar{D}\right)$$

$$xF_3 \sim xu_v + xd_v$$

CC structure functions

$$W_2^- = x(U + \overline{D}), \qquad W_2^+ = x(\overline{U} + D)$$

$$xW_3^- = x(U - \overline{D}), \qquad xW_3^+ = x(D - \overline{U})$$

Due to increased precision of data, more flexibility in functional form is allowed —> 15 free parameters

- PDFs are evolved via evolution equations (DGLAP) to NLO and NNLO (as(MZ)=0.118)
- Thorne-Roberts GM-VFNS for heavy quark coefficient functions – as used in MSTW
- Chi2 definition used in the minimisation [MINUIT] accounts for correlated uncertainties:

$$\chi^2_{tot}(\mathbf{m},\mathbf{b}) = \sum_i \frac{\left[\mu^i - m^i\left(1 - \sum_j \gamma^i_j b_j\right)\right]^2}{\delta^2_{i,stat}\mu^i m^i\left(1 - \sum_j \gamma^i_j b_j\right) + \left(\delta_{i,unc} m^i\right)^2} + \sum_j b_j^2 + \sum_i \ln \frac{\delta^2_{i,unc} m_i^2 + \delta^2_{i,stat}\mu^i m^i}{\delta^2_{i,unc}\mu_i^2 + \delta^2_{i,stat}\mu_i^2}$$

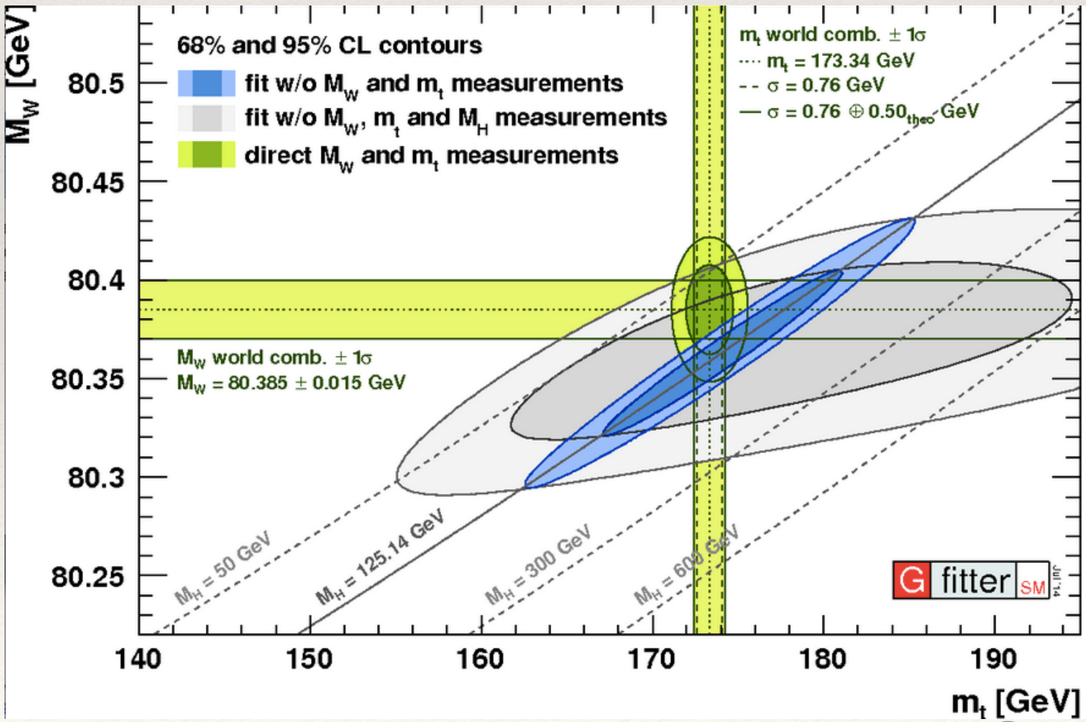
Why do we still need to care about PDFs?

- Discovery of new exciting physics relies on precise knowledge of proton structure.
- **Factorisation theorem:**
    - Cross section can be calculated by convoluting short distance partonic reactions (calculable in pQCD) with Parton Distribution Functions (PDFs):

$$d\sigma(h_1h_2 \to cd) = \int_0^1 dx_1 dx_2 \sum_{a,b} f_{a/h_1}(x_1,\mu_F^2) f_{b/h_2}(x_2,\mu_F^2) d\hat{\sigma}^{(ab\to cd)}(Q^2,\mu_F^2)$$

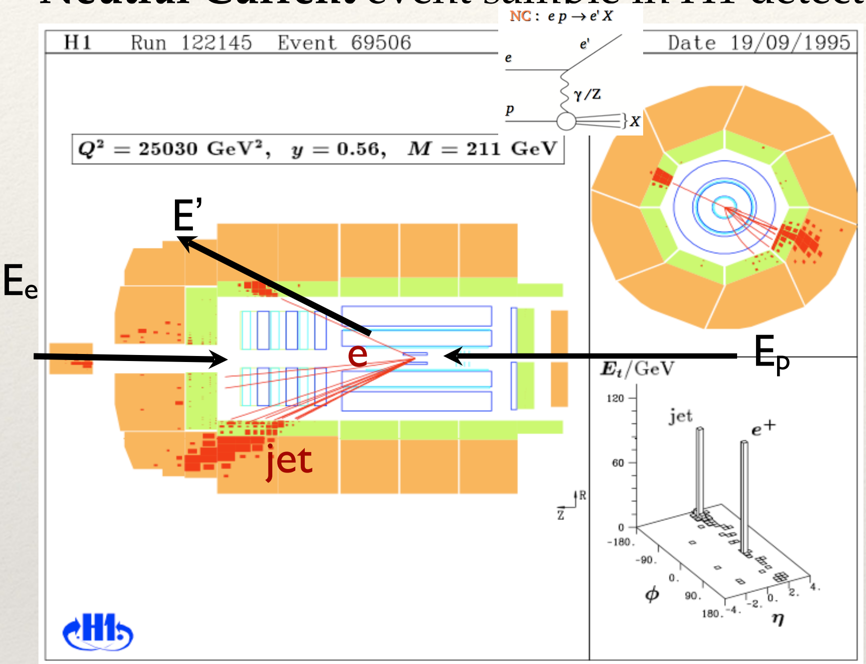
    - PDFs cannot be calculated in perturbative QCD, however they are process independent (universal) and their evolution with the scale is predicted by pQCD

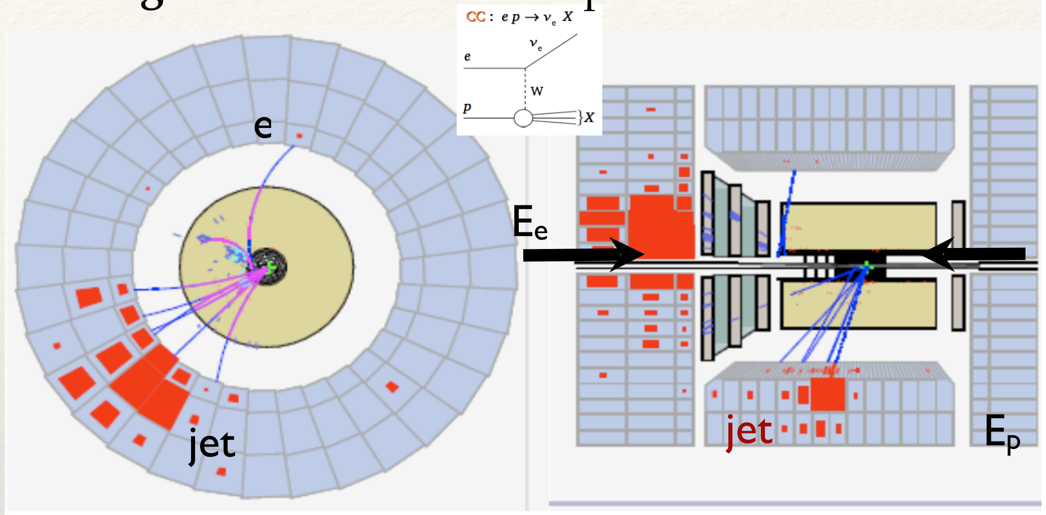
- PDFs are one of the main theory uncertainties in Mw measurement
- PDFs are one of main theory uncertainties in Higgs production:
  ~ 7-8 % PDF and αs and ~ 7-8 % scale



HERA ep collider (1992-2007) @ DESY

- H1 and ZEUS experiments at HERA collected ~1/fb of data
  - Ep=460/575/820/920 GeV and Ee=27.5 GeV
- **Neutral Current** event sample in H1 detector
- **Charged Current** event sample in ZEUS detector





$$\frac{d\sigma^{\pm}_{NC}}{dxdQ^2} = \frac{2\pi\alpha^2}{x}\left[\frac{1}{Q^2}\right]^2\left[Y_+\tilde{F}_2 \mp Y_- x\tilde{F}_3 - y^2\tilde{F}_L\right]$$

$$\frac{d\sigma^{\pm}_{CC}}{dxdQ^2} = \frac{G_F^2}{4\pi x}\left[\frac{M_W^2}{M_W^2+Q^2}\right]^2\left[Y_+\tilde{W}_2^{\pm} \mp Y_- x\tilde{W}_3^{\pm} - y^2\tilde{W}_L^{\pm}\right]$$

$$Y_{\pm} = 1 \pm (1-y)^2$$

Determination of the Event Kinematics:

- using lepton information (E_e', θ_e)
- **using hadronic final state particles**
- using both lepton and hadronic final state variables

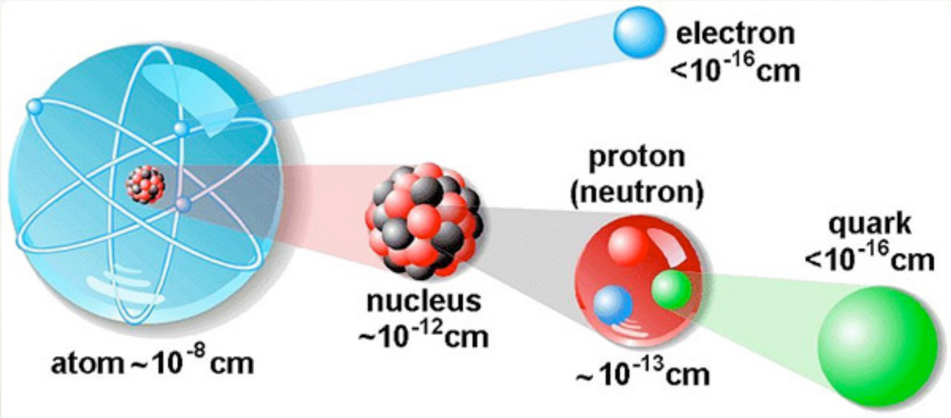
$$s = 4E_eE_p$$
$$Q^2 = E_eE'(1+\cos\theta_e)$$
$$y = 1 - \frac{E'}{E_e}\frac{1}{2}(1-\cos\theta_e)$$
$$x = \frac{Q^2}{sy}$$

$$\tilde{F}_2 \propto \sum(xq_i + x\bar{q}_i)$$
dominant contribution (all Q2 plane)

$$x\tilde{F}_3 \propto \sum(xq_i - x\bar{q}_i)$$
significant contributions at high Q2

$$\tilde{F}_L \propto \alpha_s \cdot xg(x,Q^2)$$
high y

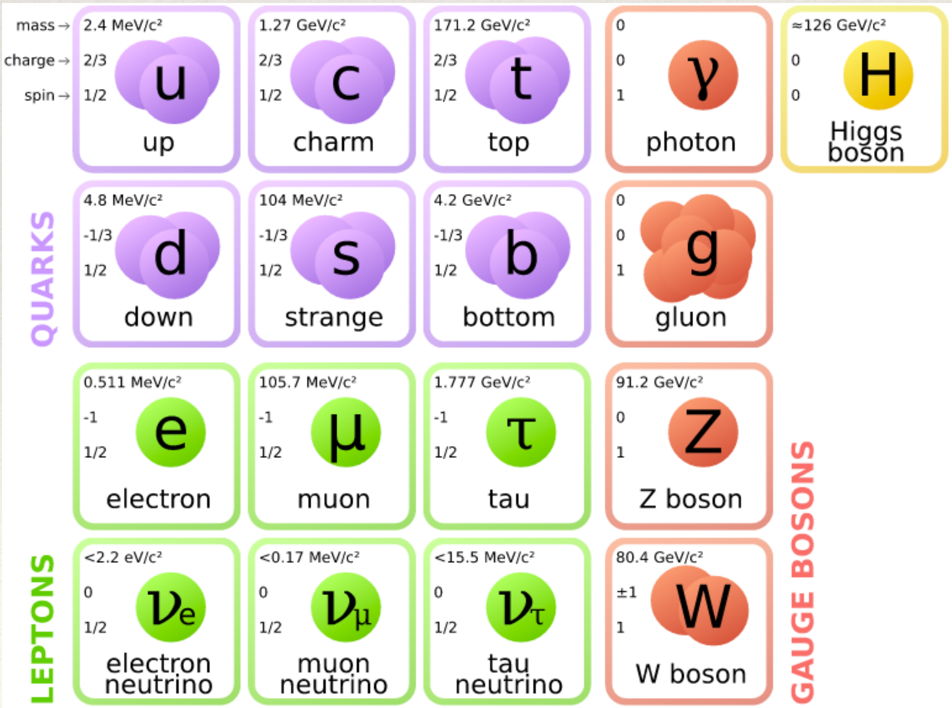
Buildings Blocks of the Standard Model



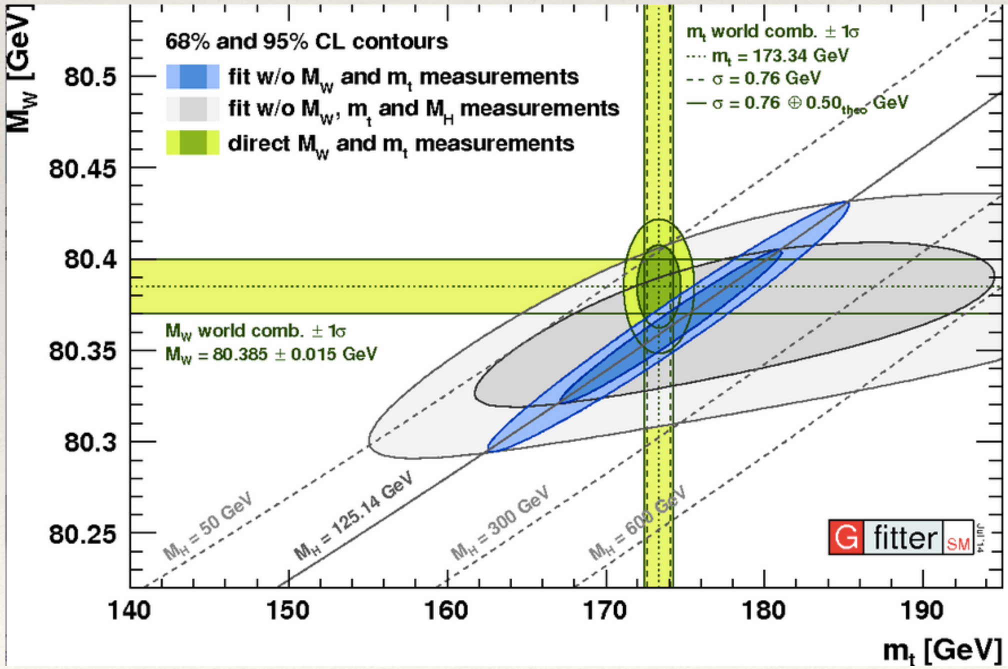
Road from nuclear physics to today's physics:

- Hadron spectroscopy: quark model
- Deep Inelastic Scattering experiments: reality of quarks
- Tests of QCD theory

Building Blocks of matter:

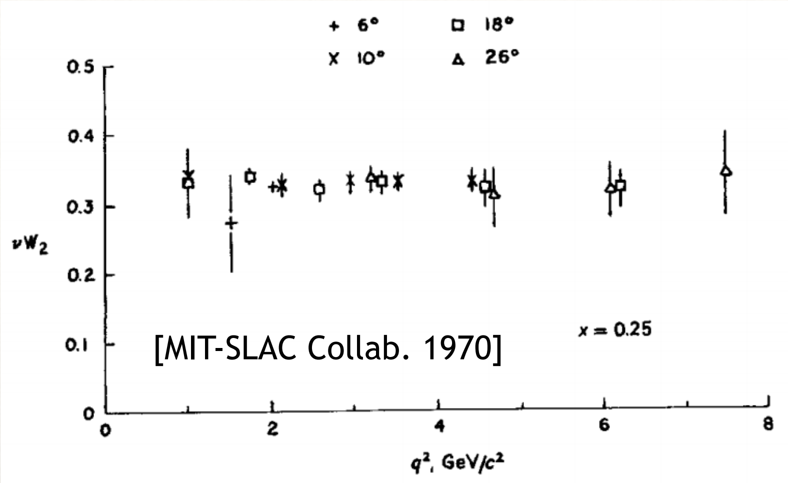


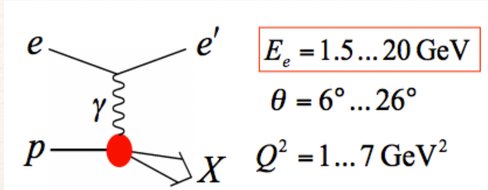
Cornering Standard Model to look for first signs of new physics



Parton Distribution Functions

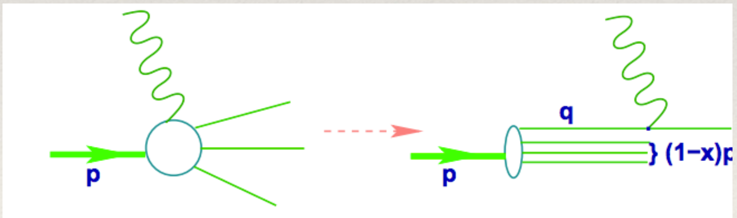
- Parton Model was introduced by Feynman (1969) to explain the Bjorken scaling:





Concept of Scaling:
if proton is made up from point-like particles, then the cross section becomes approximately independent on the scale;
[as we know, later with higher resolving power scaling violations were proven]

- inelastic scattering with nucleon is viewed as elastic scattering between lepton and a point-like constituent of the target – **partons** (non-interacting) – explicitly assumed to be spin-1/2 particles

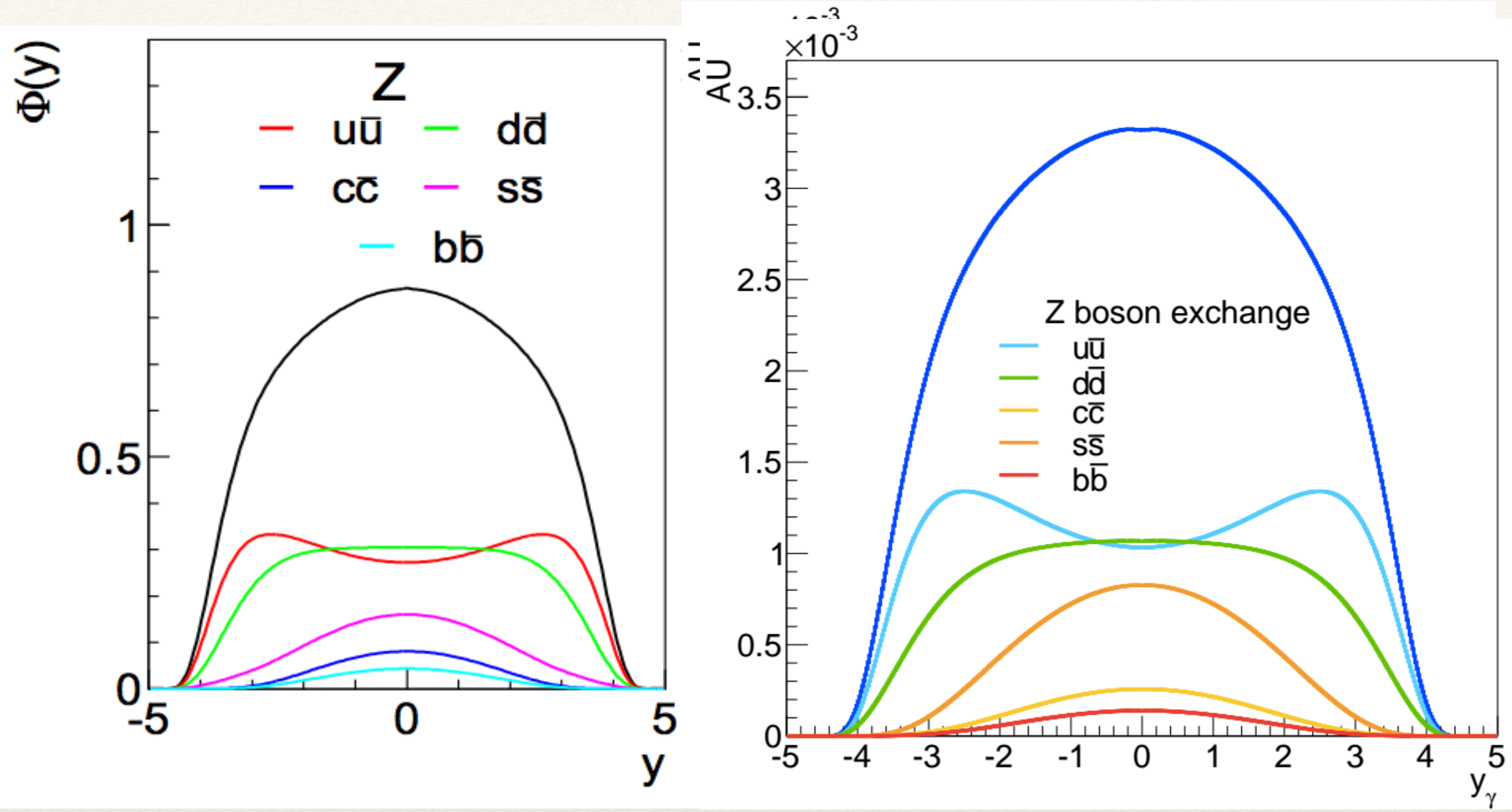


- Parton Distribution Functions (PDFs): Each parton carries the fraction x with a probability q(x)

$$\left(\frac{d\sigma}{dxdQ^2}\right)_{ep\to eX} = \sum_i \int dx e_i^2 q_i(x) \left(\frac{d\sigma}{dxdQ^2}\right)_{eq_i\to eq_i}.$$

Enhanced Strange

- CT10 vs ATLAS_epWZ



DIS Cross Sections

- Differential cross section is experimentally measured: theory meets the experiment

- Factorisable nature of interaction: Inclusive scattering cross section is a product of leptonic and hadronic tensors times propagator characteristic of the exchanged particle:

$$\frac{d^2\sigma}{dxdQ^2} = \frac{2\pi\alpha^2}{Q^4 x}\sum_j \eta_j L_j^{\mu\nu} W_j^{\mu\nu}$$

For NC: j=γ, Z, γZ
For CC: j=W+, W-

$$\eta_\gamma = 1; \qquad \eta_{\gamma Z} = \left(\frac{G_F M_Z^2}{2\sqrt{2}\pi\alpha}\right)\left(\frac{Q^2}{Q^2+M_Z^2}\right); \qquad \eta_Z = \eta_{\gamma Z}^2;$$

$$\eta_W = \frac{1}{2}\left(\frac{G_F M_W^2}{4\pi\alpha}\frac{Q^2}{Q^2+M_W^2}\right)^2,$$

Leptonic tensor: related to the coupling of the lepton with the exchanged boson

- contains the electromagnetic or the weak couplings
- can be calculated exactly in the standard electroweak U(1) × SU(2) theory.

Hadronic tensor: related to the interaction of the exchanged boson with proton

- can't be calculated, but only be reduced to a sum of structure functions:

~m_{lepton}

$$W^{\alpha\beta} = -g^{\alpha\beta}W_1 + \frac{p^\alpha p^\beta}{M^2}W_2 - \frac{i\epsilon^{\alpha\beta\gamma\delta}p_\gamma q_\delta}{2M^2}W_3 + \frac{q^\alpha q^\beta}{M^2}W_4 + \frac{p^\alpha q^\beta + p^\beta q^\alpha}{M^2}W_5 + \frac{i(p^\alpha q^\beta - p^\beta q^\alpha)}{2M^2}W_6$$

$$\frac{d^2\sigma}{dxdQ^2} = A^i\left\{(1 - y - \frac{x^2y^2M^2}{Q^2})F_2^i + y^2xF_1^i \mp (y - \frac{y^2}{2})xF_3^i\right\}$$

A^i: process dependent

DIS Cross Sections

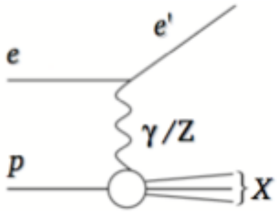
General Form for the Differential cross section:

$$\frac{d^2\sigma}{dxdQ^2} = A^i \left\{ (1 - y - \frac{x^2y^2M^2}{Q^2}) F_2^i + y^2 x F_1^i \mp (y - \frac{y^2}{2}) x F_3^i \right\}$$

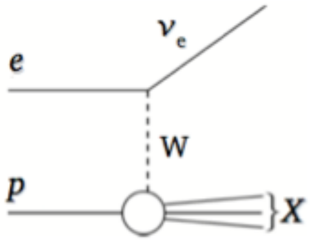
$$xF_3 \sim \sum (xq_i - x\bar{q}_i)$$
$$F_L \sim \alpha_S g$$
$$F_2 \sim \sum e_i^2 (xq_i + x\bar{q}_i)$$

A^i: process dependent

NC : $e\,p \to e'\,X$



CC : $e\,p \to \nu_e\,X$



NC:

$$\frac{d^2\sigma_{NC}^{\pm}}{dxdQ^2} = \frac{2\pi\alpha^2}{x}\left[\frac{1}{Q^2}\right]^2 \phi_{NC}^{\pm}(x, Q^2)$$

$$\phi_{NC} = Y_+ \tilde{F}_2^{\pm}(x,Q^2) - y^2 \tilde{F}_L^{\pm}(x,Q^2) \mp Y_- x\tilde{F}_3^{\pm}(x,Q^2),$$

$$\tilde{F}_2^{\pm} = F_2 - (v_e \pm P_e a_e)\kappa_Z F_2^{\gamma Z} + (v_e^2 + a_e^2 \pm 2P_e v_e a_e)\kappa_Z^2 F_2^Z, \quad \kappa_Z(Q^2) = \frac{1}{4sin^2(\theta_W)cos^2(\theta_W)} \frac{Q^2}{Q^2 + M_Z^2}$$

$$x\tilde{F}_3^{\pm} = -(a_e \pm P_e v_e)\kappa_Z xF_3^{\gamma Z} + (2v_e a_e \pm P_e(v_e^2 + a_e^2))\kappa_Z^2 xF_3^Z,$$

CC:

$$\frac{d^2\sigma_{CC}^{\pm}}{dxdQ^2} = (1 \pm P_e)\frac{G_F^2}{2\pi x}\left[\frac{M_W^2}{Q^2 + M_W^2}\right]^2 \phi_{CC}^{\pm}(x, Q^2).$$

At LO

$$e^+ : \quad \phi_{CC}^{+} = x[(\bar{u}(x) + \bar{c}(x)) + (1-y)^2(d(x) + s(x))].$$

$$e^- : \quad \phi_{CC}^{-} = x[(u(x) + c(x)) + (1-y)^2(\bar{d}(x) + \bar{s}(x))]$$